



THEORETICAL STUDIES OF SINGLE-SITE CATALYSTS FOR EFFICIENT ELECTROCHEMICAL CO₂ REDUCTION

Paulina Pršlja

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Paulina Pršlja

Theoretical Studies of
Single-Site Catalysts
for Efficient
Electrochemical CO₂ Reduction

DOCTORAL THESIS

Supervised by
Prof. Núria López Alonso

Institute of Chemical Research of Catalonia (ICIQ)
and Rovira i Virgili University (URV)



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I STATE that the present study, entitled “**Efficient Electrochemical CO₂ Reduction on Single Site Catalysts**”, presented by Paulina Pršlja for the award of the degree of Doctor in Chemical Science and Technology, has been carried out under our joint supervision at the Institute of Chemical Research of Catalonia and that it fulfills all the requirements to be eligible for the International Doctorate Award.

Tarragona, November 18th, 2020



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**”Nothing in life is to be feared
it is only to be understood.
Now is the time to understand more,
so that we may fear less.”**

Marie Skłodowska-Curie

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**”Every brilliant experiment,
like every great work of art,
starts with an act of imagination.”**

Jonah Lehrer

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List of publications

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Paulina Pršlja was responsible for the conception of idea design, carried out all calculations and formal analyses, wrote most of the text of all manuscripts, and made all computational figures.

Abbreviations

GHG	Greenhouse Gasses
CCUS	Carbon dioxide Capture, Utilization, and Storage
CCS	Carbon dioxide Capture, and Storage
CCU	Carbon dioxide Capture, and Utilization
EOR	Enhanced Oil Recovery
eCO ₂ RR	Electrochemical CO ₂ Reduction Reaction
HER	Hydrogen Evolution Reaction
OER	Oxygen Evolution Reaction
ORR	Oxygen Reduction Reaction
GTL	Gas To Liquid
CTL	Coal To Liquid
SHE	Standard Hydrogen Electrode
RHE	Reversible Hydrogen Electrode
CHE	Computational Hydrogen Electrode
DFT	Density Functional Theory
LDA	Local-Density Approximation
GGA	Generalized Gradient Approximation
VASP	Vienna Ab Initio Simulation Package
PAW	Projector-augmented wave pseudopotentials
PBE	Perdew-Burke-Ernzerhof density functional
RPBE	Revised Perdew-Burke-Ernzerhof density functional
RDS	Rate Determining Step
HK	Hohenberg and Kohn
SAC	Single-Atom-Catalysis
MOF	Metal-Organic Framework
M-N-C	Metal-Nitrogen-Carbon catalyst
CoTPP	Cobalt(II)-TetraPhenyl Porphine
CoPc	Cobalt(II)-Phthalocyanine
MWCNT	MultiWalled Carbon NanoTube
CNT	Carbon NanoTube
MO	Molecular Orbital

M-N-C	Metal-Nitrogen-Carbon catalyst
EXAFS	Extended X-ray Absorption Fine Structure
XANES	X-ray Absorption Near Edge Structure
XPS	X-ray Photoelectron Spectroscopy
XAS	X-ray Absorption Spectroscopy
HRTEM	High-Resolution Transmission Electron Microscopy
SEM	Scanning Electron Microscopy
¹ H NMR	Proton Nuclear Magnetic Resonance
HPLC	High-Performance Liquid Chromatography
GC	Gas Chromatography
J _P FE	Faradaic Efficiency
NPs	Nanoparticles

Abstract

Global climate change caused by anthropogenic emission of CO₂ is a growing threat to society. Clean CO₂ electrochemical conversion (eCO₂RR) to produce value-added chemicals and fuels using renewable electricity is a promising strategy for closing the carbon cycle. However, this early-stage technology is challenged by the low activity and selectivity and the competing hydrogen evolution reaction (HER). Thus, it is important to design catalysts with economically viable current density at low overpotential and high selectivity toward certain products. Single-site carbon-based materials are alternative candidates to typical metal catalysts for the CO₂RR. This thesis aims to model carbon materials under electrochemical conditions using Density Functional Theory (DFT) for CO₂ reduction. DFT can be applied to systems under electrochemical conditions coupled with Computational Hydrogen Electrode (CHE) formalism to investigate their selectivity and activity. The main goal is to establish experimental and theoretical correlations between physicochemical and catalytic properties for the eCO₂RR over different types of catalysts. One of the promising single-site catalysts is metal-nitrogen-carbon (MNC) material with metal atom present as atomically dispersed metal-N_x centers and it was investigated as a model catalyst. The distinct activity for CO formation (J_{CO}) observed along the series of catalysts is attributed to the nature of the transition metal in MN_x moieties. Furthermore, I identified the volcano trend between their activity toward CO formation and the CO binding energy for different MN_x sites, with Fe and Co at the top of the volcano. Among the atomically dispersed moieties the Nickel Nitrogen Carbon (NiNC) single-atom catalyst exhibits the highest efficiency for producing CO at different potentials. The variation in the material synthesis produces defects with coordinatively saturated and unsaturated N-doped cavities, and once the metal is placed there these single atoms can present different metal oxidation states depending on the cavity nature. I have computationally evaluated the activity and selectivity of a wide potentially active sites for single atoms and nanoparticles, where the best performing reactivity is found for N-doped models. In addition, the electrochemical stability of the reconstruction and redispersion of supported nanoparticles is addressed. To dissect the possible disintegration of Ni nanoparticles by Ni(CO)₂-like species Ostwald formalism was applied. I developed process of Ni nanoparticles reconstruction by forming Ni(CO)₄ species that can redisperse into active single atoms at high CO coverages. Moving to different systems, I also unravelled the production of

C₁ hydrocarbons with molecular catalysts. I focused on the potential and temperature influence of the product selectivity of CO₂ reduction reaction over Cobalt(II)-tetraphenylporphine and Co(II)-phthalocyanine deposited on multiwalled carbon nanotubes and carbon nanotubes. I identified the key differences between the bonding energies of CH₂O via O atom or CHOH via C atom, the binding strength of CO, and the protonation of *CHO. Taken together, this thesis presents a computational understanding of activity and selectivity of single-site carbon-based catalysts for CO₂ reduction.

Chapter 1

Introduction

1.1 Fossil fuels dependency

Fossil fuels (coal, oil, and natural gas) have been the driving force behind the industrialized world and its economic growth. Such energy has grown from insignificant levels in 1800 to an annual output of nearly 10,000 million tons of oil equivalents Mtoe. At present, about 80% of all primary energy in the world is derived from fossil fuels with oil accounting for 32.8%, coal for 27.2%, and natural gas for 20.9%. Combustible biomass and waste (10.2%), nuclear power (5.8%), and hydroelectric dams (2.3%) are the largest contributors to the global energy system after fossil based energy, but they account for only a minor share of the global primary energy supply.¹ Fossil fuels have powered our world for 250 years, they are inexpensive and accessible energy.² There are plenty of coalfields and petrol stations, besides they are still used today to make the majority of plastic products.³ However, such non-renewable sources lead to energy crises and concomitant environmental issues:

- Acid rain is formed by a chemical reaction between sulfur dioxide (SO₂) and nitrogen oxides (NO_x) combined with water, oxygen (O₂), and other chemicals high in the atmosphere. Most of the SO₂ originates from coal combustion, while NO_x can be derived from natural gas and motor fuels.^{4,5} As CO₂ dissolves into carbonic acid, pH of rain is about 5.6, while Acid rain has a pH between 4.2 and 4.4.
- Combustion of coal also enhance the emission of Mercury (Hg), which is toxic and extremely damaging to the environment. The main driving force for increased emissions is the expansion of coal-fired electricity generation in the developing world. From 2008 to 2050 elemental Hg emission will only reduce from 65% to 50-55%.⁶

- The major cause of greenhouse gasses (GHG) emission into the atmosphere is global warming or climate change.⁷ The burning of coal is thought to contribute 44% of the world's carbon dioxide emissions. Around the world, petrol (gasoline) alone is said to be responsible for a third of the carbon emissions. Though cleaner than both coal and crude oil, natural gas is responsible for around 20% of our carbon emissions. Since 2015 400ppm of CO₂ was released into the atmosphere.⁸

Only 0.8% of the world's primary energy is derived from geothermal, wind, solar, or other alternative energy sources. More specifically, wind power accounted for only 0.2% of the global primary energy supply with its 23 Mtoe contribution while direct solar energy accounted for 0.1% with a 12 Mtoe output.¹ Making the switch to clean energy sources will help us to reverse the effects of global warming.

1.2 CO₂ capture, utilization and storage

CO₂ is a greenhouse gas that originates as waste from the burning of fossil fuels, the production of electricity, fertilizers, chemicals, steel, and cement.

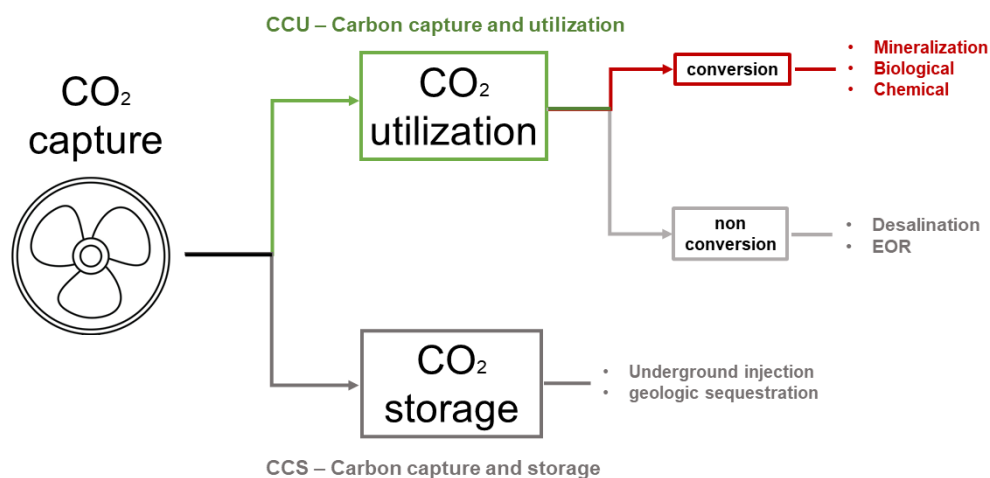


Figure 1.1: The schematic representation of carbon capture, utilization and storage (CCUS) concept, adapted from Ref.⁹

Carbon dioxide capture, utilization, and storage (CCUS) is a challenging demanding issue, see scheme in the Figure 1.1. The captured CO₂ is diluted at about 95% purity, therefore it can be injected underground where the earth can store it. The earth has around 300 billion tons of storage capacity in oil and gas reservoirs, un-mineable coal seams, and deep saline reservoirs. Other than storing CO₂ it can also be utilized, where it can be divided into two parts. First, CO₂ alone without any conversion has certain uses, such as, enhanced oil recovery (EOR) by CO₂ flooding.¹⁰ Indeed, the injection of CO₂ into an oil reservoir increases the production due to the high mutual dissolving

capability of supercritical CO₂ and hydrophobicity of oil. Also, the increase in pressure lowers the viscosity of the CO₂ oil mixture. Consequently, CO₂ flooding can raise production by about 15%. The oil containing dissolved CO₂ is then brought to the surface and the gases are flashed and CO₂ is separated for reinjection. In the second class of utilization, CO₂ is converted to chemicals and fuels⁹ via carboxylation or electrochemical avenues.

Indeed, using CO₂ as a viable feedstock for the chemical industry has been the prediction of visionary scientists and their studies of converting CO₂ to C₁ building block chemicals, which significant achievements have already been established. Already, around 130 Mt (mega or millions of tons) CO₂ are used annually to manufacture urea, salicylic acid, cyclic carbonates, and polycarbonates. Among them, the urea process consumes most of the CO₂ industrially. Urea is produced at around 185-190°C and a pressure range of 180-200 atm by reacting CO₂ with ammonia. Several other, potentially more sustainable alternatives to fossil fuels are being researched. For example, significant progress is being made in the production of biofuels from algae and woody biomass.^{11,12} Wind, solar, and tidal powers are also very popular and expanding rapidly along with increased interest in expanding nuclear energy.^{13,14}

1.3 Electrochemical CO₂ reduction

One of the solutions is an electrochemical conversion of CO₂ into useful chemicals (see Figure 1.2). The result is that greenhouse gas is transformed into products that can replace plastics and chemicals that are now produced from fossil feedstock. Heterogeneous electroreduction of carbon dioxide is an important energy conversion reaction of chemical products using renewable energy as an input.¹⁵ Important is that the electricity used for conversion must be renewable, or at least from a carbon-neutral source such as nuclear; otherwise, more CO₂ would be emitted in producing the electricity than would be used in the process. A result of CO₂ reduction is a variety of possible products such as carbon monoxide (CO),¹⁶⁻²⁰ formate (HCOO⁻)^{19,21-25}, formaldehyde (CO₂H), methane (CH₄),^{19,26,27} methanol (CH₃OH),²⁸⁻³⁰ and C₂ hydrocarbons^{19,31-35} and oxygenates, where these products can be used as commodity chemicals as well as fuels, thus allowing CO₂ to be recycled into compounds that can act as energy carriers (see Table 1.1).¹⁵ The cathodic process of the CO₂RR can be coupled to the oxygen evolution reaction at the anode side or to the anodic chlorine production, which is technologically used for Cl₂ production from NaCl via aqueous electrolysis.³⁶

Table 1.1: CO₂ reduction products with equilibrium potential E⁰ in V vs. RHE

Reaction	E ⁰
$2\text{H}^+ + 2\text{e}^- \rightarrow \text{H}_2(\text{g})$	0
$\text{CO}_2 + 2\text{H}^+ + 2\text{e}^- \rightarrow \text{HCOOH}(\text{aq})$	-0.12
$\text{CO}_2 + 2\text{H}^+ + 2\text{e}^- \rightarrow \text{CO}(\text{g}) + \text{H}_2\text{O}$	-0.10
$\text{CO}_2 + 6\text{H}^+ + 6\text{e}^- \rightarrow \text{CH}_3\text{OH}(\text{aq}) + \text{H}_2\text{O}$	0.03
$\text{CO}_2 + 8\text{H}^+ + 8\text{e}^- \rightarrow \text{CH}_4(\text{g}) + 2\text{H}_2\text{O}$	0.17
$2\text{CO}_2 + 12\text{H}^+ + 12\text{e}^- \rightarrow \text{C}_2\text{H}_5\text{OH}(\text{aq}) + 3\text{H}_2\text{O}$	0.09
$2\text{CO}_2 + 12\text{H}^+ + 12\text{e}^- \rightarrow \text{C}_2\text{H}_4(\text{q}) + 4\text{H}_2\text{O}$	0.08
$2\text{CO}_2 + 14\text{H}^+ + 14\text{e}^- \rightarrow \text{C}_2\text{H}_6(\text{q}) + 4\text{H}_2\text{O}$	0.14
$3\text{CO}_2 + 18\text{H}^+ + 18\text{e}^- \rightarrow \text{C}_3\text{H}_7\text{OH}(\text{aq}) + 5\text{H}_2\text{O}$	0.10

The production of syngas (synthesis gas) will play an important role in the 21st century. Syngas is a fuel gas mixture consisting of hydrogen and carbon monoxide, which can be produced from many sources, including natural gas, coal, biomass, or virtually any hydrocarbon feedstock, by reaction with steam (steam reforming), carbon dioxide (dry reforming), or oxygen (partial oxidation). Syngas is a crucial intermediate resource for the production of hydrogen, ammonia, methanol, and synthetic hydrocarbon fuels. Natural gas to liquid (GTL) and coal to liquid (CTL) technology yield syngas, which is then converted into liquid fuels via the Fischer-Tropsch process. CTL and GTL would allow the use of relatively abundant coal and natural gas to decrease our dependence on imported oil for transportation fuels, however the solution does not address the non-renewable nature of fossil fuels.¹⁵

For CO₂ electrochemical conversion to become feasible, two norms need to satisfy, (i) high energy efficiency and (ii) high reaction rates. The energetic efficiency defines the energy cost of producing the product, where it can be achieved through a combination of high selectivity towards a certain product (Faradaic or current efficiency) and low overpotentials. The equilibrium potentials for most of the carbon dioxide reduction half-reactions are close to 0 V vs. RHE, making the HER an additional competing reaction beyond those toward unwanted carbon-based by-products. The main challenge lays in increasing the energy efficiency, where the limiting step is the formation of a $\bullet\text{CO}_2^-$ radical anion intermediate. This step has a standard potential of -1.9 V versus SHE and is the reason for the high overpotentials.^{37,38}

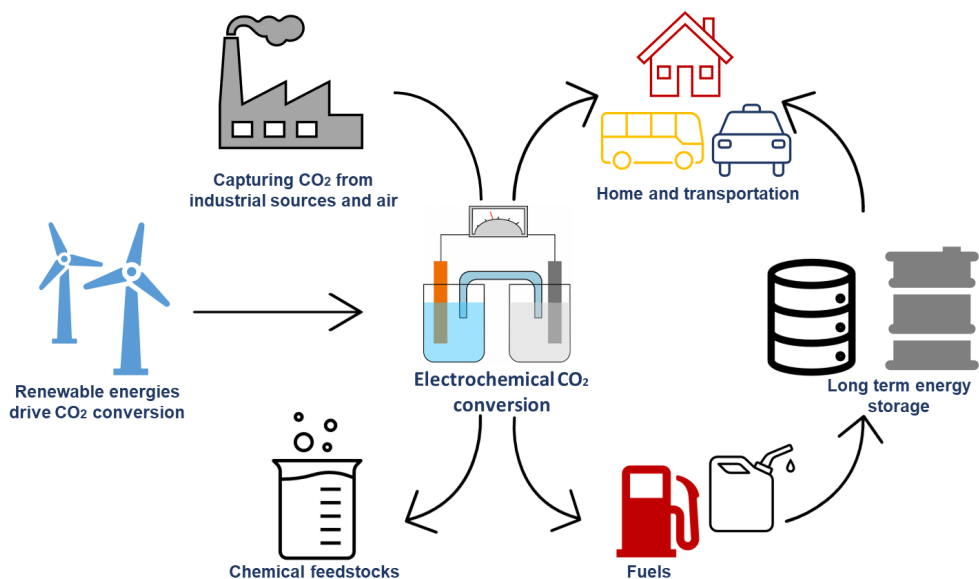


Figure 1.2: Overall picture of capturing CO₂ to electrochemically converting into commodity chemical feedstocks or fuels. Adapted from Ref.³⁹

The major challenge of CO₂ utilization is its high thermodynamic stability and so reactions that involve CO₂ conversion require high activation energy inputs. Therefore a catalyst is required, so it lowers the kinetic barrier and allow the reaction to occur under reasonable reaction conditions. For this reason focus of this issue is on CO₂ catalysis to valorize this waste material into potential new products, and to improve the efficiency of already existing technologies. Research contributions to this issue, have been supplied by some of the leading researchers in this area, considering either heterogeneous or homogeneous catalysis approaches.

1.4 Heterogeneous catalysis

In 1835 Baron J. J. Berzelius described catalysis as the property of substances that facilitate chemical reactions without being consumed in them.^{40,41} Catalytic reaction controls the changes of the route to equilibrium, where the focus is on reaction kinetics, and not thermodynamics. Catalysts do not change the thermodynamic equilibrium, the reaction thermodynamics describe your energy at your starting point and your destination.⁴² Catalysts can be either heterogeneous or homogeneous, depending on whether a catalyst exists in the same phase as the substrate. In heterogeneous catalysis, the catalyst is in a solid phase and then acts on substrates in a gaseous or liquid reaction mixture. The chemical reactions take place at the surface of the material, where at the interphase between the solid surface and the gas or liquid phase, atoms and molecules are exchanged. The catalyst has to be able to: adsorb reactants from the gas phase, make a bond and create a product molecule. And the newly formed molecule has to be able to desorb from the surface. Interaction between

a molecule and a surface can be classified as physisorption and chemisorption. Physisorption is a weak intermolecular (van der Waals) interaction, which do not involve a significant change in the electronic orbital patterns of the species involved. Chemisorption means that atoms form a chemical bond with the surface upon adsorption.

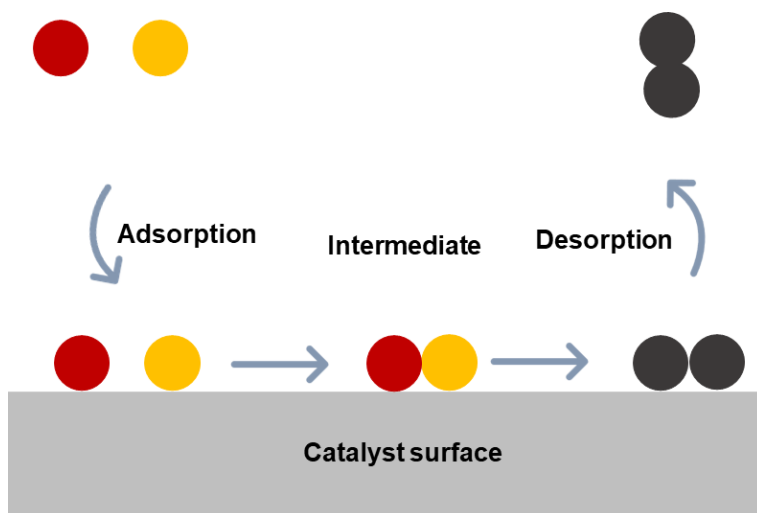


Figure 1.3: Schematic presentation of reaction steps in heterogeneous catalysis.

Catalytic reaction only happens on the active sites of the catalyst. A schematic representation is shown in Figure 1.3. Consequently, heterogeneous catalysts should have large surface area. So then, the catalytic materials can be made with a high surface area or the same support, can be used to stabilize active material (nanoparticles). Pure metals and metal oxides are the representatives of typical heterogeneous catalysts. The advantages of heterogeneous catalysis include, amongst others, easy catalyst separation in gas-liquid/solid systems (the catalyst is easily cleaned, or filtered), convenient handling in a fixed-, fluidised- or moving bed reactors, regeneration of catalyst and above all established reactor design principles. These are the reasons why heterogeneous catalysts are important for industry.⁴³

1.5 Motivation and objectives

The main objective of this thesis is to theoretically study CO₂ electrochemical conversion towards CO and higher products over single-site catalysts, in particular well defined and undefined carbon N-doped materials. In particular, I focused on understanding the activity and selectivity of these materials and the possibility of supported Ni nanoparticle reconstruction and disintegration. The objectives for each chapter are detailed below.

Chapter 3 Single atom catalyst for eCO₂RR

Supported metal nanostructures are the most widely used type of heterogeneous catalyst in industrial processes. The size of metal particles is a key factor in determining the performance of such catalysts. The ultimate small-size limit for metal particles is the single-atom catalyst (SAC), which contains isolated metal atoms singly dispersed on supports. SACs maximize the efficiency of metal atom use, which is particularly important for noble metal catalysts. Moreover, with well-defined and uniform single atom dispersion, SACs offer great potential for achieving high activity and selectivity.

- To describe the importance of single-atom catalysts as an alternative to metal nanoparticles; due to structural motifs, they may increase the selectivity of the CO₂RR over the competitive HER.
- To explain the application of coordination chemistry to homogeneous-like defects.

Chapter 4 Volcano trend in electrocatalytic CO₂ reduction activity over atomically dispersed metal sites on Nitrogen-doped carbon

The development of catalysts for electrochemical reduction of carbon dioxide with high activity and selectivity remains a grand challenge to render the technology usable. As promising candidates, metal-nitrogen-carbon (MNC) catalysts with metal atoms present as atomically-dispersed metal-N_x moieties (M = Mn, Fe, Co, Ni and Cu) were investigated as model catalysts. The distinct activity for CO formation observed along the series of catalysts is attributed to the nature of the transition metal in MN_x moieties, due to otherwise similar composition, structure, and morphology of the carbon matrix.

- To obtain the oxidation states, on which eCO₂RR occurs, using computational models of MN₄ moieties with computed Pourbaix diagram.
- To explain experimental activity with CO binding energy gained from DFT simulations as a descriptor.
- To obtain RDS of calculated eCO₂RR to CO Gibbs free reaction energies for each active site.
- To explain selectivity with the relationship between the adsorption(*CO₂⁻, *H)/desorption (CO) as a function of the d-antibonding occupation at electrochemical conditions.

Chapter 5 Stability and redispersion of Ni nanoparticles supported on N-doped carbons for the electrochemical CO₂ reduction

Electrocatalysts based on Ni supported on N-doped carbon materials have been computationally analyzed for the reduction of CO₂ considering both nanoparticles and single atom supports. In particular, there is a need for understanding the processes that happen under electrochemical conditions.

- To obtain the optimal activity and selectivity over different saturated and unsaturated N-doped carbon materials.
- To obtain surface energies of Ni NPs for nanoparticle shape.
- To investigate if the disintegration of Ni nanoparticles with Ni(CO)₂ can happen on studied materials, using Ostwald formalism.
- To develop a mechanism of Ni(211) surface reconstruction that is possible by the formation of Ni(CO)₄ species and which could redisperse into active single atoms.
- To study the influence of high CO coverages on the limitation of the HER.

Chapter 6 Temperature dependent selectivity of CO₂ electrochemical reduction on CoTPP/MWCNT catalyst

Electrochemical reduction of CO₂ to hydrocarbons on molecular catalysts is possible due to their structural tunability, selectivity, and low amount of loading on the support. Here we focus on the potential and temperature influence of the CO₂ reduction reaction product selectivity over cobalt porphyrin (CoTPP)/multiwalled carbon nanotube (MWCNT) hybrid catalyst.

- To describe the interaction between the CoTPP complex and graphene layer with DFT.
- To obtain Gibbs free reaction energies of calculated eCO₂RR to CH₃OH and CH₄ for CoTPP/MWCNT.
- To understand activity and selectivity with DFT simulations, where the difference between the binding energies of CH₂O oxygen bonding, the strength of *CO, and protonation of *CHO intermediate are analyzed.
- To explain the difference in selectivity with increasing negative potential.
- To explain the temperature influence on the product distribution.

Chapter 2

Theoretical methods

2.1 Schrödinger equation

The modern approach to describe atomistic properties of systems is based on quantum mechanics, namely molecules and atoms can be described with wave function, where we get the Schrödinger equation solution. It describes how the quantum state of a physical system changes in time and it has form of eigenvalue equation. For a set of atomic nuclei and electrons, according to Born Oppenheimer approximation the motions of nucleus and electrons are treated separately.⁴⁴ Thereby a fundamental time independent electronic equation follows as:

$$\hat{H}\Psi = E\Psi \quad (2.1)$$

Where \hat{H} is the Hamiltonian operator, Ψ -many-body wave function depending on the spatial coordinates of the electrons and are a set of solutions (Ψ_n) or eigenstates of the Hamiltonian, E presents electronic energy of the system (to each set of solutions correspond an eigenvalue (E_n)).

Because of the electron - nuclei, and electron - electron interaction the complete Schrödinger equation is written as:

$$\left[-\frac{\hbar^2}{2m_e} \sum_{i=1}^N \nabla_i^2 + \sum_{i=1}^N V(\mathbf{r}_i) + \sum_{i=1}^N U(\mathbf{r}_i, \mathbf{r}_j) \right] \Psi = E\Psi \quad (2.2)$$

m_e - electron mass

1st term - kinetic energy of each electron

2nd term - the interaction energy between electrons and the atomic nuclei

3rd term - the interaction energy between electrons

$\Psi = \Psi_1(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_n)$ is the electronic wave function which is a function of each coordinates of all N electrons and can be approximated as a product of individual wave functions. If we want to solve the Schrödinger equation for CO_2 molecule, which has $N=22$ electrons in total, the full wave function is going to be 66 dimensional function. When we adsorb the CO_2 molecule to a solid surface with 66 Carbon atoms, the full wave function becomes 1254

dimensional many-body function. Therefore Schrödinger equation still remains a many-body issue which is too complex to solve for a simple catalytic systems. The term that describes the electron-electron interaction is one of the most critical and it is worth realizing that the wave function for any particular set of coordinates cannot be directly observed and is also difficult to give the wave function any physical significance. A basic idea to simplify the N-electron problem consists in finding some physical quantity that defines the system uniquely without growing in complexity as a function of the number of electrons N. Here is where the Density functional theory is introduced, which focuses on measuring the probability of electron density at a certain position in space. The equation is then written as individual electron wave function:

$$n(\mathbf{r}) = 2 \sum_i \psi_i^*(\mathbf{r})\psi_i(\mathbf{r}) \quad (2.3)$$

2.2 Density Functional Theory

2.2.1 The Hohenberg-Kohn approach

In 1964 Density functional theory was developed.⁴⁵ The method of applying electron density to solve Schrödinger equation was developed by Hohenberg and Kohn (HK). Thereby the number of variables is reduced from $3N$ to 3 .

Theorem 1 The first theorem proved by HK is: The ground state energy from Schrödingers equation is a unique functional of the electron density. Re-stated the ground state energy can be expressed as a functional of the electron density $n(\mathbf{r})$:

$$E_0 = F[n(\mathbf{r})_0] = E[n(\mathbf{r})_0] \quad (2.4)$$

Therefore the ground state electron density uniquely determines all properties including the energy and wave function. The theorem says that such a functional F of the electron density exists, but does not tell which mathematical expression it has.

Theorem 2 The second HK theorem defines a property of the functional: The electron density that minimizes the energy of the overall functional is the true electron density corresponding to the full solution of the Schrödinger equation:

$$E_0 = \min E[n(\mathbf{r})] \quad (2.5)$$

2.2.2 The Kohn-Sham approach

A convenient way to rewrite a functional described by Hohenberg-Kohn theorem is in terms of single electron wave functions $\Psi_i(\mathbf{r})$.

$$\hat{H}_{KS}\Psi_i = E_i\Psi_i \quad (2.6)$$

The energy functional can be now written as:

$$E_{\text{KS}}[\Psi_i] = -\frac{\hbar^2}{m} \sum_{i=1} \int \Psi_i^* \nabla_i^2 \Psi_i d^3 \mathbf{r} + \int V(\mathbf{r}) n(\mathbf{r}) d^3 \mathbf{r} + \frac{e^2}{2} \int \int \frac{n(\mathbf{r}) n(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d^3 \mathbf{r} d^3 \mathbf{r}' + E_{\text{XC}}[\Psi_i] \quad (2.7)$$

Where 1st term-is the electron kinetic energy

2nd term-the Coulomb interaction between electrons and nuclei (external potential)

3rd term-the Coulomb interactions between pairs of electrons

$E_{\text{XC}}[\Psi_i]$ is the exchange correlation functional that include all the many body effects that are not included in the 'known' terms. But however the exact form of the 3rd term in equation 2.8 is not known. This problem could be solved by applying the next approach. Kohn and Sham turn the many body problem to single electron problem.⁴⁶ Hence the task of finding the right electron density can be expressed in a way that involves solving a set of equations and each equation only involves a relation for a single electron:

$$\left[-\frac{\hbar^2}{2m} \nabla_i^2 + V(\mathbf{r}) + V_{\text{H}}(\mathbf{r}) + V_{\text{XC}}(\mathbf{r}) \right] \Psi_i(\mathbf{r}) = \varepsilon_i \Psi_i(\mathbf{r}) \quad (2.8)$$

This is similar to equation 2.2 but without summation terms since are single electron wave functions and depend on only three spatial variables.

Where, $V(\mathbf{r})$ -the potential which defines the interaction between an electron and the atomic nuclei.

V_{H} -the Coulomb repulsion between the considered electron and the total electron density defined by all electrons in the problem (Hartree potential):

$$V_{\text{H}}(\mathbf{r}) = e^2 \int \frac{n(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d^3 \mathbf{r}' \quad (2.9)$$

Therefore corrections of the many body effects of electrons have to be included in the final potential V_{XC} which defines exchange and correlation contribution to the single electron equations:

$$V_{\text{XC}}(\mathbf{r}) = \frac{\delta E_{\text{XC}}(\mathbf{r})}{\delta n(\mathbf{r})} \quad (2.10)$$

2.2.3 Exchange correlation energy

In the Kohn Sham scheme the expression for the exchange correlation functional E_{XC} is unknown, for which several approximations has been developed.

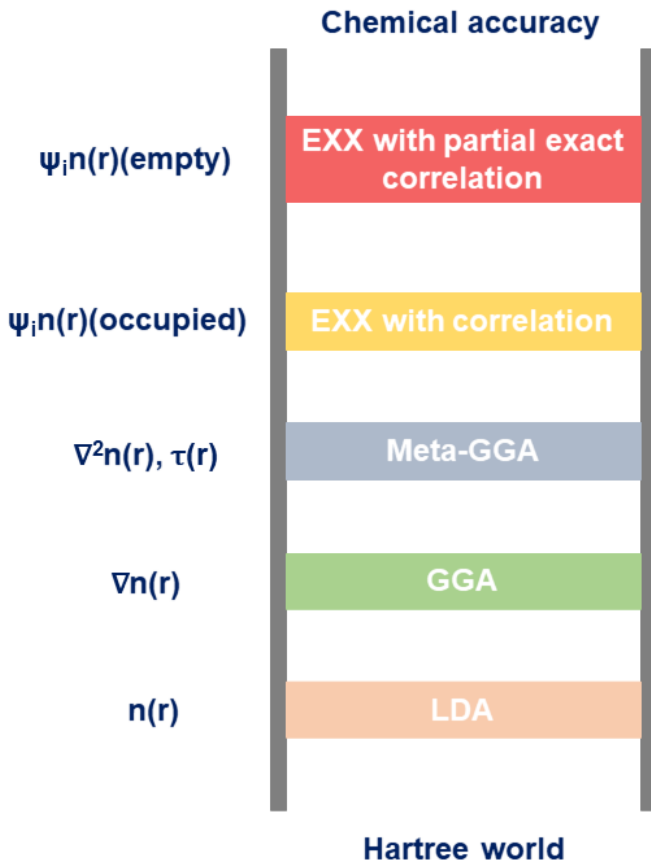


Figure 2.1: Perdew’s classification ladder of DFT functionals. Adapted from Ref.[⁴⁵]

In Figure 2.1 the hierarchy of accuracy in the exchange correlation functionals has been defined by John Perdew as ”the Jacobs ladder of DFT”.^{45,47}

Local density approximation (LDA), exchange correlation energy is approximated to the one of an homogeneous electron gas, where the electron density is constant at all points in space, $n(\mathbf{r})$. E_{XC} energy can be written as:

$$E_{XC, LDA} = \int \varepsilon_{XC}^{\text{hom}}(n(\mathbf{r}))n(\mathbf{r})d\mathbf{r} \quad (2.11)$$

Where $\varepsilon_{XC}^{\text{hom}}(n(\mathbf{r}))$ is the energy of the exchange correlation hole in the homogeneous electron gas of density n . The LDA describes properties such as lattice constants, vibrational frequencies and equilibrium geometries for bulk materials. Other properties such as dissociation energies of molecules, cohesive energies are not very well described using this functional. The spin-polarized version (LSDA) can be applied for magnetic materials.

$$E_{XC, LSDA} = \int \varepsilon_{XC}^{\text{hom}}(n(\mathbf{r}))n_{\uparrow}, n_{\downarrow}(\mathbf{r})d\mathbf{r} \quad (2.12)$$

Generalized gradient approximation (GGA), uses information about the local electron density and the local gradient in the electron density. GGA

has a good accuracy and reasonable computational cost, so this approximation is used to describe many systems in homogeneous and heterogeneous catalysis. Also, it corrects the overbinding tendency of LDA functionals. With GGA approach E_{XC} energy can be written as:

$$E_{XC, GGA} = \int \varepsilon_{XC}^{\text{hom}}(n(\mathbf{r}), \nabla n(\mathbf{r})) d\mathbf{r} \quad (2.13)$$

The gradient of the electron density can be included in a GGA functional in different ways, therefore a large number of functionals has been developed. The most widely used functionals for bulk solids and surfaces are Perdew-Wang (PW91),⁴⁸ Perdew-Burke-Ernzerhof (PBE)^{49,50} and revised RPBE.⁵¹ For molecules, hybrid functional B3LYP can be used. Thereby the best choice of functional will depend on the system of interest.

Meta-Generalized gradient approximation (Meta-GGA), calculations are computationally demanding, and the improvement of the results from GGA-type functionals is often negligible. E_{XC} energy can be written as:

$$E_{XC, GGA} = \int \varepsilon_{XC}^{\text{hom}}(n(\mathbf{r}), \nabla n(\mathbf{r}), \nabla^2 n(\mathbf{r})) d\mathbf{r} \quad (2.14)$$

2.2.4 Implementation of DFT

In order to find the right electron density an iterative method is applied:

1. We start with a trial electron density $n(\mathbf{r})$
2. The trial $n(\mathbf{r})$ is used to solve the Kohn-Sham equations to obtain the single electron wave functions $\Psi_i(r)$.
3. The single electron wave functions is used to calculate an electron density, $n_{KS}(\mathbf{r}) = 2 \sum_i \Psi_i^*(\mathbf{r}) \Psi_i(\mathbf{r})$
4. Finally the calculated electron density $n_{KS}(\mathbf{r})$ is compared with the trial electron density $n(\mathbf{r})$. If they are different, the trial electron density is updated and the loop is repeated again. If they are the same (with small tolerances), the ground state density is obtained.

2.2.5 Periodic systems

Plane wave basis sets

A basis set is a set of functions that is used to represent the electronic wave function in density-functional theory. The one-electron Kohn-Sham orbitals $\Psi_i(\mathbf{r})$ (Equation 2.6), can be expanded using different basis:

- spatially localized functions (Slater- or Gaussian-type orbitals)

- spatially extended functions (plane waves)

For the periodic crystalline solids, plane wave approach is used. In order to satisfy Bloch's Theorem,⁵² the electronic wave function can be expressed in form of a plane wave $e^{i\mathbf{k}\cdot\mathbf{r}}$ with wave vector \mathbf{k} multiplied by a periodic function $u_{n\mathbf{k}}(\mathbf{r})$ with the same periodicity as the crystalline lattice:^{52,53}

$$\Psi_{i, \mathbf{k}}(\mathbf{r}) = e^{i\mathbf{k}\cdot\mathbf{r}}u_{i\mathbf{k}}(\mathbf{r}) \quad (2.15)$$

The cell periodic function ($u_{i\mathbf{k}}(\mathbf{r})$) is periodic in space with the same periodicity of the supercell. The electronic wave function is written as a sum of plane waves (Equation 2.16):

$$u_{i,\mathbf{k}}(\mathbf{r}) = \sum_{\mathbf{G}} C_{i,(\mathbf{k}+\mathbf{G})} e^{i\mathbf{k}\cdot\mathbf{r}} \quad (2.16)$$

Where $C_{i\mathbf{k}+\mathbf{G}}$, are the expansion coefficients, \mathbf{k} is the wave vector in the first Brillouin zone and \mathbf{G} is the reciprocal in the reciprocal wave vectors. The total plane wave expansion can now be written as:

$$\Psi_{i, \mathbf{k}}(\mathbf{r}) = C_{i,\mathbf{k}+\mathbf{G}} e^{i\mathbf{k}+\mathbf{G}\cdot\mathbf{r}} \quad (2.17)$$

Reciprocal space

A crystal is defined as a three dimensional periodic structure and can be described as a Bravais lattice, which shows its periodicity as a display of repeated discrete points with lattice vector \mathbf{R} (Equation 2.18). The lattice vectors \vec{a}_i are linearly independent. The parameters n_i are integer number, because the translations are discrete.^{53,54}

$$\vec{R} = n_1\vec{a}_1 + n_2\vec{a}_2 + n_3\vec{a}_3 \quad (2.18)$$

With operations shown in Equation 2.19 we can define reciprocal lattice with cell vectors \vec{b}_j for a Bravais lattice with cell vectors \vec{a}_j . Reciprocal space (k-space) is periodic for translations over \vec{b}_1 , \vec{b}_2 , and \vec{b}_3 . The primitive cell of the k-space is known as Brillouin zone.

$$\begin{aligned} \vec{b}_1 &= 2\pi \frac{\vec{a}_2 \times \vec{a}_3}{\langle \vec{a}_1, \vec{a}_2, \vec{a}_3 \rangle} \\ \vec{b}_2 &= 2\pi \frac{\vec{a}_3 \times \vec{a}_1}{\langle \vec{a}_1, \vec{a}_2, \vec{a}_3 \rangle} \\ \vec{b}_3 &= 2\pi \frac{\vec{a}_1 \times \vec{a}_2}{\langle \vec{a}_1, \vec{a}_2, \vec{a}_3 \rangle} \end{aligned} \quad (2.19)$$

k-points

In order to choose appropriate k-points for calculations of periodic structures, Monkhorst-Pack method should be used.^{55,56} Larger lattice vectors in real

space correspond shorter lattice vectors in reciprocal space. Therefore increasing the volume of the cell reduces the number of k-points. Exploring the symmetries in the k-points can significantly reduce the computational time and cost.

Energy cutoff

Another way to lower computational cost is by limiting the number of plane-waves by setting a kinetic energy cutoff. The kinetic energy of the systems is:

$$E = -\frac{\hbar^2}{2m}|\mathbf{k} + \mathbf{G}|^2 \quad (2.20)$$

Next step is to reduce the infinite sum from Equation 2.20 and to include only the solutions that have the kinetic energy lower than a certain value:

$$E_{\text{cut}} = -\frac{\hbar^2}{2m}G_{\text{cut}}^2 \quad (2.21)$$

And the infinite sum from equation (2.21) reduces to:

$$\Psi_{\mathbf{k}}(\mathbf{r}) = \sum_{|\mathbf{G}+\mathbf{k}| < G_{\text{cut}}} C_{\mathbf{k}+\mathbf{G}} e^{i(\mathbf{k}+\mathbf{G})\mathbf{r}} \quad (2.22)$$

This parameter has to be defined whenever a DFT calculation is performed. Convergence tests should be done for each system. In general many packages provide default values. A key point to remember is whenever a DFT calculation is performed for multiple systems that are going to be compared, the same energy cutoff should be used for all systems.

Pseudopotentials

Electrons in atoms can be divided into valence and core electrons. Valence electrons responsible for most of the chemical bonding, while the core electrons are tightly bound to the nuclei. Pseudopotential approximation is introduced to replace the electron wave function from a set of core electrons with a smoothed wave function. This is established so as to match various important physical and mathematical properties of the true ion core. This approach is greatly necessary to reduce the computational cost when plane waves are used as basis set. The properties of the core electrons are fixed in all calculations, which is known as the frozen core approximation. In the DFT codes typically a library of pseudopotentials that includes an entry for almost each element in the periodic table is provided. There has to be defined a minimum energy cutoff that has to be used in calculations. Most used families of pseudopotentials are the Vanderbilt ultrasoft pseudopotentials (US-PP)⁵⁷ and the Projector Augmented Wave ones (PAW).^{58,59} In this thesis PAW approximations are used, as they describe transition metals better.⁶⁰

General computational details

All electronic structures were performed using spin-polarized Density Functional Theory (DFT) as implemented in the Vienna Ab initio Simulation Package (VASP).^{58,61} The GGA PBE-D3 was the functional of choice, where D3 stands for Grimme dispersion.^{62,63} Core electrons were described by projector augmented wave (PAW),^{58,59} and valence electrons were expanded by plane wave basis sets with a kinetic energy cut-off of choice. In each chapter of the thesis the details of calculations are given. They includes the sizes of the unit cell, k-point sampling, kinetic energies cutoffs and the cutoff for the density grid. For all the systems described before, we obtained the energies of the relevant intermediates in the electrochemical environment under eCO₂RR conditions by considering the computational hydrogen electrode (CHE). Gas-phase molecules that were used in the thesis as a reference were calculated with PBE functional. Zero-point, entropic, and heat capacity contributions are shown in Table 2.1.⁶⁴ The solvation correction to adsorbate energy was applied to *COOH (-0.25 eV) and *CO (-0.1 eV).⁶⁴ Corrections of DFT-based gas-phase errors for standard PBE Gibbs free energies of CO₂ ($\epsilon=-0.19$ eV) and CO ($\epsilon=0.24$ eV) were applied. The corrections were calculated as:

$$\Delta G_{\text{DFTcorr}}^{\circ} = \Delta G_{\text{DFTcorr}}^{\circ} - \left(\sum \varepsilon_{\text{T}}^{\text{P}} - \sum \varepsilon_{\text{T}}^{\text{R}} \right) \quad (2.23)$$

where the $\varepsilon_{\text{T}}^{\text{R}}$ is the sum of all errors for reactants and $\varepsilon_{\text{T}}^{\text{P}}$ for the products.⁶⁵ solvation and error correction were added for studied materials in Chapters 4 and 5. The optimized structures can be retrieved from ioChem-BD.⁶⁶ A unique ioChem-BD link of the DFT calculations for each of the projects will be provided in according chapters.

Table 2.1: Zero-point, entropic, and heat capacity contributions (in eV) for gas-phase molecules used as a reference.

Molecule	ZPE	TS	C _p
CO ₂	0.31	-0.66	0.10
CO	0.13	-0.61	0.09
H ₂ O	0.57	-0.67	0.10
H ₂	0.28	-0.40	0.09
NiCO ₄	0.81	-0.45	

2.3 Basic concepts of computational electrochemistry

2.3.1 The Computational Hydrogen Electrode

Electrochemical reactions, where electrons and protons are transferred, can be modelled by applying the computational hydrogen electrode (CHE),^{67–69} where simplification of combined chemical potential of protons and electrons is applied. Schematic representation of CHE is presented in Figure 2.2. CHE approach is based on the following hypotheses:

1. The Standard Hydrogen Electrode (SHE) is taken as reference electrode and is define as 0 V and pH=0. The chemical potential for the reaction $\text{H}^+ + \text{e}^-$ can be related to the Gibbs free energy of $\frac{1}{2}\text{H}_2$. Therefore, at standard conditions (298 K, 1 bar of gas phase H_2)

$$\Delta G_{\text{H}^+} = \frac{1}{2}G_{\text{H}_2, \text{DFT}} \quad (2.24)$$

2. Proton Gibbs free energy is corrected by pH:

$$G_{\text{pH}} = -k_{\text{B}}T\ln(\text{H}^+) = k_{\text{B}}T\ln 10 \cdot \text{pH} \quad (2.25)$$

3. The effect of the electric potential is included adding to a thermodynamic state the energy of each transferred electron, $-neU$, where U is the applied potential.
4. The electrostatic field caused by the electrical double layer is neglected.

Hydrogen can easily be calculated and then referenced back to the chemical potentials of the electrons and protons in solutions. Once the hydrogen molecule energy is calculated and then add the potential U_{RHE} , free energy G_{CHE} can be written as:

$$G_{\text{CHE}}(n, U_{\text{RHE}}) = \Delta G = \Delta E - T\Delta S + \int_0^T C_{\text{p}}dT + \Delta \text{ZPE} - n\left(\frac{1}{2}\mu_{\text{H}_2} - eU_{\text{RHE}}\right), \quad (2.26)$$

where n is the number of proton-electron transferred and G_{CHE} is a function of both potential and proton-electron pair transfers. The thermodynamic contributions to G_{CHE} from the support are typically neglected, and only the contributions from the gas phase and vibrational entropy of the adsorbate are included. The simplicity of the computational hydrogen electrode is obtained by the choice of references and the assumption of electrolyte and electrode in equilibrium. One could include solvation contribution either through explicit water model⁶⁸ or implicit model based on solvent dielectric permittivity.⁷⁰ As will be shown in the next sections, the CHE is applicable for calculating many

interesting properties in electrochemical reactions, such as surface Pourbaix diagrams and free energy diagrams. The application of CHE model to decouple proton-electron transfers, such as CO₂ adsorption, has intrinsic limitations.⁷¹ An alternative model derived from Marcus theory (Ref.⁷²) has been proposed to account for the transition between sequential and concerted proton-electron transfer.⁷³⁻⁷⁵

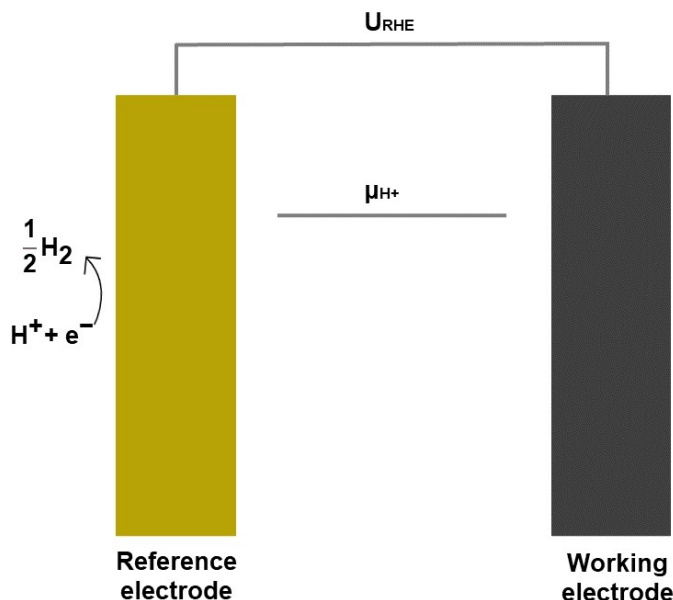


Figure 2.2: Schematic representation of the computational hydrogen electrode (CHE). The hydrogen reaction is in equilibrium with the reference electrode. The protons for both the working and the reference electrode have the same chemical potential, while the electrons at the working electrode are affected by the applied potential (U_{RHE}). In this case U_{RHE} is negative. Adapted from Ref.⁶⁹

Pourbaix diagram

Classical Pourbaix diagram, also known as a potential/pH diagram, is a plot of possible thermodynamically stable states of a compound as a function of pH and potential. As noted above, these diagrams are essentially phase diagrams that map these conditions, most typically in aqueous solutions where different redox species are stable (see Figure 2.3).

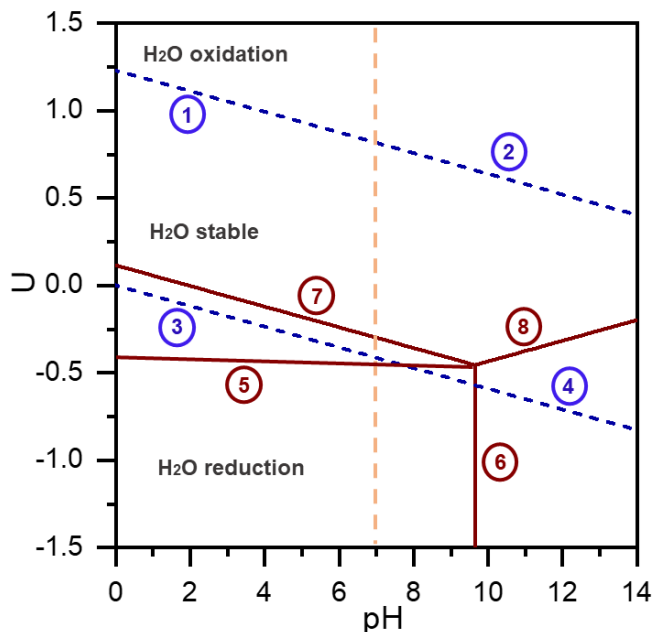
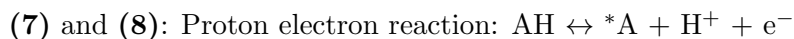
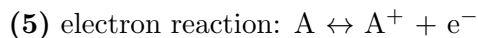
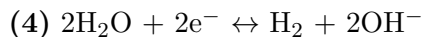
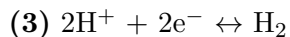
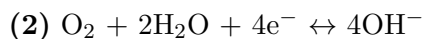
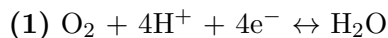


Figure 2.3: Pourbaix diagram. The blue lines correspond to the water oxidation and reduction reactions, the red lines separate the different phases, and the orange line corresponds to pH 7.

Areas in the Pourbaix diagram mark regions where a single species is stable, while lines mark places where two species exist in equilibrium. Reaction of numbered lines in Figure 2.3 are presented below;



The proton reaction (vertical line) is reaction that only depend on pH. Pure redox reaction (horizontal line) depends on applied potential. The proton-electron reaction depends on pH and potential and has a slope of -0.0592 V/pH (RHE scale). To construct the surface Pourbaix diagram, we take into consideration DFT formation energies of the clean surfaces with according

experimental oxidation states. The aim was to identify the possible oxidation states on which CO₂ reduction reaction could happen at experimental condition. Furthermore, the surface Pourbaix diagrams assume that no other reaction occur, for example HER, where on metals such as Au, Ag, or Cu hydrogen gets produced. However on single-site materials, because of their unique structure, HER is limited.

Computed surface Pourbaix diagrams were constructed for FeN₄⁻, CoN₄⁻, and Ni₄⁻ species (see Chapter 4). Figure 2.4 shows an example of computed diagram for FeNC catalyst, for which two different oxidation states were observed (2+ and 3+). Therefore, Fe(II)N₄-H₂O, Fe(II)N₄-O, and Fe(III)N₄-OH species were chosen.

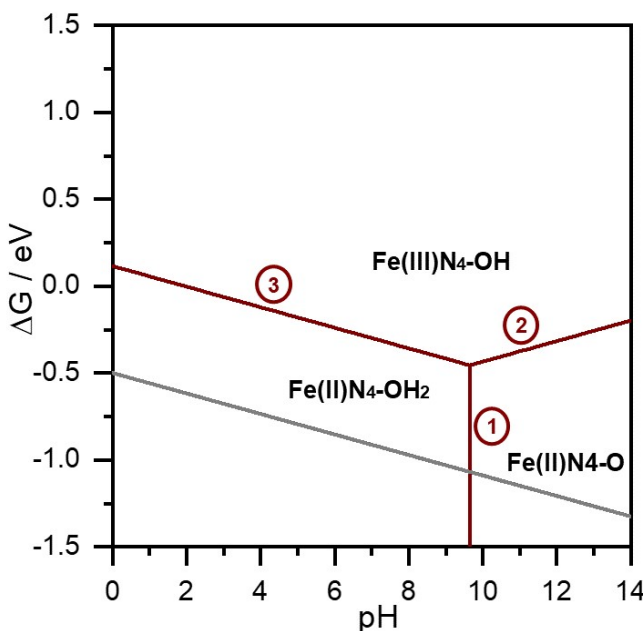


Figure 2.4: Computed Pourbaix diagram The red lines separate the different states of FeN₄⁻, and the orange line corresponds to pH 7. Gray line corresponds to -0.5 V vs. RHE.

The oxidation states were determined based on calculated magnetization of the metal using crystal field theory (see Chapter 3). Once the formation energies of the the chosen species were obtained the Pourbaix diagram was build. Numbered reactions in Figure 2.4 are presented below:

- (1) $\text{Fe(II)N}_4\text{-H}_2\text{O} \leftrightarrow \text{Fe(II)N}_4\text{-O} + 2\text{H}^+$
- (2) $\text{Fe(III)N}_4\text{-OH} + \text{e}^- \leftrightarrow \text{Fe(II)N}_4\text{-O} + \text{H}^+$
- (3) $\text{Fe(III)N}_4\text{-H}_2\text{O} \leftrightarrow \text{Fe(II)N}_4\text{-O} + \text{H}^+ + \text{e}^-$

This approach is very simple and in particular it can be applied to screening projects and to study trends. With these trends we show the correlation between experimental and theoretical results.

2.3.2 Explanation of electrocatalytic reactions and trends

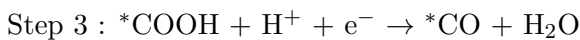
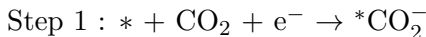
Gibbs free energy diagrams

A way to compare the experimental result with theoretical prediction is the Gibbs free energy diagram. It allows us to interpret different reaction pathways by calculating the reaction intermediates. Once the intermediates for certain reaction are calculated we can get theoretical onset potential, which is numerically equivalent to the additive inverse of the largest positive reaction energy considering electrochemical steps. Onset potential can be calculated by

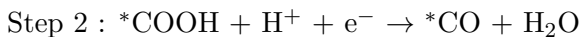
$$U_{\text{onset}} = -\max(\Delta G_1, \Delta G_2)/e^- \quad (2.27)$$

where the ΔG is Gibbs free energy for certain intermediate.^{67,76} For constructing free diagram for CO₂ reduction towards CO, we can follow two different mechanisms;

- First mechanism, CO₂ may be activated by an electron transfer prior to its adsorption, and the adsorbed species is stabilized by a proton close to the surface (one decoupled and coupled proton-electron transfers):⁷⁷⁻⁸¹



- Second mechanism (two coupled proton-electron transfers).⁸²



In order to construct Gibbs free energy diagram (see Figure 2.5), the energy of each reaction step was calculated with respect to the energies of reactants, CO₂ and support. ΔG for each step was calculated as shown in Equation 2.26:

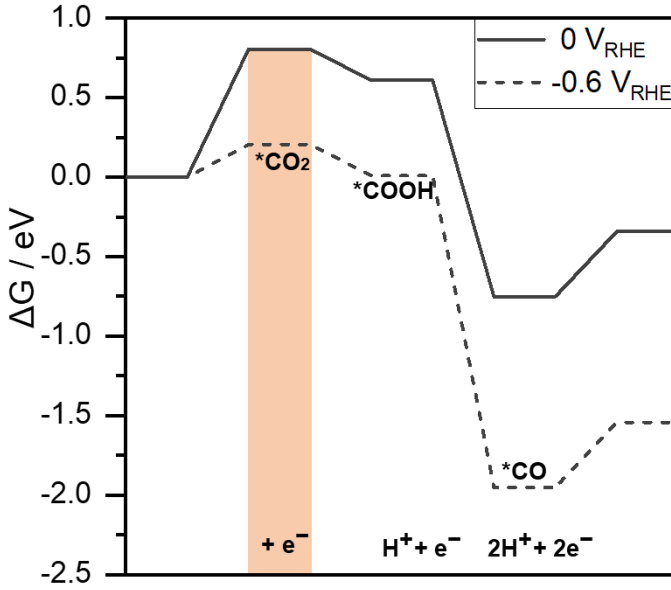


Figure 2.5: The Gibbs free energy diagram for CO₂ reduction towards CO, at 0 V (solid gray line) and -0.6 V (dashed gray line) vs. RHE. In addition, the light orange area indicates *CO₂ intermediate, which is not part of CHE formalism.

$$G_1 = \Delta G_{*CO_2} - \Delta G_{CO_2} - \Delta G_* + (eU_{RHE}) \quad (2.28)$$

$$G_2 = \Delta G_{*COOH} - \Delta G_{CO_2} - \Delta G_* + \left(\frac{1}{2}H_2 + eU_{RHE}\right) \quad (2.29)$$

$$G_3 = \Delta G_{*CO} - \Delta G_{CO_2} - \Delta G_* + (H_2 + 2eU_{RHE}) \quad (2.30)$$

$$G_{des} = \Delta G_{CO} + \Delta G_* - \Delta G_{*CO} \quad (2.31)$$

Since the modelling of decoupled proton-electron transfers is challenging from a DFT point of view, in this thesis second mechanism was used only when theoretical onset potential was calculated. An example of how it was calculated is shown below:

$$G_{1ons} = \Delta G_3 - \Delta G_2 \quad (2.32)$$

$$G_{2ons} = \Delta G_{des} - \Delta G_3 \quad (2.33)$$

Furthermore, the free energies calculated for the free energy diagram can be used to find scaling relations and construct Sabatier activity volcanoes as a function of reactivity descriptors.

Scaling relation

Once the adsorption energies of intermediates are calculated, the question is whether different adsorption energies are related to one another and how. One of the solution for understanding relation would be scaling relationship, which corresponds to sets of binding energies through linear functions (example in Figure 2.6), and have been found between C-, N-, S-, and O- derived adsorbates on catalyst surfaces.^{83–85} Furthermore scaling relationships are not limited to the binding energies of adsorbed atoms and their hydrogenated species but they hold for a broader group of adsorbates that fulfil certain conditions related to the electronic structure of the surface and the adsorbates.⁸⁶

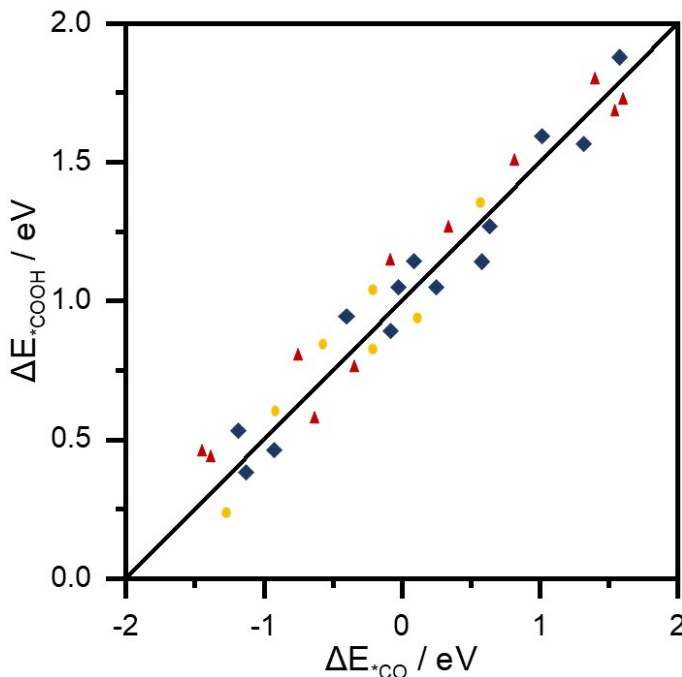


Figure 2.6: Linear scaling relation between *COOH and *CO binding energy.

Toward the electrocatalytic CO₂RR, it is difficult to attain a low overpotential because of the adsorbate (particularly, *COOH, *CO and *CHO or *COH) scaling relations. Cu electrodes are good example because the protonation of *CO is the rate determining step toward CH₄ with the highest free energy change. Protonation from *CO to *CHO or *COH is determined by the electrode surface structure, electrode potential, acidic or alkaline environment, standard strategy to reduce the free energy change of this step is to strengthen the binding energy of *CHO or *COH and weaken the binding strength of *CO. The optimal conversion of CO₂ with the lowest limiting potential can be achieved by tuning the free energy change of each elementary reaction to minimum. The way to break the above scaling relations, would be with introducing p states, reducing metal coordination numbers, tethering active ligands, alloying metal catalysts, and designing dual functional sites.⁸⁷

On the other hand, we can use the scaling relations, without breaking them by using single-site catalysts. They have a structural advantages over typical metal and metal oxide catalyst, because intermediate can only bind on the top site of the catalyst. Mostly it can be useful for eCO₂RR, as the competitive HER can be limited by lacking the hollow site.⁸⁸

Volcano plot

French chemist Paul Sabatier qualitatively described the adsorption properties of the optimal catalyst, where the best catalysts should bind adsorbate with an intermediate strength: not too weakly in order to be able to activate the reactants, and not too strongly to be able to desorb the products. Today this statement is known as the Sabatier principle of catalysis,⁸⁹ however the quantitative description was portrayed in Balandin-Sabatier volcano-shaped curves, which consider the implication of the rate (catalyst activity) relative to a descriptor of substrate binding.^{90,91} Balandin developed a theory mainly concerned with the steric relation between reactant and the surface structure of the catalyst. Figure 2.7 presents illustration of volcano plot, where the top of the volcano corresponds to the highest activity and the intermediate bond strength.

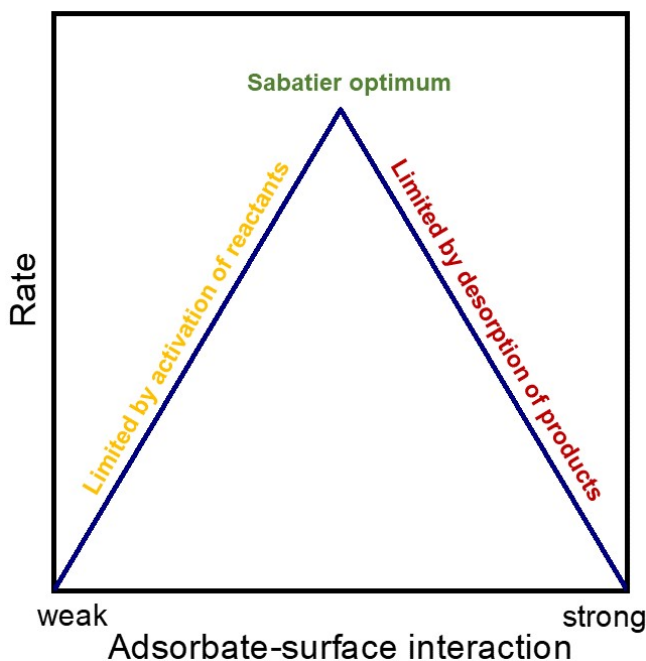


Figure 2.7: The catalytic volcano plot. Illustration of the qualitative Sabatier principle.

For electrochemical reactions such as oxygen evolution reaction (OER)/oxygen reduction reaction (ORR) and HER, Sabatier principle coupled with scaling relations can be used to describe activity with overpotential or current density as function of a potential-determining binding energy of intermediates.^{43, 69, 92–99} Volcano trends have been generated for numerous other catalytic reactions such as NO_x decomposition, ammonia synthesis, and CO oxidation.⁹⁴

Chapter 3

Carbon based single atom catalyst

The green conversion of CO₂ to fine chemicals calls for a new generation of efficient and robust catalysts. Herein I present single-atom catalysts (SACs), in which all metal species are atomically dispersed on a solid support, and which often consist of well-defined mononuclear active sites and enables 100% atom utilization. They are a bridge between homogeneous and heterogeneous catalysts and often show outstanding performance in many industrially important reactions, including thermochemical, electrochemical and photochemical conversions. The relationship between catalytic performance and the active site structure is discussed in terms of the valence state, coordination environment and anchoring chemistry of single atoms to the support, in an effort to guide the rational design of SACs, which has traditionally been dominated by homogeneous catalysis.

3.1 Decreasing size particle toward single atom

In heterogeneous catalysis, enormous efforts have been devoted toward improving the performance of supported metal catalysts by downsizing the metal particles. The supported metal particles have multiple active sites with different performances. Such heterogeneity affects the efficient utilization of metal active sites and thus reduces selectivity toward a specific product. The size of metal particles therefore becomes one of the most critical factors determining the reactivity and specificity of supported metal catalysts. The bulk materials can be turned into nanoparticles, subnanoclusters, and finally single metal atoms (see Figure 3.1). The size reduction generates an increase of unsaturated coordination environment of the metal species. Accordingly, the surface free energy of the metal components increases, and the metal sites become more and more active for chemical interactions with the support and adsorbates, which accounts for the size effects of metal nanocatalysts.

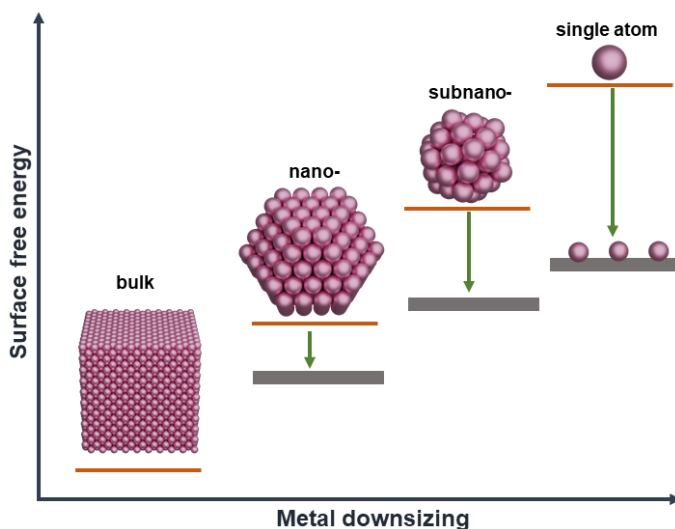


Figure 3.1: Schematic diagram illustrates minimizing the metal size from bulk to single atom.

Despite their small sizes, subnanoclusters still contain multiple active centers and are not always the most desirable active sites for specific catalytic processes.¹⁰⁰ Extreme downsizing of particles would be SACs, which incorporate many advantageous features of homogeneous and heterogeneous catalysts. Typical homogeneous catalysts have well defined structures but are not easily recycled because they are difficult to separate from reaction media. Heterogeneous catalysts can have high stability and are easily separated from reactants and products. However, their active sites may be non-uniform in nature, and only now there are technologies emerging that enable their in operando study.¹⁰¹ Thanks to SACs unique structural properties, not only they have a maximal atom utilization efficiency but also offer opportunities for tuning reaction rates and selectivity, thus presenting them as bridge between ho-

mogeneous and heterogeneous catalysts. Therefore, SACs may provide a good platform for understanding structure-activity relationships on an atomic scale. Searching for catalysts with well defined single active centers is necessary for improving the performance and understanding of catalytic mechanisms.

3.2 Single sites design for eCO₂RR

Ion-exchanged metals supported on porous materials such as metal-organic frameworks (MOFs) and zeolites, as well as organometallic complexes anchored to supports, can also be viewed as SACs.^{?,102–104} Furthermore the development of the concept of single-atom catalysis is confirmed with both atomic-resolution characterization techniques and theoretical modelling. The use of these methods provides a clear picture of the dispersion of the atoms, the bonding between single atoms and the support, and even in operando dynamics of single atoms during catalysis. The stabilizing effect of bonding between single transition metal atoms and/or ions and the support is not limited to materials containing oxo ligands. Other donor atoms can also be effective, particularly when carbon-rich supports are concerned. For example, N atoms in carbon-rich supports not only strongly anchor individual metal centres but also modify the electronic properties of the metal (see Figure 3.2). Relative to metals deposited on undoped supports, metal sites on N-doped carbon materials can exhibit higher activity and/or selectivity in a variety of electrocatalytic processes.

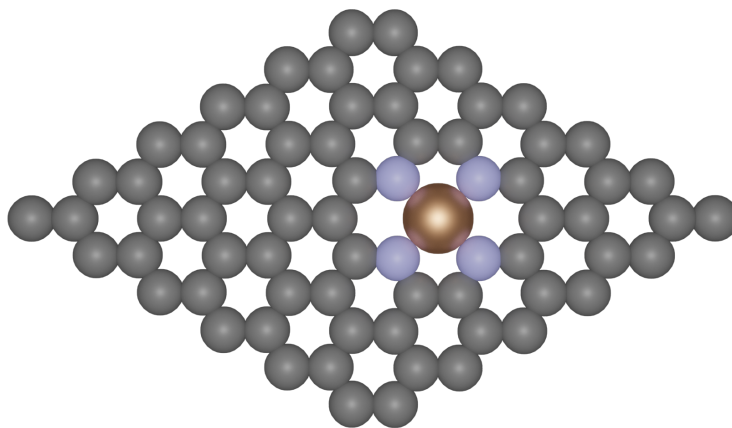


Figure 3.2: Model for carbon-rich single-site catalyst, particularly MNC materials.

In this thesis, the research focus on designing the models for (i) defined MOFs, (ii) various unsaturated N-doped carbon based materials, and (iii) molecular catalyst. MOFs have been one of the most researched type of single site catalyst.^{105–109} They can be constructed through the periodic coordination of metal ions and organic ligands to form one- to three-dimensional networks. The organic linkers in MOFs are enormously diverse and possess a variety of configurations. In addition, MOFs possess ordered pore structures, uniform

pore sizes and high specific surface areas, making them ideal substrates for the deposition of single atoms. For the preparation of single atoms, the pyrolysis of MOFs can serve as an important route to fabricate N-doped carbon materials, which are ideal substrates for the dispersion. Such materials were initially developed for catalyzing the oxygen reduction reaction (ORR), in particular in acidic medium, with focus on the FeN_x and CoN_x since other MN_x motives have much lower ORR activity.^{106,110,111} This has for a long time impeded the elucidation of the MN_x site structures and their ORR catalytic activities.¹⁰⁶ In the context of eCO₂RR, the presence of metal clusters along with MN_x moieties may, depending on the nature of the metal, catalyze the HER,^{112–115} and/or other pathways of eCO₂RR. This has prevented establishing clear experimental structure-to-property relationships between the nature of the metal in MN_x and their activity and selectivity toward eCO₂RR.¹¹⁶

Validation of DFT predictions on a reasonably large series of model materials is highly important to speed up the rational identification of more promising MNC candidates for eCO₂RR,¹¹⁷ given that the library of metal demonstrated to form MN_x sites in carbon has been recently expanded also to heavier elements such as Sn,¹¹⁸ Ru,¹¹⁹ Pt,¹²⁰ Pd,¹²⁰ etc. The investigation of model MNC catalysts exclusively comprising atomically-dispersed MN_x sites is thus critical to understand their structure-performance relationships and to identify the proper activity/selectivity descriptor(s). While numerous studies have now reported the eCO₂RR selectivity of some atomically-dispersed MNC catalysts prepared by different synthetic approaches,^{120–123} only one single study, to the best of our knowledge, has compared the trends in activity and selectivity for a narrow family of three MNC catalysts prepared identically except for the nature of the metal element.¹²⁴ The synthesis leaned on silica templating of o-phenylenediamine and Fe, Co or Ni salts. While metal particles formed in parallel to MN_x sites during the first pyrolysis, they could be removed during subsequent NaOH etching and acid-wash steps. However, the formation of metallic particles during the first pyrolysis lead to markedly different pore size distributions, in particular for CoNC. This rendered the comparison of the materials more complex due to expectedly different utilizations of the MN_x sites for different pore size distributions.

For an example, in Chapter 4 a robust synthesis approach involving a sacrificial Zn(II) zeolitic imidazolate framework and a metal salt (Mn, Fe, Co, Ni and Cu) was applied, which resulted in the complete integration of the metal atoms as atomically-dispersed MN_x moieties in the N-doped carbon matrix derived from the MOF, without formation of metallic particles. Diversity of these ligand pockets changes with a different type of synthesis depending on the reaction precursors and conditions. In addition, in Chapter 5 Ni saturated and unsaturated N-doped materials are investigated, where Ni-N coordination number can be tuned by changing pyrolysis temperatures in which different atomically dispersed Ni catalysts with different NiN_x numbers can be selectively prepared at 923 and 1223°C. Once the metal is embedded in to saturated and unsaturated N-doped cavities these single atoms can present different metal oxidation states depending on the cavity nature. Nickel Nitrogen Car-

bon (NiNC) single-atom catalyst exhibits the highest efficiency for producing CO, at different potentials.^{121,125–137} Molecular catalyst with a well-defined metal coordination counts as another type of SACs. Organometallic complexes can be anchored either covalently or non-covalently to the support. Macrocyclic ligands can be divided into different classes, such as phthalocyanine, porphyrin, cyclam, and so on. A number of transition-metal complexes have been utilized as electron-transfer mediators to achieve electrochemical reduction of CO₂.^{20,138–145} In Chapter 6 a model of Cobalt(II)-tetraphenyl porphine (CoTPP) complex deposited on the multiwalled carbon nanotube (MWCNT) is constructed along with Cobalt(II)-phthalocyanine (CoPc) anchored to carbon nanotube (CNT).

3.3 Basic concept of crystal field theory

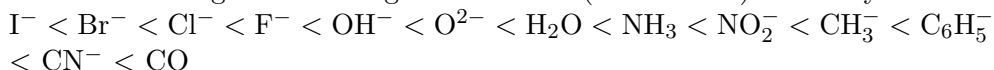
Single site carbon materials, which contain homogeneous-like defects, namely MN_x centers, resemble coordination compounds. The d-block metal (Mn, Fe, Co, Ni, Cu, and Zn) complexes can be described by bonding theories that rationalize experimental facts such as electronic spectra and magnetic properties. Once the geometry of the catalyst is determined (see section d-block geometries), computed magnetic data, particularly magnetization of embedded metal atom in SACs material, is used to determine unpaired electrons, where coupled with crystal field theory oxidation state of the metal atom is established. Continuing with the description of bonding in d-block metal complexes can be described with applications of valence bond theory, the electrostatic model and molecular orbital (MO) theory. However valence bond theory may rationalize stereochemical and magnetic properties, but only at a simplistic level. Furthermore, the model implies a distinction between high- and low-spin complexes that is actually misleading and it cannot tell us why certain ligands are associated with those formation. Molecular orbital and the electrostatic model concept differentiated once we move to complexes with ML π -bonding.

Therefore I used a second approach, the crystal field theory,^{146,147} which describes the bonding in coordination compounds of d-block metals. It is an electrostatic model and uses the electrons from the ligands to create an electric field around the metal centre. Ligands are considered as point charges and there are no metal-ligand covalent interactions. This may lead to a change in magnetic properties as well as color. Hybridization schemes are explained for octahedral, square planar, square pyramidal, and trigonal planar geometries, where applications of these hybridization schemes that can be used to describe the bonding in d-block metal complexes. The shapes of the molecular orbitals may be obtained without difficulty from the shapes of the fragment orbitals. Given the greater electronegativity of the ligands, the bonding MOs are more concentrated on them, whereas the antibonding MOs are more concentrated on the metal. Among them, seven involve contributions from the metal d orbitals: the two bonding MO 1e_g, the three nonbonding t_{2g} MO and the two

antibonding MO $2e_g$. The term 'd block' of an octahedral complex (and also for other types of complexes) is used to refer to the five molecular orbitals that are mainly concentrated on the d orbitals of the metal. The three nonbonding MO of t_{2g} symmetry clearly belong to the d block, since they are formed from pure d orbitals (xy, xz, and yz). For the MO of e_g symmetry, the bonding combinations ($1e_g$) are mainly concentrated on the ligands while the largest coefficient for the antibonding combinations ($2e_g$) is on the metal.

3.3.1 d-block geometries

An isolated metal atom has five d orbitals with the same energy. In **octahedral complex** the five fold degeneracy of the d orbitals is lifted, to give three nonbonding degenerate (t_{2g}) MO and two antibonding degenerate MO (e_g). The three nonbonding orbitals are then xy, xz, and yz, and the two antibonding orbitals are x^2-y^2 and z^2 (see Figure 3.3a). The energy separation (Δ^0) between the t_{2g} orbitals and the e_g orbitals depends on the strength of the σ interactions between the metal and the ligands. The value of this energy gap allows us to distinguish, in the family of octahedral complexes, strong field (large Δ^0) from weak-field complexes (small Δ^0). Measurement of the energies of d-d transitions allows us to estimate the value of Δ^0 in a large number of complexes, and to establish a spectrochemical series, in which the ligands are ranked according to the strength of the field (value of Δ^0) that they create:



Organometallic complexes, which contain one or several metal-carbon bonds, are thus strong-field complexes.

One can therefore consider, at least formally, that a **square-planar complex** is formed by removing two ligands from an octahedral complex, for example, those situated on the z-axis (see Figure 3.3b). To establish the structure of the d block, it is convenient to start from the results already obtained for octahedral complexes. Notice that while the three nonbonding orbitals xy, xz, and yz are degenerate from the energetic point of view, only two of them (xz and yz) are degenerate by symmetry. The main difference between the d blocks of octahedral and square planar complexes concerns the number of nonbonding or weakly antibonding orbitals: there are three in the former but four in the latter.

A complex which adopts a **square-based pyramidal** geometry, four ligands are located at the corners of a square which is the base of the pyramid, while the fifth ligand, or apical, is placed on the summit of the pyramid (see Figure 3.3c). In the octahedral complex, the coefficients of the xy, xz, yz, and x^2-y^2 orbitals are zero for the ligands located on the z-axis. The removal of one of these ligands therefore makes no change to the shape or energy of these four orbitals: in the d block of the square-based pyramidal complex, we therefore find that the three orbitals xy, xz, and yz are nonbonding, while x^2-y^2 is a strongly antibonding orbital. This result is just the same as that found when studying the square-planar geometry as a derivative of the octahedron.

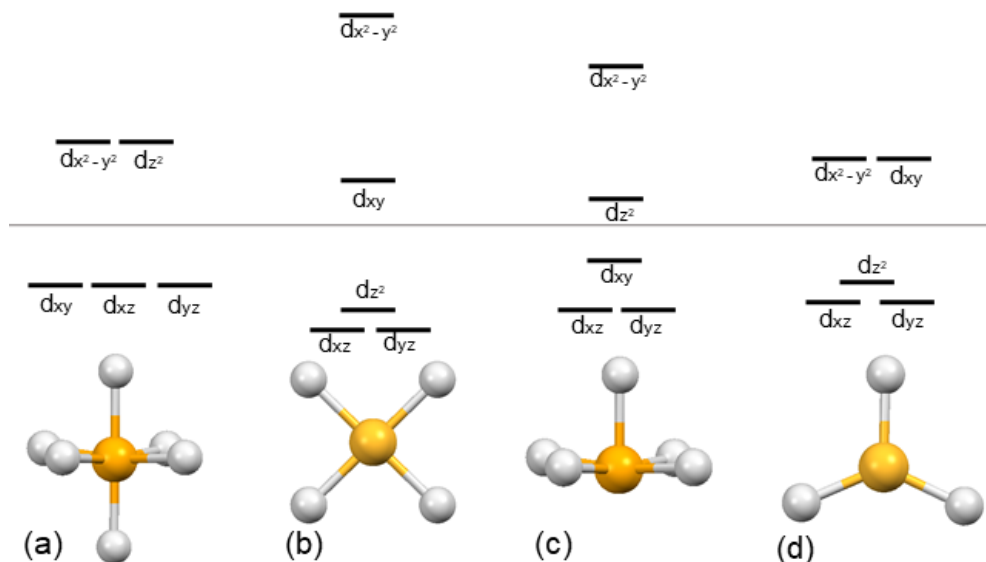


Figure 3.3: Derivation of the orbitals of the d block for (a) octahedral (b) square planar (c) square-based pyramidal and (d) trigonal-planar configurations. Gray line shows center of gravity. Ligands are represented with white spheres, and metal with yellow sphere.

The z^2 orbital is stabilized by the elimination of one of the two antibonding interactions with the ligands placed on the z-axis. But this stabilization is not as large as that observed when passing from an octahedral to a square-planar complex, since in this latter case, both the antibonding interactions along the z-axis are removed.

A **trigonal-planar** complex can be formed by removing the two axial ligands from a trigonal-bipyramidal complex. The d-block orbitals may therefore be readily deduced from those established in the preceding section. The variations in energy for the five d orbitals are shown in Figure 3.3d. The xz and yz orbitals of a trigonal-bipyramidal complex have zero coefficients on the axial ligands, as do xy and x^2-y^2 . The removal of these two ligands does therefore not change either the shapes or the energies of these four orbitals. But the z^2 orbital, which was strongly antibonding due to interactions with the axial ligands, is substantially stabilized by their removal, becoming almost nonbonding. Only three weak antibonding interactions with the ligands in the xy plane are left, and the amplitude of the z^2 orbital in this plane is small. The d-block for trigonal-planar geometry thus contains two nonbonding degenerate orbitals (xz and yz), a very weakly antibonding orbital (z^2), and two weakly antibonding degenerate orbitals (xy and x^2-y^2). These five orbitals are sufficiently low in energy to be doubly occupied.

Chapter 4

Volcano trend in electrocatalytic CO₂ reduction activity over atomically dispersed metal-sites on Nitrogen-doped carbon

In this project, theoretical and experimental results are joint to explain electrochemical carbon dioxide reduction reaction trends for Metal-Nitrogen-Carbon (M-N-C) like catalyst. High-throughput computing for series of atomically dispersed moieties MN_x (M = Mn, Fe, Co, Ni, and Cu) were performed. In order to model oxidation states of metal atoms ligands such as H₂O, OH, and O were added. Computational models then led to the identification of M²⁺N₄-H₂O as the most active centers in FeNC and CoNC, whereas Ni¹⁺N₄ was predicted as the most active one in NiNC at the experimental potentials of -0.5 and -0.6 V versus the reversible hydrogen electrode. The experimental activity and selectivity could be rationalized from DFT results, identifying in particular the difference between the binding energies for *CO₂⁻ and *H as a descriptor of selectivity toward CO. This work has been published in the Paper I and the optimized structures can be retrieved from ioChem-BD link (<https://doi.org/10.19061/iochem-bd-1-115>).

4.1 Methodology

4.1.1 Computational details

The valence electrons were expanded by plane waves basis sets with a kinetic energy cut-off of 500 eV. To model atomically-dispersed MN_x carbon structures, we used graphene layer expanded in a (6×6) supercell replacing four C atoms by nitrogens and removing two C atoms to compensate for the valence. These layers were interleaved by 12 Å vacuum along the Z direction. Various transition metal atoms (M = Mn, Fe, Co, Ni, Cu, Zn) were placed in the cavity left on the carbon sheet. In this configuration, two axial positions are empty. However, most of these atoms prefer octahedral coordinations, and thus ligands (L = none, H₂O, OH, O) typically occupy one of the axial positions in aqueous electrolyte, see Figure 4.1. The empty coordination site is then employed to investigate the eCO₂RR and competitive HER. Structures were relaxed with a force threshold of 0.05 eV/Å. The Brillouin zone was sampled using a 3×3×1 k-point mesh generated with the Monkhorst–Pack method.⁵⁵ The speciation of the different oxidation states of the metals in the carbon lattice was performed via a Pourbaix diagram derived with Phonopy.¹⁴⁸ The computed extended X-ray absorption fine structure (EXAFS) of MNC models were obtained with Demeter software.¹⁴⁹

4.2 Structural characterization

The Metal-Nitrogen-Carbon (MNC) materials were characterized by different experimental techniques. Multiple areas of the samples were examined, and no metallic particles could be observed. To confirm experimental findings I modelled MNC catalyst, where the single metal atom (M = Mn, Fe, Co, Ni, Cu, and (Zn)) is coordinated with four pyridinic nitrogen atoms, and the ligands (L = none, H₂O, OH, O) were used (see Figure 4.1). The (Zn)NC catalyst was used as a control sample and it was synthesized according to used synthesis process, except that no separate metal acetate was added. The models for ZnN₄- species did not present low binding energies of eCO₂RR intermediates. Extended X-ray absorption fine structure spectra were fitted with in-plane MN₄ structures, with an axial oxygen atom needed to obtain a best-fit only for FeNC, and obtained results were confirmed with computed EXAFS (see Figure 4.2). The metal-nitrogen bond distance and K-edge EXAFS spectra were also calculated ab initio for MN₄C model sites (10 in-plane carbon atoms in the second coordination sphere) in graphene sheets, with or without axial oxygen adsorbate. Considering water coordination via O atom, MnN₄ (2.40 Å), FeN₄ (2.33 Å), CoN₄ (2.31 Å) and (Zn)N₄ (2.35 Å) exhibited similar M-O bond distance between 2.31-2.40 Å. No water could be stabilized on the CuN₄ and NiN₄ models. The DFT-calculated EXAFS spectra shown in Figure 4.2 have characteristic features analogous to those of the experimental FT-EXAFS spectra with a main peak assigned to N or O back-scattering from the first coordination sphere.

The calculated and experimental M-N distances are also in good agreement, except for Fe, for which the EXAFS fitting is complicated by the presence of axial oxygen ligands (see Table 4.1).

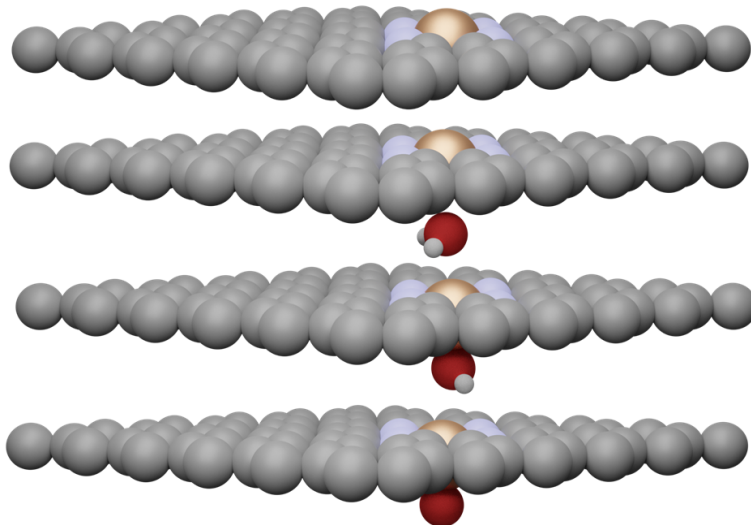


Figure 4.1: Computational models for the atomically-dispersed MN₄ sites, and when axial ligand is present (L = H₂O, OH, O). C is represented by grey spheres, N by purple, the metal in gold, O in red and H in white.

Table 4.1: Structural parameters obtained from the analysis with MN_x moieties of the metal K-edge EXAFS spectra and DFT calculations. Errors are given in brackets (1.89(1) means 1.88-1.90). CN stand for coordination number.

	$R_{(M-N)}/\text{\AA}$	CN	$R_{\text{DFT}(M-N)}/\text{\AA}$	CN_{DFT}
MnNC	1.89(1)	4	1.91	4
FeNC	2.01(2)	5/6	1.89	5
CoNC	1.96(2)	3.8(6)	1.89	4
NiNC	1.92(1)	4.3(3)	1.88	4
CuNC	1.94(1)	3.9(3)	1.93	4
ZnNC	2.00(1)	4.1(4)	1.96	4

The oxidation state of the atomically-dispersed metal centers before the reaction were investigated with X-ray photoelectron spectroscopy (XPS). They identify the coexistence of M²⁺ and M³⁺ species for FeNC, and the coexistence of M¹⁺ and M²⁺ species for CuNC.^{150,151} Only the 2+ oxidation state is observed in CoNC, NiNC and ZnNC, while for MnNC, the three different

oxidation states 2+, 3+ and 4+ seem to be present.^{152–154} The existence of different oxidation states of the metal in a same catalyst reveals a structural and/or a topological heterogeneity of MN_x moieties. For example, higher oxidation states of the metal cation might be triggered by oxygen (O₂, OH) adsorbates, which is possible only for MN_x moieties located on the top surface, while those located in the bulk of the nitrogen-doped carbon matrix are not accessible. Chemical heterogeneity might also lead to different oxidation states for a same structure. For example, while an average MN₄ structure was identified from our EXAFS analysis for all materials except FeNC (with additional oxygen adsorbates on top of the FeN₄ site itself), EXAFS cannot distinguish between N, O and C atoms in the first coordination sphere, implying that e.g. MN₂O₂ moieties or any other combination with a total of 4 in-plane ligating atoms among C, N and/or O are also possible structural candidates from an EXAFS viewpoint. For MnNC in particular, the presence of high valence-state metal (4+) suggests that there might be some O atoms replacing N atoms in MnN₄ moieties even during the reaction. This will be re-discussed later.

4.3 Selectivity and activity towards eCO₂RR

Over all studied MNC catalyst CO was the major eCO₂RR product produced, where the highest FEs were observed over the Fe-, Ni- and Mn-NC materials. High FEs have been reported previously for other FeNC^{155–158} and NiNC^{121, 127, 133, 159, 160} materials synthesized via different approaches but sharing a key common structural feature, namely the atomically dispersed MN_x moieties. Also in agreement with other recently published results, the major product over CoNC is H₂.^{124, 157, 161} On the contrary, high selectivity (more than 90%) toward CO in a wide potential range (-0.57 to -0.88 V vs. RHE) over atomically dispersed CoN_x was reported as well^{162, 163} but with slightly different coordination environment. Since the present series of MNC and (Zn)NC catalysts has, except for the nature of the metal element in MN_x moieties, negligible difference in morphology and structure, the different activity and selectivity toward eCO₂RR to CO can be mainly ascribed to the nature of the metal in MN_x sites. Last, (Zn)NC shows a selectivity for CO being in-between the low value of CoNC and the high values of MnNC, FeNC and NiNC. A catalytic synergy between MN_x and (Zn)N₄ moieties is also unlikely when the main considered product of CO₂ reduction is CO. The process requires only a 2 electron transfer and the steps are stated in Chapter 2, section explanation of electrocatalytic reactions and trends for reaction steps.

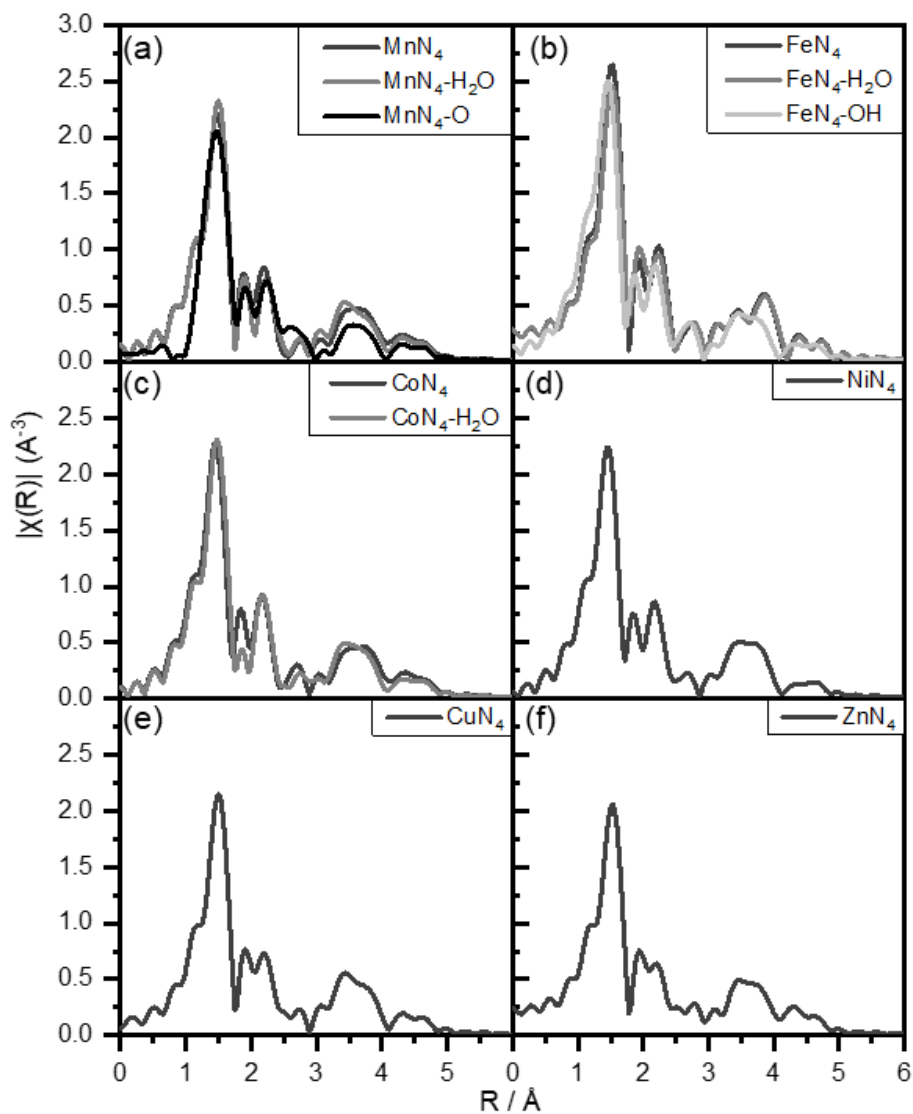


Figure 4.2: Simulated metal K-edge EXAFS. (a) MnNC, (b) FeNC, (c) CoNC, (d) NiNC, (e) CuNC, and (f) (Zn)NC. For Mn, Fe, and Co all the relevant computed structures are presented.

If a synergistic effect between MN_x and (Zn)N₄ moieties occurs, there are two possibilities: a) MN_x catalyzes step 1 ($* + \text{CO}_2 + \text{H}^+ + \text{e}^- \rightarrow * \text{COOH}$) and (Zn)N₄ catalyzes step 2 ($* \text{COOH} + \text{H}^+ + \text{e}^- \rightarrow * \text{CO} + \text{H}_2\text{O}$); or b) the opposite. Assuming the synergy case a), the FE for formate on FeNC, CoNC and NiNC should be significantly higher than that on (Zn)NC, which is not the case (Figure 4.3a). Assuming the synergy case b), (Zn)NC should have a J_{HCOO⁻} value comparable to J_{CO} observed on the active catalysts NiNC, FeNC and CoNC. However, this is also in contrast with experimental results, where J_{HCOO⁻} on (Zn)NC being negligible compared to J_{CO} on NiNC, FeNC and CoNC (Figure 4.3b). Thus, we can rule out any significant catalysis synergy between MN_x and (Zn)N₄ moieties.

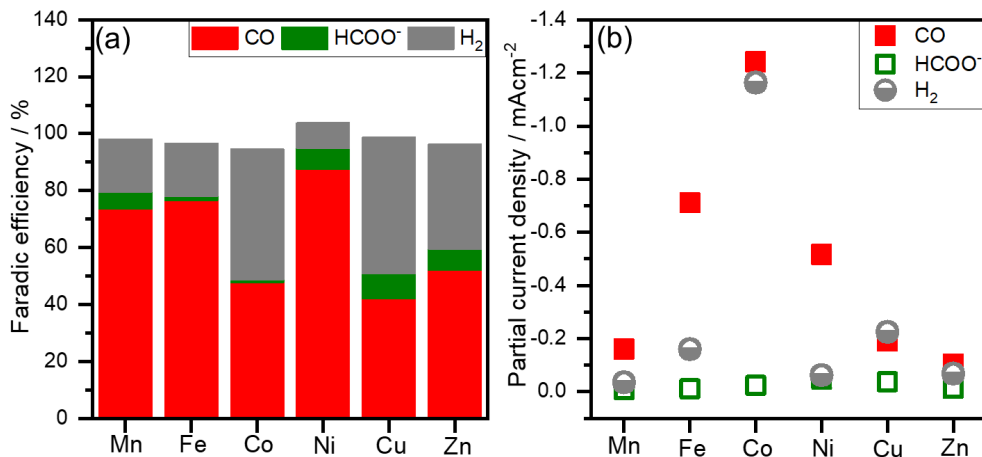


Figure 4.3: Product analysis over the MNC catalysts (M = Mn, Fe, Co, Ni, and Cu), plotted against the embedded metal center sorted by the atomic number. (a) Faradaic efficiency and (b) partial current densities at -0.6 V vs. RHE for CO, HCOO⁻ and H₂ products. The data were obtained by Chronoamperometry (CA) for 90 min, in 0.1 M KHCO₃ saturated with CO₂ (pH 6.7) at room temperature.

Operando X-ray absorption near edge structure (XANES) spectroscopy was employed to track possible changes in the oxidation state of MN_x sites during eCO₂RR. The DFT simulations were performed on the most likely oxidation state and environment for each metal. Combined with the ex situ XPS, we can infer that the oxidation state of Fe was modified from 2+/3+ to only 2+, that of Ni reduced from 2+ to 1+, and that of Cu reduced from 2+/1+ to 0.^{161,164} The oxidation state of Mn remained a mixture of 2+/3+/4+, and that of Co remained as 2+, even at -1.0 V vs. RHE. It can further be observed that Fe was reduced from a mix of 2+/3+ to only 2+ already at 0 V vs. RHE, in line with the redox potential identified for FeNC at ca 0.75 V vs. RHE in acidic electrolyte.^{165,166} This may explain the high onset potential for eCO₂RR over FeNC, while NiNC is possibly inactive for eCO₂RR as long as it is in 2+ oxidation state. The changes in oxidation states of the three most active MNC catalysts (namely Fe, Co and Ni) with applied potential agree well with the calculated Pourbaix speciation diagrams (see Figure 4.4), which

show that in a pH 7 electrolyte, the oxidation state of Fe and Co in MN_4 corresponds to M^{2+} at any potential 0 V vs. RHE, while there is a reduction from Ni^{2+} to Ni^{1+} below 0 V vs. RHE at pH 7, in line with the Ni-reduction experimentally observed between 0 and -0.5 V vs. RHE experimentally (see Figure 4.5).

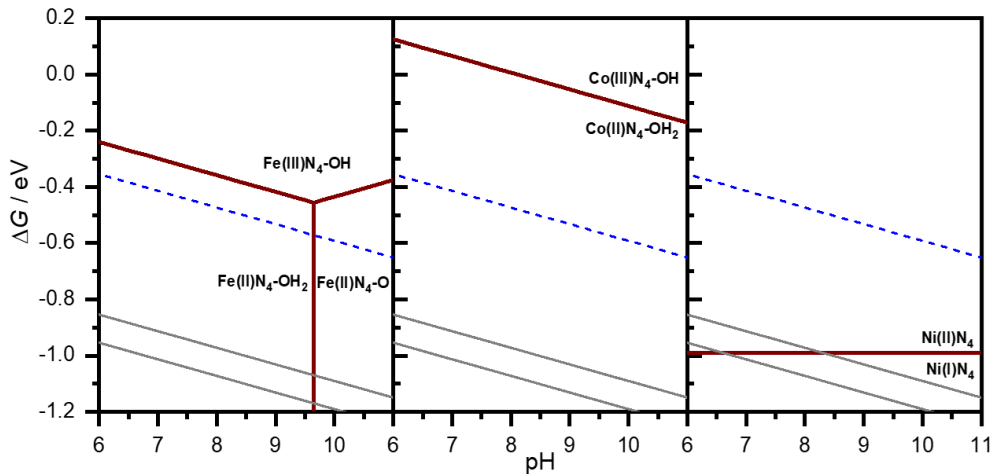


Figure 4.4: Simulated Pourbaix diagrams. Computed Pourbaix diagram for FeN_4^- , CoN_4^- and NiN_4^- species. The blue lines corresponding to the water oxidation reactions, the wine red lines separate the different phases, and the area between two gray lines corresponds to the experimental potential window at which eCO_2RR was evaluated (-0.5 and -0.6 V vs. RHE).

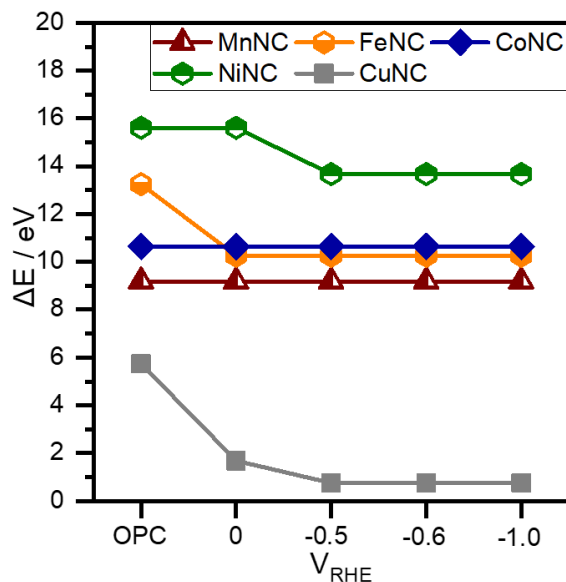


Figure 4.5: Change in the absorption edge of metal K-edge XANES spectra of MNC catalysts with applied potential. The electrolyte was CO₂-saturated 0.1 M KHCO₃, and the OCP was around 0.7 V vs. RHE. The x-axis is the difference between the X-ray energy and the absorption edge of MNC (E_0).

4.4 Explanation of observed volcano trends

The observation of volcano trends in the electrochemical activity reactions among a series of metal-based catalysts with similar site structures is generally explained by Sabatier's principle. With the binding energy of the metal sites being tuned by the metal's atomic number, and leading to too weak, optimum and too strong binding of the key intermediate adsorbed species when moving in a given direction in the series of transition metal.^{167,168} We performed DFT calculations with MN_4C model ($M = Mn, Fe, Co, Ni$ and Cu , with axial oxygen adsorbates for all the metals) with the CHE model at -0.6 V vs. RHE. The reaction energy profiles for each elementary step of eCO_2RR show that the particular transition metal present in the N_4 lattice has a tremendous effect on reactivity. I will focus on the Gibbs free energy diagrams at -0.6 vs. RHE (Figure 4.6a and 4.7), potential at which the eCO_2RR activity and selectivity of MNC catalysts are experimentally measured. Overall, the calculated Gibbs free energy diagrams demonstrate that MnN_4C and FeN_4C bind CO too strongly and their CO production rate is determined by the *CO desorption step. In contrast, NiN_4C and CuN_4C have weak CO binding energy, and the CO_2 activation and first electron transfer then becomes the RDS (Table 4.2). In this series of model sites, CoN_4 shows the optimum balance between the energy barriers for CO_2 activation and CO desorption (steps 1 and 4 (desorption) in Table 4.2, respectively).

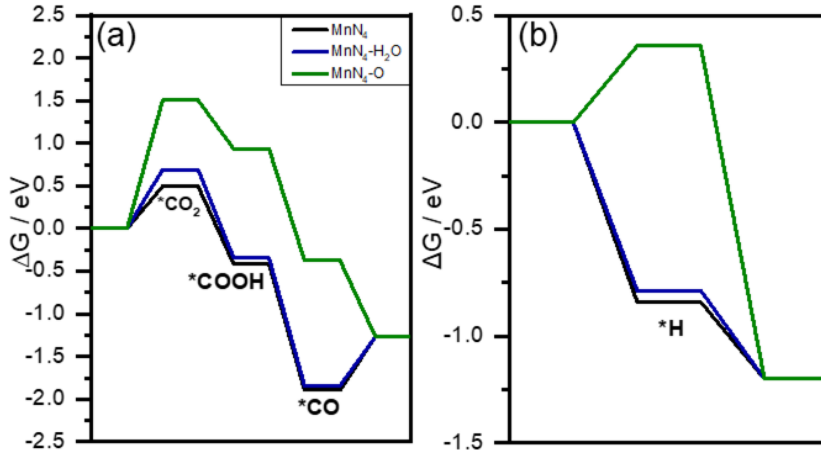


Figure 4.6: Gibbs free energy profile over MnN_4 - model for (a) eCO_2RR and (b) HER at -0.6 V vs. RHE (pH 7) according to the different oxidation state and ligand coordination.

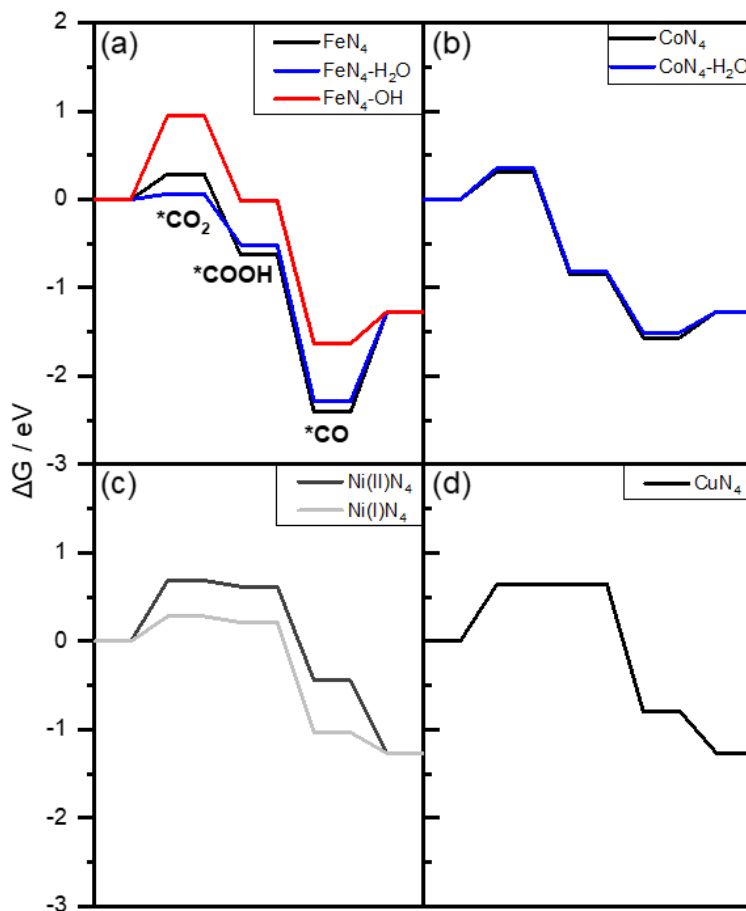


Figure 4.7: Gibbs free energy profiles for eCO₂RR over (a) FeN₄-, (b) CoN₄-, (c) NiN₄-, and (d) CuN₄ model sites at -0.6 V vs. RHE (pH 7) according to the different oxidation state and ligand coordination.

We have also investigated how axial adsorbates (H₂O, OH and O) affect the reaction energies on MnN₄, FeN₄ and CoN₄ (no oxygen adsorbates could be stabilized on CuN₄ and NiN₄). Particularly, the H₂O molecule axially adsorbed on FeN₄ lowers the reaction barrier for the first electron transfer to form $*CO_2$ (Figure 4.7a) and thus promotes the eCO₂RR to CO over FeNC. In this case, FeN₄ sites are approachable from both axial directions (one side for adsorbing H₂O, while the other side for activating CO₂ gas molecule). Moreover, for the case of NiNC, NiN₄C with oxidation states of 1+ and 2+ are both used in the DFT simulations, since a reduction of Ni²⁺ to Ni¹⁺ was observed from operando X-ray absorption spectroscopy (XAS) (Figure 4.5). In contrast to the high energy barrier for the CO₂ activation on Ni²⁺N₄C (the first step for eCO₂RR), Ni¹⁺N₄C binds CO₂ more strongly (Figure 4.7c). This explains that the onset potential of eCO₂RR over NiNC is highly correlated to the Ni²⁺/¹⁺N₄ redox potential (Figure 4.5). The CO₂ activation energies are also highly correlated to the geometries of the CO₂ coordination over MN₄ sites, as depicted in Figure 4.8 showing how CO₂ is activated in the Fe, Co and

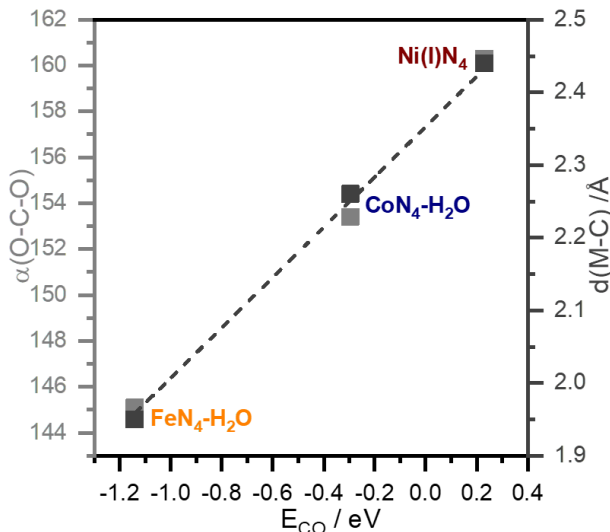


Figure 4.8: Linear scaling relation. Linear angle $\alpha(\text{O-C-O})$ and distances $d(\text{M-C})$ of CO₂ scaling relations for adsorption of CO molecule. Fe presents an $\eta^2\text{-C,O}$ hapticity, while Co and Ni prefer the $\eta^1\text{-C}$ configuration.

Ni sites. While both a M-C and a M-O bond can simultaneously form with the FeN_4 active site, only a M-C bond appears for the CoN_4 and NiN_4 sites. In addition, the O-C-O angle departs from the 180° value in the gas phase to only 145° when adsorbed on $\text{FeN}_4\text{-H}_2\text{O}$. This explains the lower activation energy of CO₂ over FeNC vs. other MNC catalysts and the higher onset potential for FeNC . The activity trends for CO₂ reduction to CO can be estimated from DFT calculations on the basis of the exergonic value of the reaction energy at the RDS. Table 4.2 reports the reaction energy calculated for each elementary step for $\text{MN}_4\text{-H}_2\text{O}$ moieties according to the Gibbs free energy profiles for eCO₂RR over MNC catalysts ($M = \text{Mn, Fe, Co, Ni}$ and Cu) at -0.6 V vs. RHE. For each MN_4 site, the structure (with/without H₂O adsorbate) with lowest reaction energy at the RDS is considered in Table 4.2. To compare DFT results to the experimental eCO₂RR activity trend, we assumed an Arrhenius type law of the form $A \exp(-E_{\text{RDS}}/RT)$, where E_{RDS} is the free energy change at the RDS and A is a pre-exponential factor, assumed the same for all MNC catalysts investigated due to similar metal content and similar carbon porosity, chemistry and morphology.

Table 4.2: Gibbs free energies at -0.6 V vs. RHE for the elementary steps of eCO₂RR to CO. The RDS values correspond to energies corrected by CO and CO₂ gas entropies. All metal cations in 2+ oxidation state unless otherwise indicated. FeN₄-H₂O with a high spin state (S=2) is used

	Step 1	Step 2	Step 3	Step 4	RDS
MnN₄	0.49	-0.91	-1.48	0.63	0.63
MnN₄-O	1.51	-0.57	-1.31	-0.81	1.51
FeN₄-H₂O	0.06	-0.58	-1.77	1.02	1.02
CoN₄-H₂O	0.36	-1.17	-0.69	0.24	0.36
NiN₄	0.68	-0.07	-1.06	-0.83	0.68
Ni(I)N₄	0.28	-0.07	-1.25	-0.24	0.28
CuN₄	0.64	0.01	-1.44	-0.48	0.64

Figure 4.9 shows that the trend of reaction rate predicted from DFT according to this approach matches well the experimentally observed volcano trend in the present work, except for MnN₄ and CuN₄. Among all the MN₄(-H₂O) active sites, Ni¹⁺N₄ is predicted to show the highest activity toward eCO₂RR to CO, with lowest maximum for the free reaction energy at any of the elementary steps. However, the experimental oxidation states of NiN₄ sites are most likely in a mix of 1+/2+ at -0.6 V vs. RHE (see Figure 4.5). Thus, the experimental activity of NiNC should fall between those predicted for Ni¹⁺N₄ and Ni²⁺N₄. Experimentally, the eCO₂RR activity of NiNC is even lower than that estimated from DFT for Ni²⁺N₄. We assign this to a lower onset potential for NiNC than for CoNC, but once the reaction has started, recent studies have shown that NiNC catalysts with single metal atom sites are the most active and selective MNC catalysts reported to date.^{7, 133, 159, 160, 169} The calculated activity trends for MN₄ sites are as following FeN₄-H₂O < MnN₄ < CoN₄-H₂O < Ni¹⁺N₄ > Ni²⁺N₄ ≈ CuN₄, with metals on the left hand-side of Co or Ni binding CO too strongly, and metals on the right hand-side binding CO₂ too weakly.

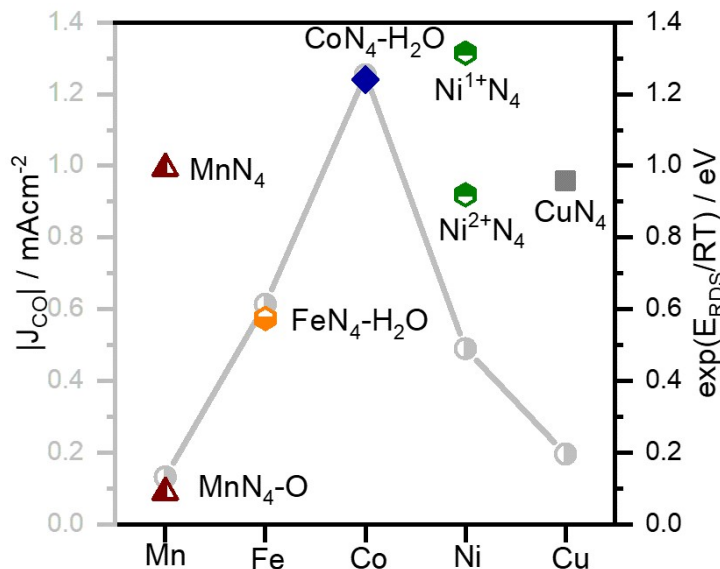


Figure 4.9: Relation between theoretical and experimental data. The comparison of CO partial current density at -0.6 V vs. RHE for pyrolyzed MNC materials and DFT calculated trends (at -0.6 V vs. RHE). E_{RDS} is the calculated free reaction energy at the specific RDS of $e\text{CO}_2\text{RR}$ to CO for each active site (see Table 4.2).

One discrepancy between the experimental and theoretical results is MnNC, for which our experimental data show much lower CO partial current density than FeNC while the DFT calculations predicted that MnN_4 and FeN_4 sites are both limited by the $^*\text{CO}$ desorption step, but with lower desorption barrier for MnN_4 . This results in a predicted $e\text{CO}_2\text{RR}$ activity higher for MnN_4 than FeN_4 moieties (Figure 4.9). Given that a higher oxidation state of Mn (a mixture of 4+, 3+ and 2+) was observed by ex situ XPS and operando XANES than can be assumed with a MnN_4 model site, and given that MnN_x has a higher predicted affinity to oxygen,¹⁵⁰ we re-evaluated the possible active site structure of MnNC. We assumed MnN_2O_2 as an alternative model site (Figure 4.10b). Both MnN_4 and MnN_2O_2 structures converged and optimized under the same conditions, confirming similar energy stabilities of these two structures. Moreover, the experimental FT-EXAFS spectrum of MnNC could be properly fitted with the MnN_2O_2 model site, and its DFT calculated EXAFS spectrum exhibits the same characteristic peak at 1.5 Å (Figure 4.10a) assigned to both N and O back-scattering from the first coordination sphere. According to our DFT calculations, this MnN_2O_2 model can not coordinate CO_2 and is therefore predicted to be inactive toward $e\text{CO}_2\text{RR}$. Even with a single oxygen adsorbate, the MnN_4 site becomes poorly active toward $e\text{CO}_2\text{RR}$, as indicated by the very high energy barrier for CO_2 adsorption for the $\text{MnN}_4\text{-O}$ model (Figure 4.6a). Of all the Mn models investigated, only Mn^{2+} without oxygen ligands can effectively coordinate CO_2 . However, as identified by ex situ XPS and operando XANES (see Figure 4.5), this species is uncommon at the investigated potentials, explaining why

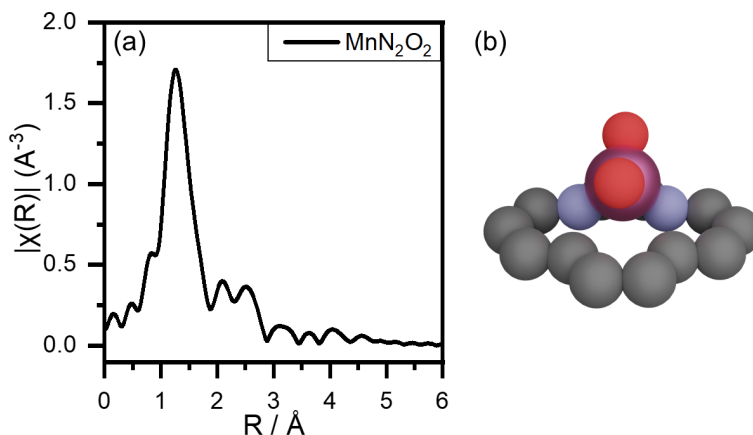


Figure 4.10: MnN₂O₂ site. (a) Computed simulated metal K-edge EXAFS spectrum and (b) computational model of MnN₂O₂ model site.

no activity is experimentally found for MnNC. Last, for CuNC, our operando XANES data show that Cu switches from 2+ to zero oxidation state early on (Figure 4.5), explaining the mismatch between DFT predictions and the experimental CuNC activity. In summary, with the MnN₄-O model site for MnNC and MN₄(-H₂O) model sites for other MNC catalysts (M = Fe, Co, Ni), the experimental volcano trend in the eCO₂RR to CO activity of MNC catalysts can be nicely explained by our DFT calculated Gibbs free energy for the RDS (Figure 4.9).

4.5 Selectivity toward CO over MN₄

As another critical factor for a good eCO₂RR catalyst, selectivity is highly correlated to the *H and *CO₂ binding energies over MN_x active sites. Thus, we investigated the Gibbs free energy profiles of HER over MN₄C model sites (M = Mn, Fe, Co, Ni and Cu) at -0.6 V vs. RHE (Figure 4.6b and 4.11). These binding energies can be rationalized from simple coordination chemistry rules. The adsorption behaviour for the metals in the first transition series depends on the occupation of antibonding d-orbitals coming from the crystal field splitting for the resting square pyramidal structure. Therefore, we use the number of d-antibonding electrons (0 for FeN₄-H₂O, 1 for CoN₄-H₂O, 2 for Ni¹⁺N₄ and 3 for Ni²⁺N₄) to explain the selectivity toward eCO₂RR to CO over MN₄ sites (Figure 4.12). The population of the antibonding states results from counting the number of electrons in the d-orbitals of the ground state configuration according to the crystal field splitting theory. Figure 4.12 indicates that there is a linear scaling relationship between the *CO₂ and *CO binding energies, which determines the CO activity. Thus, the CO binding energy can serve as a key descriptor for eCO₂RR to CO activity.

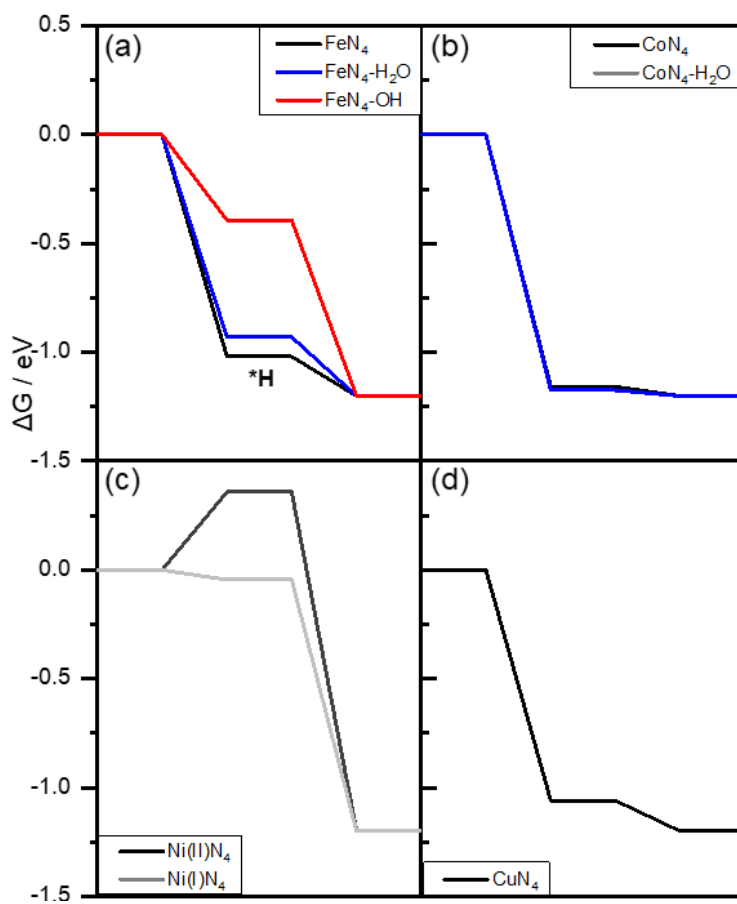


Figure 4.11: Gibbs free energy profiles for HER over (a) FeN₄-, (b) CoN₄-, (c) NiN₄-, and (d) CuN₄ model sites at -0.6 V vs. RHE (pH 7) according to the different oxidation state and ligand coordination.

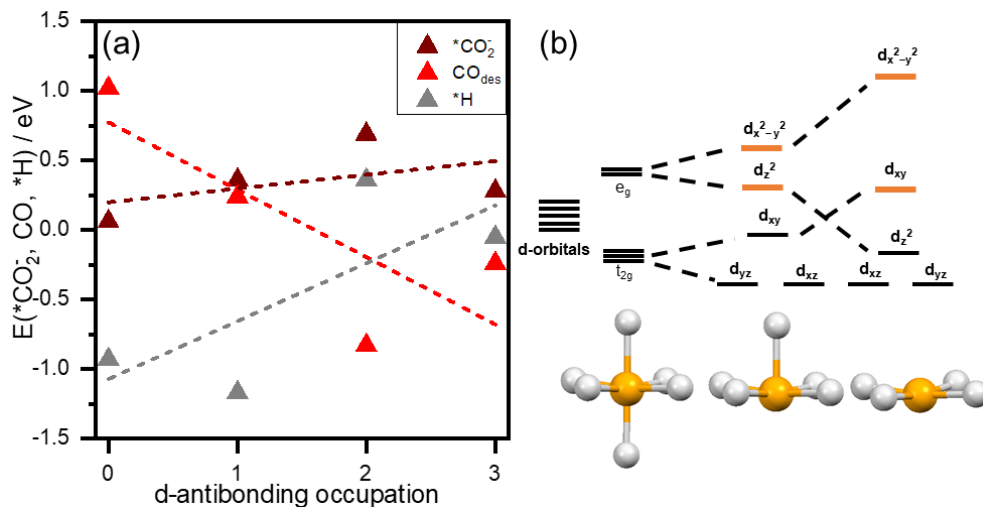


Figure 4.12: Relationship between the adsorption/desorption energies and the local electronic configurations on the Fe, Co and Ni (1+ and 2+). (a) $*\text{CO}_2^-$ (wine), $*H$ (grey) adsorption energies and CO (red) desorption (at pH 7 and -0.6V) as a function of the d-antibonding occupation in the active site. In (b) the symmetry descent from the octahedral configuration to either square pyramidal or square planar are shown (orange indicates the antibonding states). For the most common Fe species under potential Fe and Co are in $\text{MN}_4\text{-H}_2\text{O}$ (square pyramidal with 0 and 1 antibonding electrons). Instead Ni is in the square planar configuration with already two electrons in the antibonding state. At high negative potential the highest antibonding $d_{x^2-y^2}$ would be occupied, which could trigger the distortion of the planar pocket (Ni^{1+} state).

The volcano trend with CO binding energy as a descriptor (J_{CO} measured experimentally vs CO binding energy calculated from DFT simulations) is shown in Figure 4.13, with Fe located at the strong binding branch, whereas Ni and Cu are on the weak binding branch. $\text{CoN}_4\text{-H}_2\text{O}$ sites possess the optimum binding energies with $^*\text{CO}$, which are neither too strong nor too weak. On the other hand, good selectivity toward CO requires high $^*\text{CO}_2^-$ binding energy and low $^*\text{H}$ binding energy (the more negative value in Figure 4.12a, the higher the binding energy). Thus, we utilized the difference of $^*\text{CO}_2^-$ and $^*\text{H}$ binding energies ($E(^*\text{CO}_2^-) - E(^*\text{H})$) as the descriptor for selectivity (Figure 4.14). In that figure, we show only the experimental data for $\text{FeN}_4\text{-H}_2\text{O}$, $\text{Ni}^{2+/1+}\text{N}_4$, and $\text{CoN}_4\text{-H}_2\text{O}$ and did not include the other MNC materials, whose eCO_2RR activity and selectivity can originate in whole or in part from metal-free nitrogen functional groups.¹⁷⁰ The $\text{CoN}_4\text{-H}_2\text{O}$ site has an optimum binding energy with $^*\text{CO}_2^-$, but, unfortunately, the binding energy of $^*\text{H}$ is much higher than that of $^*\text{CO}_2^-$ which leads to a low selectivity toward CO, in line with the experimental result on CoNC (48%, Figure 4.14).

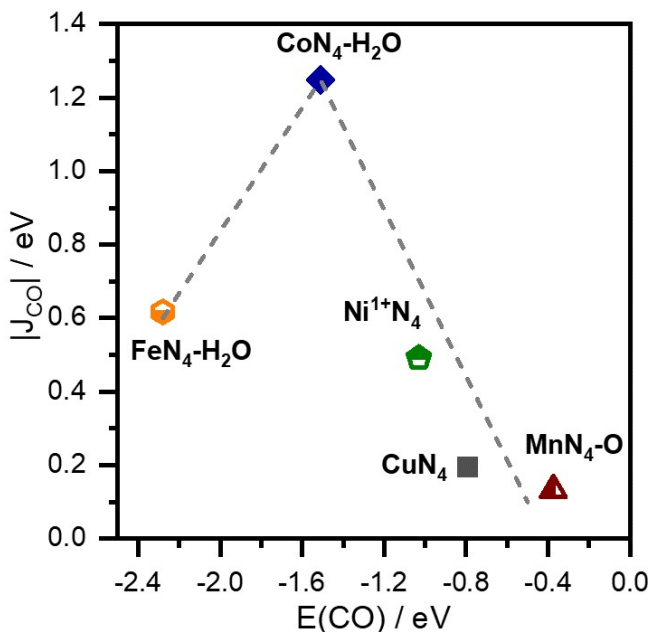


Figure 4.13: Volcano plot of CO partial current density at -0.6 V vs. RHE with CO binding energy obtained from DFT simulations as descriptor.

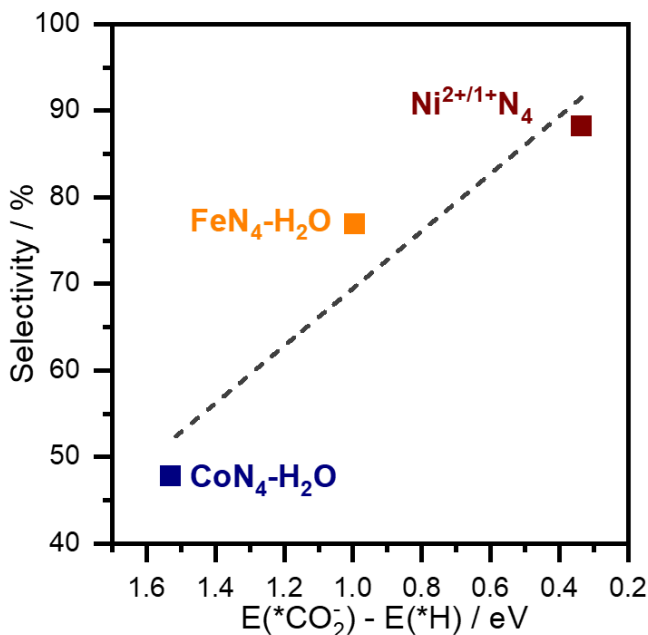


Figure 4.14: Selectivity analysis. The experimental selectivity at -0.6 V vs RHE over FeNC, CoNC, and NiNC vs. the DFT simulated $E(*CO_2^-) - E(*H)$.

The high selectivity of eCO₂RR to CO over NiN₄ benefits from the lower *H binding energies compared to other metal MN₄ sites (Figure 4.14). For FeN₄-H₂O, even though the high binding energy to *H suggests a high HER activity (Figure 4.12), the reduction of CO₂ to *CO also occurs with zero energy barrier, and the binding energy of *CO on FeN₄-H₂O is much stronger than that of *H on FeN₄-H₂O. Thus the majority of FeN₄ sites are occupied by *CO in CO₂-saturated electrolyte. This site-blocking effect probably suppresses the binding of such active sites with HER related species, resulting in high experimental selectivity toward CO over FeNC. As can be seen in Figure 4.14, the selectivity of eCO₂RR toward CO over MN_x moieties can be explained with the DFT-calculated *CO₂ and *H binding energies, that in turn are highly correlated to the number of d-antibonding electrons. The latter descriptor can thus be used to screen for more promising eCO₂RR candidate MNC materials.

4.6 Conclusions

We have employed Density Functional Theory to understand experimental results and establish clear correlations between physicochemical and catalytic properties for the eCO₂RR of atomically-dispersed MN₄ centres (M = Mn, Fe, Co, Ni, and Cu). With computed Pourbaix diagrams we supported the ex situ XPS and in situ XANES results that were used to identify and monitor changes of the metal oxidation state at rest and under operating conditions, respectively. Mn and Co did not change oxidation state down to -1.0 V vs. RHE while Fe and Ni were partially reduced and Cu was largely reduced to metallic state. Over these highly controlled systems, catalytic tests revealed a volcano-like dependency between eCO₂RR activity and the atomic number of the transition metal, with Fe and Co as the most active centres. Computational models led to the identification of active centres and finally to descriptors rationalizing the distinct activity and selectivity patterns observed. To this end, different configurations were studied according to the speciation of the metal atom, its charge, spin, and coordination sphere. Among the best performing materials, Fe²⁺N₄-H₂O and Co²⁺N₄-H₂O centres were found as the most likely active catalytic centres at considered potentials (-0.5 and -0.6 V vs. RHE), whereas Ni¹⁺N₄ was predicted as the most active Ni-based one. The Gibbs free energy change at the RDS accounted for the experimentally determined activity volcano, with Co-, Fe-, and Ni-based systems showing the best compromise between *CO and *CO₂⁻ binding energies. Regarding selectivity, the difference between the binding energies for *CO₂⁻ and *H, directly related to the number of d-antibonding electrons, described successfully the high selectivity observed on Fe- and Ni catalysts and the more favoured HER observed over the Co-based one.

Chapter 5

Stability and redispersion of Ni nanoparticles supported on N-doped carbons for the electrochemical CO₂ reduction

The Nickel-Nitrogen-Carbon (Ni-N-C) catalyst exhibits the highest efficiency for producing CO, at different potentials. The variation in the material synthesis produces defects with coordinatively saturated and unsaturated N-doped cavities, and once the metal is placed there these single atoms can present different metal oxidation states depending on the cavity nature. The research in this study was based on experimental findings from article^[137]. We explored the possibility of Ni nanoparticle redispersion by Ni(CO)₂-like disintegration, when supported on saturated and unsaturated N-doped carbon materials. The binding energies of intermediates towards CO on seven different models (NiN₄, NiN₃C, NiN₃, NiNC₃, NiNC₂, NiC₄, NiC₃, Ni(100), Ni(111), and Ni(211)) were obtained from DFT taking the most stable conformations for each intermediate on each support. Besides nanoparticle shape was obtained from Ni(100), Ni(111), and Ni(211) facets. This work has been published in the Paper 2 and the optimized structures can be retrieved from ioChem-BD link (<http://dx.doi.org/10.19061/iochem-bd-1-168>).

5.1 Methodology

5.1.1 Computational details

This project is divided in two parts. The first part contains calculations of various N-doped carbon materials (see Figure 5.1), where the kinetic energy cut-off was set at 500 eV. N-doped carbon materials are represented with a graphene layer expanded in a (6×6) supercell and replacing carbon atoms by nitrogen atoms to design desired models. These layers were interleaved by 12 Å vacuum along the Z direction. A nickel atom was placed in the cavity left on the carbon sheet. Structures were relaxed with a force threshold of 0.050 eV/Å. The Brillouin zone was sampled using a (3×3×1) k-point mesh generated with the Monkhorst–Pack method.¹⁷¹

The second part focuses on Ni nanoparticles represented by three metal surfaces with the lowest Miller indexes Ni(100), Ni(111), and Ni(211). High quality surface energy calculations ((10×10×1) k-point mesh) were performed to build the nanoparticle structure using the Wulff model^{172,173} with VESTA.¹⁷⁴ The vacuum between slabs was larger than 10 Å and the structures were relaxed with a force threshold of 0.025 eV/Å. They contained five layers, where the two uppermost were fully relaxed and the rest were fixed to the bulk structure.

CO can adsorb in very dense phases on Ni^{175,176} thus different coverages were considered: from 1/12 to 1 ML for Ni(100) and Ni(211) and from 1/16 to 1 ML for Ni(111). The Brillouin zone was sampled using a (3×4×1), (3×3×1) and (5×3×1) k-point meshes for Ni(100), Ni(111) and Ni(211) respectively. The Wulff construction was also employed on the CO-dense phases to explain the restructuring of the nanoparticle. Ni(CO)_x species can be formed in rich CO environments, which can then redisperse and reconstruct the Ni nanoparticle. The Ni(211) surface edge has enough space to accommodate more than one CO molecule, allowing the formation of Ni(CO)₄ complexes. We thus calculated the formation and elimination of Ni(CO)₄ with different CO coverages on that surface, where the elimination of this species forms a vacancy at the edge of the Ni(211) steps.

5.1.2 Ostwald disintegration formalism

We applied the energetic formalism from Ouyang and co-workers¹⁷⁷ to evaluate whether the disintegration of Ni(CO)_x from Ni NPs is possible on N-doped carbon materials.

Energetics of Ni nanoparticle

To study this, we start from the energetics of a metal particle in the absence of reactants. Average energy ΔE_{NP} (per atom) with respect to infinite size particle (bulk) can be calculated by

$$\Delta E_{\text{NP}} = \frac{3\Omega\gamma}{R} \quad (5.1)$$

where Ω is the molar volume of bulk metal atom, γ is the overall surface energy of the Ni nanoparticle, and R stands for radius of the nanoparticle. Considering that metal particles may expose different facets i with surface energy γ_i and corresponding area ratio f_i over the whole surface area, the overall surface energy γ could be rewritten as:

$$\gamma = \sum_i f_i \times \gamma_i \quad (5.2)$$

The chemical potential $\Delta\mu_{\text{NP}}$ of the supported metal particle can be derived

$$\Delta\mu_{\text{NP}} = \frac{2\Omega\gamma}{R} \quad (5.3)$$

Under reaction conditions, CO may adsorb on the Ni particles, and corresponding Gibbs free energy of the adsorption would reduce the surface energy and stabilize the metal particles. The reduction of surface energy $\Delta\gamma_i$ on the facet i at given temperature T and partial pressure p can be calculated

$$\Delta\gamma_i(T, p) = \theta \times [E_{\text{CO}}(\theta) - \Delta\mu_{\text{CO}}(T, p)]/A_i \quad (5.4)$$

where θ and A_i is the coverage of adsorbed CO and surface unit area of the facet i . $E_{\text{CO}}(\theta)$ is the average binding energy of CO and coverage dependent. The effective surface energy γ' of supported metal particles with CO adsorbate becomes

$$\gamma'(T, p) = \sum_i f_i [\gamma_i + \Delta\gamma_i(T, p)]. \quad (5.5)$$

By substituting γ_i in equation 5.1 and 5.3 average energetics $\Delta E'_{\text{NP}}$ and chemical potential $\Delta\mu'_{\text{NP}}$ of supported metal particles under reactants can be obtained, respectively.

CO induces disintegration

Following the criteria of reactant induced dynamics, the stability of Ni nanoparticle can be explained by the Gibbs free energy of disintegration ($\Delta G_{\text{NP}}^{\text{dis}}$). First, we compute the formation energy of isolated metal adatom (E_{Ni}^{f}) with respect to Ni bulk reference supported on the different carbon defects by

$$E_{\text{Ni}}^{\text{f}} = E_{\text{sur-Ni}} - E_{\text{sur}} - E_{\text{B}} \quad (5.6)$$

where $E_{\text{sur-Ni}}$ is the energy of anchored Ni on the support, E_{sur} is the energy of support, and E_{B} is the energy of Ni bulk reference. From the results (see Table 5.1), we can see that formation energies are exothermic, as Ni atoms are anchored to a defect. Second, we calculate the energy of the Ni atom coordinated to the N-doped carbon substrate, where they should be able to form a chemical bond with the Ni atom, and corresponding binding energy $E_{\text{nCO}}^{\text{f}}$ must be negative.

$$E_{\text{nCO}}^{\text{f}} = E_{\text{sur-nCO}} - E_{\text{sur-Ni}} - n \times E_{\text{CO}} \quad (5.7)$$

Table 5.1: Formation Energy of Ni anchored atom (E_{Ni}^{f}) (with respect to Ni bulk energy), Binding Energy of CO on Ni atom ($E_{\text{nCO}}^{\text{f}}$), and Formation Energy of the Ni-CO complexes on the support with respect to the infinite size metal particle (CO in gas phase) ($E_{\text{surf-nCO}}^{\text{f}}$). All energies in eV.

	E_{Ni}^{f}	$E_{\text{1CO}}^{\text{f}}$	$E_{\text{surf-1CO}}^{\text{f}}$	$E_{\text{2CO}}^{\text{f}}$	$E_{\text{surf-2CO}}^{\text{f}}$
NiN ₃	-1.97	-0.55	-2.51	-0.84	-2.80
NiNC ₂	-2.71	-0.89	-3.60	-1.33	-4.04
NiC ₃	-3.51	-0.72	-4.23	-0.84	-4.35
NiC ₃ C	-3.52	0.21	-3.31	1.81	-1.76
NiN ₄	-3.00	0.77	-2.24	2.18	-0.82
NiNC ₃	-3.25	-0.38	-3.62	1.04	-2.20
NiC ₄	-1.69	-0.45	-2.14	-0.12	-1.81

where n is the number of CO molecules coordinated to the Ni adatom, $E_{\text{sur-Ni}}$ is the total energy of the Ni-CO complexes on support, $E_{\text{sur-Ni}}$ is the total energy of the Ni adatom on support, and E_{CO} is the total energy of CO in gas phase. Furthermore we calculate Gibbs free energy for the CO-decorated nanoparticle, ($\Delta G_{\text{NP}}^{\text{dis}}$), following the equation:

$$\Delta G_{\text{NP}}^{\text{dis}}(\text{R}, \text{T}, \text{p}) = E_{\text{surf-nCO}}^{\text{f}} - n\text{CO} - n\Delta\mu_{\text{CO}}(\text{T}, \text{p}) - \Delta' E_{\text{NP}}(\text{R}) - \text{TS} \quad (5.8)$$

where, $E_{\text{surf-CO}}^{\text{f}} = E_{\text{nCO}}^{\text{f}} + E_{\text{Ni}}^{\text{f}}$ is the formation energy of the Ni-CO complexes on support with respect to the infinite size metal particle and reactants in gas phase, and $\Delta\mu_{\text{CO}}(\text{T}, \text{p})$ is the excess chemical potential of CO (T, p). In our case, the configurational entropy S was not included because the number of cavities is difficult to assess; however, it will favor dispersed species. Therefore, the values obtained here for the transition correspond to a maximum threshold. The excess chemical potential was set to $\Delta\mu_{\text{CO}}(\text{T}, \text{p}) = -0.76 \text{ eV}$ ¹⁷⁸ corresponding to experimental conditions of 300 K and 0.10 mbar. The binding energy of nCO ($E_{\text{nCO}}^{\text{f}}$) needs to be negative and lower than $\Delta\mu_{\text{CO}}(\text{T}, \text{p}) = -0.76 \text{ eV}$. When both requirements are fulfilled, nanoparticle disintegration is thermodynamically possible by realising Ni(CO)₂ complexes.

5.2 Types of N-doped carbon materials

We modelled saturated and unsaturated N-doped carbon materials under reaction conditions. The DFT simulations were performed on a number of N-doped defects where the Ni atoms can be anchored, as well as Ni nanoparticles can be supported, to represent the species characterized in experiments.¹³⁷ In the Figure 5.1 we show three families of models: (a) saturated N-doped carbon materials (NiN_xC) representing NiN₄ and similar undefined N-doped materials made at low doping temperature, (b) unsaturated N-doped (NiNC_x) and N-free materials (NiC_x) typically occurring for carbons at high doping temperatures, when the nitrogen content decreases, and (c) Ni metal surfaces composing the supported nanoparticles on the unsaturated N-doped supports (NiNP-NC_x). This approximation allows us to provide a holistic chain of models that accounts for the complexity of the N-doped carbon support and the metallic nature of the nanoparticle. The ultimate aim is to show the role of the speciation (charge, oxidation state, and coordination environment) on the electrochemical activity and on the dynamic behaviour or the different synthetic preparations.

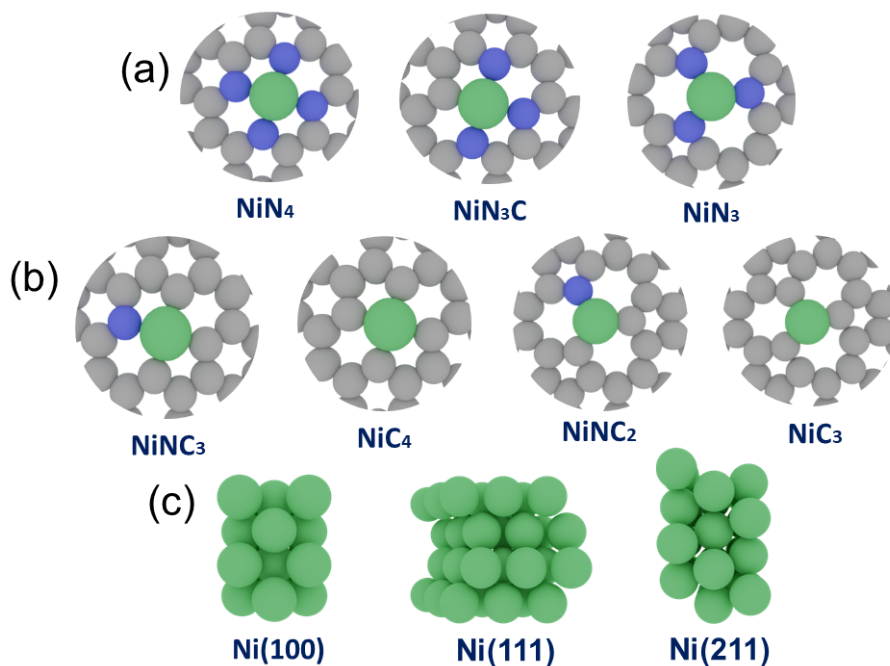


Figure 5.1: Computational models for different systems. (a) Saturated N-doped carbon material. (b) Unsaturated N-doped or N-free carbon material. (c) Metal surfaces. C in grey spheres, N in blue, and Ni in green.

5.3 Activity and selectivity of the materials

The activity and selectivity of the modelled catalysts can be discussed in terms of the corresponding Gibbs free energy profiles presented in Figure 5.2. The electrochemical activity was analyzed by calculating the binding energies of the intermediates according to the most agreed mechanism for CO₂ to CO reduction (see Chapter 2). Notice that once CO₂ is adsorbed, activated, and polarized with partial charge transfer (see Table 5.2) from the host material and thus step (i) does not correspond to any extra electron added to the simulation box. In agreement with the previous results, we can see that the highest activity towards CO is shown by unsaturated N-doped materials¹²⁵ and metal surfaces, as they bind CO strongly. Regarding the selectivity, it can be discussed by comparison to the competing HER (Figure 5.2d–f). Saturated N-doped materials reach the highest selectivity towards CO, where NiN₄ and NiN₃ are representative models, while NiN₃C model binds *H strongly. Therefore, NiN₃C model is ruled out, because it does not account for the experimental selectivity.¹³⁷

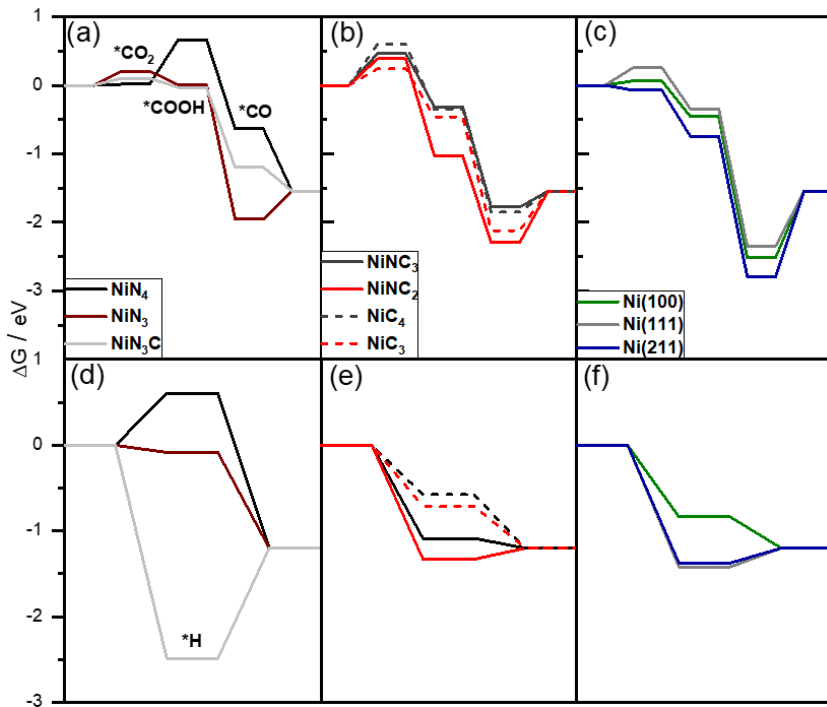


Figure 5.2: Gibbs free energy profiles for (a, b, and c) eCO₂RR and (d, e, and f) HER over different model sites at -0.6 V vs. RHE (pH 7). (a,d) Saturated nitrogen-doped carbon materials. (b,e) Unsaturated or free N-doped carbon materials. (c,f) Metal surfaces.

For a catalyst to be selective towards CO, the binding of *COOH and *CO₂ intermediates needs to be strong to reduce the onset potential. Also, selectivity can be compared with *H binding energy and CO desorption energy. The correlation between adsorption energies of *COOH and *H is shown in Figure 5.3. We can see that NiN₄ and NiN₃ bind *H weakly, therefore saturated N-doped materials are selective. NiN₄ binds *CO₂ and *COOH weakly, which means it has a higher onset potential, while the opposite trend is seen for NiN₃. The computed onset potentials are presented in Table 5.3, where NiN₃ has a lower onset potential than NiN₄ because it binds *COOH stronger. Ni(211) and NiNC₂ bind strongly *H and *COOH, which means they will not be selective toward CO. Ni(100) would be the best candidate among the metal surfaces as it binds *H weakly and *COOH relatively strongly.

Table 5.2: Bader charge of adsorbed *CO₂⁻ molecule on all computational model catalysts.

	$q(\text{CO}_2)(e^-)$
NiN₄	-0.05
NiN₃C	-0.08
NiN₃	-0.11
NiNC₃	-0.34
NiC₄	-0.09
NiNC₂	-0.08
NiC₃	-0.27
Ni(100)	-0.59
Ni(111)	-0.55
Ni(211)	-0.47

Figure 5.3b represents a correlation between adsorption energies of *CO₂ and desorption of CO, where CO₂ activation can relate to Bader charges of the adsorbed CO₂ molecule (see Table 5.2). NiN₃, NiNC₃, and Ni(100) bind *CO₂ relatively strong, where Ni(100) binds *CO strongly, which means the rate of CO formation would be low. The strong adsorption energy of *H for Ni(100) and NiNC₃ can be related to low selectivity, where HER and eCO₂RR compete. The most interesting systems is NiN₃, because of the optimal *CO binding and the weakly *H binding. From an electronic structure point of view, the Ni atom in the NiN₄ square planar confinement is in the Ni²⁺ oxidation state, while that in the NiN₃ trigonal planar configuration is in the Ni¹⁺ oxidation state (see Chapter 3).

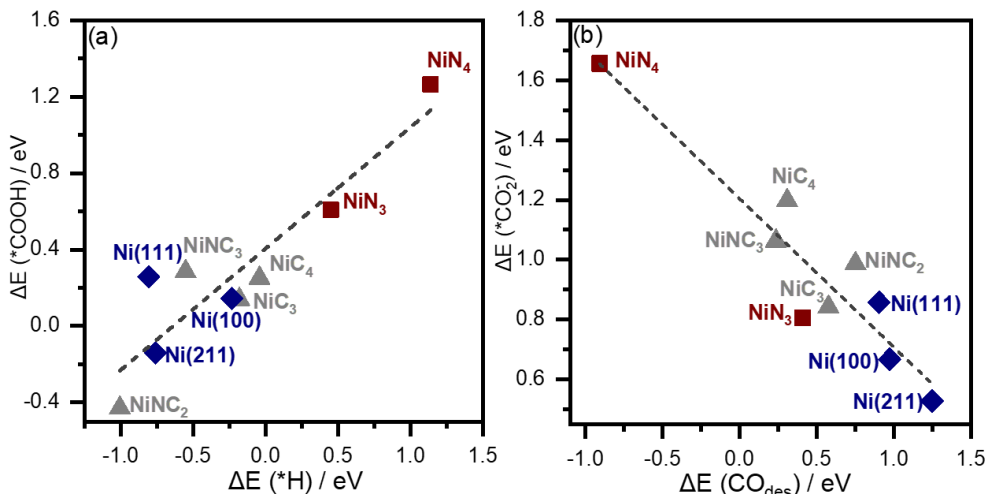


Figure 5.3: Linear scaling relationship. (a) Linear COOH binding energy scaling relations for adsorption of H atom. (b) CO₂ binding energy and corresponding linear-scaling relationships for desorption of CO molecule. Saturated N-doped carbon materials are represented with red symbols, unsaturated N-doped carbon materials with grey symbols, and metal surfaces with blue symbols.

This special role of Ni¹⁺ is long known in organometallic chemistry¹⁷⁹ and even in thermal catalysis.¹⁸⁰ The change in the oxidation state has an impact on the geometries of the activated CO₂ intermediate as both the O-C-O angle and Ni-C(CO₂) distance indicate the degree of activation of the CO₂ molecule in the N-doped saturated materials. In Figure 5.4, the linear scaling relationship of CO₂ geometry and adsorption of *CO shows that the activation in NiN₃ correlates with medium binding of *CO, while the lowest activation appears in NiN₄ with a O-C-O angle of 170.2°. All these fingerprints point out to the relevance of understanding the organometallic structures when trying to identify the most suitable defects for a given metal and the importance of the charge state for the encapsulated metal ions.

Table 5.3: DFT calculated onset potentials (in eV vs. RHE) for eCO₂RR to CO on N-doped carbon materials and Ni surfaces.

	U_{onset}
NiN ₄	-1.27
NiN ₃ C	-0.57
NiN ₃	-0.61
NiNC ₃	-0.29
NiC ₄	-0.31
NiNC ₂	-0.38
NiC ₃	-0.29
Ni(100)	-0.49
Ni(111)	-0.45
Ni(211)	-0.63

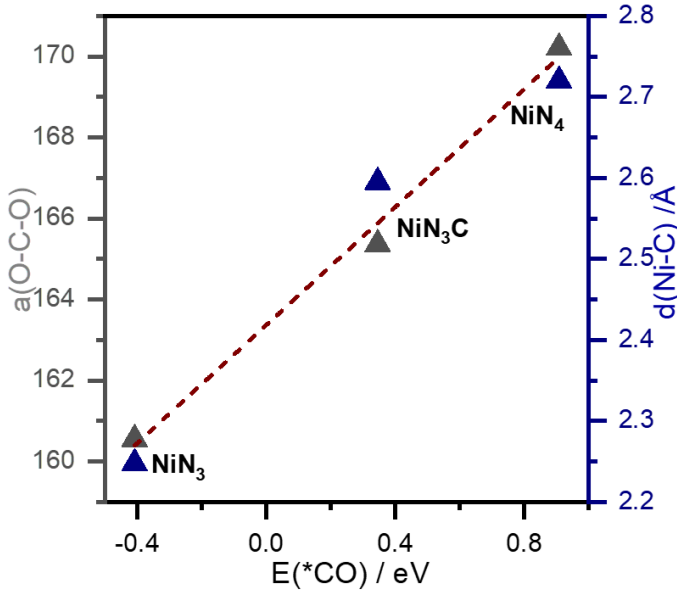


Figure 5.4: Linear scaling relationship. Linear angle $\alpha(\text{O-C-O})$ and distances $d(\text{Ni-C})$ of CO₂ scaling relations for adsorption of CO molecule.

5.4 Reconstruction of Ni NPs on N-doped carbon materials

Electrocatalysts can be affected by the presence of different variables. In our case, because of disintegration of Ni(CO)₂ complexes later the reconstruction of Ni NPs influences the activity and selectivity. Once the reconstruction happens experimentally, the partial current density of H₂ drastically drops compared to that of CO, thus favoring the disintegration process. Such dynamic behavior of electrocatalysts is crucial for their performance and thus we here address these aspects by simplified models. To investigate the role of the reaction products in the stability of the Ni NPs, we model a nanoparticle based on the lowest-energy surfaces through a Wulff construction. This is straightforward for metal nanoparticles and can be extended when adsorbates, like CO, are present.^{175,181} In addition, we apply the atomistic version of Ostwald ripening theory^{177,182} by the formation of Ni(CO)₂-like complexes to show when the disintegration of nanoparticles can occur on such N-doped carbon materials.

5.4.1 Nanoparticle shape

To address the role of nanoparticles we first focus on their structure. The size of the synthesized Ni NPs was 2.5 to 7 nm, and thus they can be represented via the Wulff construction that corresponds to the equilibrium structure for a medium-to-large nanoparticle and it is obtained from the surface energies. We calculated the surface energies for three orientations Ni(111), Ni(100), and Ni(211), and, employing the Wulff construction, we then extracted the corresponding areas for each facet (see Figure 5.5). The average surface energy for the nanoparticle is 0.126 eV/Å. In the Table 5.4 surface energies and ratios of corresponding facet are presented. To assess the change in the nanoparticle induced by the environment, we consider CO adsorption on all three facets separately.¹⁷⁵ We calculated average binding energies E'_{CO} , for a range of coverages 1/12 to 1 ML for Ni(100) and Ni(211) facets, and Ni(111) from 1/16 to 1 ML (Figure 5.6). In all cases, at low to medium coverages, all the modified surface energies (γ'_x) are smaller, but the reduction in surface energies depends on the particular facet orientation. The detailed balance between adsorption sites and CO-CO repulsion contributes to the surface energy in a non-linear way, as to achieve higher CO densities the molecules need to go to less favourable sites. At 0.25 ML we are left with Ni(100) and Ni(111) surfaces, while at full coverage of 1 ML the extent of Ni(211) facet in the nanoparticle increases and Ni(111) reconstructs (see Appendix Chapter B).

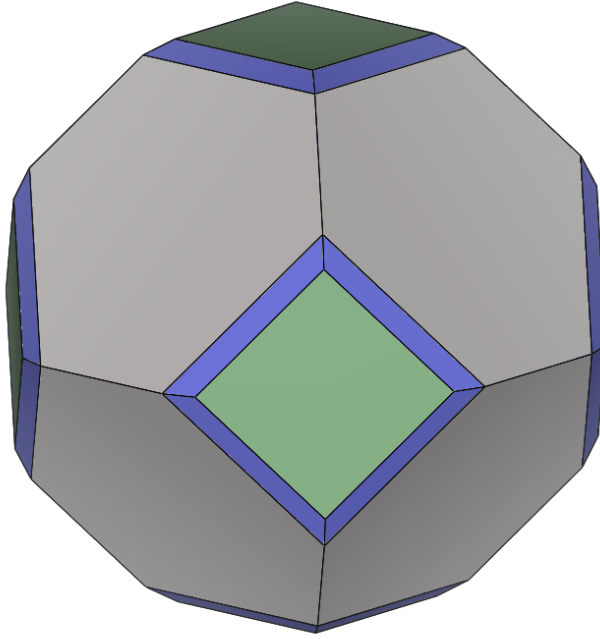


Figure 5.5: Wulff construction based on the surface energies calculated by DFT. Surface legend: Ni(100) in green, Ni(111) in gray, and Ni(211) in blue.

In summary, the average surface energy for the decorated nanoparticle (γ') at CO coverage 0.25 ML is $0.095 \text{ eV}/\text{\AA}^2$, which is $0.031 \text{ eV}/\text{\AA}^2$ less than for the bare one (see Appendix Chapter B). Therefore, the nanoparticle under product-rich conditions has more surface Ni atoms with less Ni-Ni average coordination.

Table 5.4: Calculated surface energies γ_i (in $\text{eV}/\text{\AA}^2$) of Ni nanoparticle and ratios f_i (in %) of considered facets of infinite size Ni, based on Wulff construction.

Facet	γ_i	f_i
Ni(100)	0.139	18
Ni(111)	0.120	70
Ni(211)	0.138	12
γ' (T,P)	0.126	

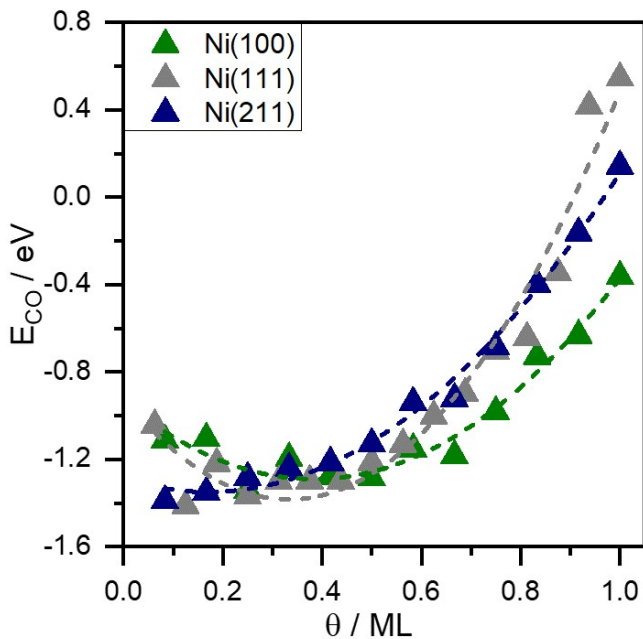


Figure 5.6: The calculated average binding energy for a CO molecule on the different Ni surfaces versus CO coverage. The DFT results are polynomial fitted.

In addition, this high product coverage influences the selectivity as it limits the HER reaction. The effect can be explained with the Gibbs free energies of the $^*\text{H}$ intermediate for the clean and on the CO covered ($\theta = 0.92\text{ML}$) Ni surfaces (see Figure 5.7). Once Ni surfaces are covered with CO, $^*\text{H}$ binds weaker (lowering the selectivity for HER), see Figure 5.8. This agrees with the experimental observations showing a drastic drop of H_2 partial current density. Thus, high CO coverage minimizes the HER path and increases selectivity.

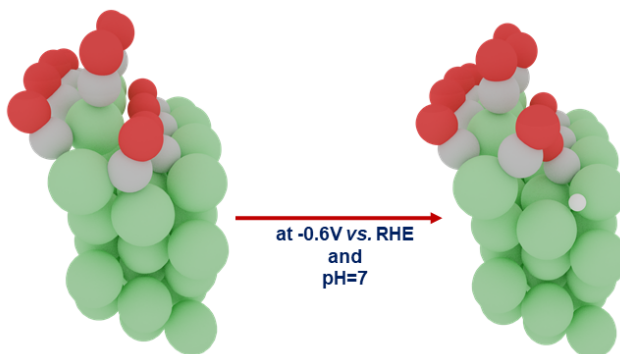


Figure 5.7: Schematic representation of $^*\text{H}$ intermediate on Ni(211) surface at 0.92 ML CO coverage.

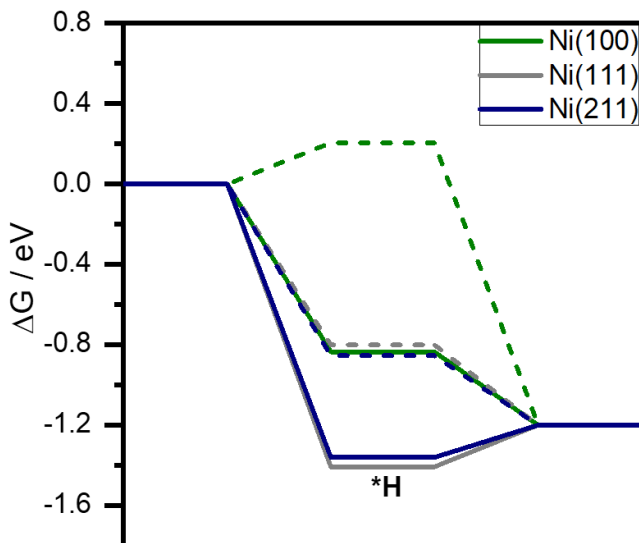


Figure 5.8: Gibbs free energy profiles for HER over clean metal (solid lines) and 0.92 ML CO coverage (dashed lines) sites at -0.6 V vs. RHE (pH 7).

5.4.2 Gibbs free energy of Ni(CO)₂ complexes disintegration on N doped carbon materials

We applied the energetic formalism (see the section Methodology) from Ouyang and co-workers¹⁷⁷ to evaluate whether the disintegration of Ni(CO)₂ from Ni NPs is possible on N-doped carbon materials. The nature of the cavities in the support provides a large versatility regarding the stability against disintegration. In our investigations, we have considered all 7 types of defects in Figure 5.1. The anchored Ni atom in saturated N-doped models is coordinated either to pyridinic or pyrrolic N atoms, which are strong ligands able to easily trap metal atoms. This is in line with experiments where they do not observe nanoparticles on those materials. We observe an opposite trend for the unsaturated N-doped materials, where Ni NPs can redisperse into single atom cavities. The reason for that is the decreased N-content and increased C-content in the materials, which makes Ni dynamic and prone to aggregate and disintegrate in the presence of dense CO phases and electrochemical condition. Therefore, single Ni metal atoms with little aggregation are expected in agreement with experiments.¹²⁸ Finally, we show that the redispersion of Ni NPs on NC₂ defect (unsaturated N-doped model) is thermodynamically possible. NiNC₂ is also one of the most active and non-selective eCO₂RR configurations, with strong bindings for *CO, *COOH, and *H intermediates. Once Ni NPs are supported on NC₂ they can disintegrate, as the experimentally partial current density drops drastically compared to CO partial current densities. That is how the combination of the two results (activity and stability for dispersion) can support the fact that enhanced Faradaic efficiency was observed when the reaction was performed

five times in a row for nanoparticles supported on unsaturated N-doped carbon materials.¹³⁷

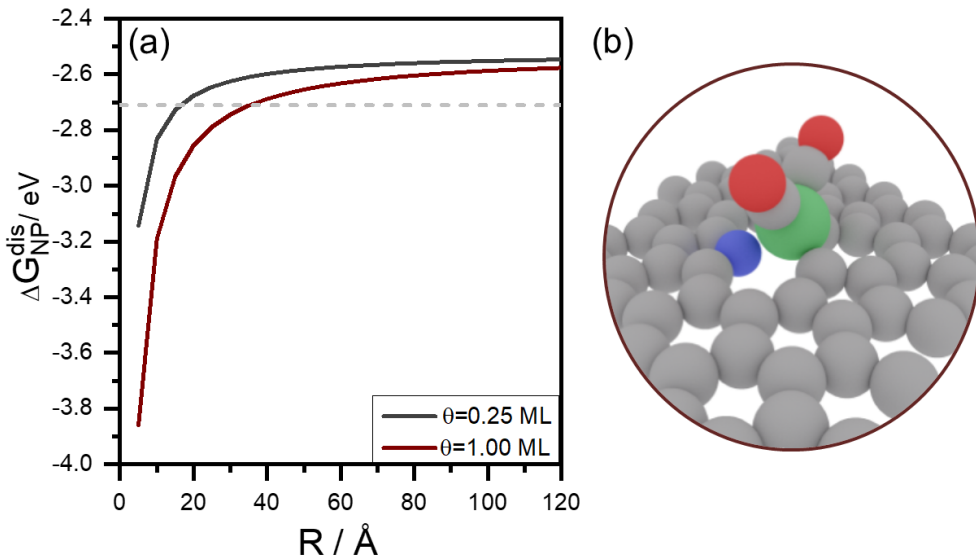


Figure 5.9: Gibbs free energy of Ni(CO)₂ complexes disintegration. (a) Size dependence of the Gibbs free energy ($\Delta G_{\text{NP}}^{\text{dis}}$) of Ni(CO)₂, when Ni NPs is supported on NC₂ at 300 K and 0.1 mbar. The horizontal dashed line indicates the limit for stability. (b) Schematic representation of Ni(CO)₂ bound to the NC₂ cavity.

Redispersion is thus possible on unsaturated materials; however, the nature of the carrier would affect catalytic performance. The electrochemical cycling from the nanoparticles does enhance the selectivity in unsaturated hosts but not up to the levels found for the saturated materials. To investigate the size (R) dependence of the disintegration induced by CO, we have used NiNP-NC₂ model as cavity and have calculated the $\Delta G_{\text{NP}}^{\text{dis}}(R, T, p)$ considering the formation of Ni(CO)₂ sitting at these pockets at experimental temperature and pressure. Figure 5.9 presents the $\Delta G_{\text{NP}}^{\text{dis}}$ corresponding to Ni(CO)₂ vs. the radius R of a Ni nanoparticle at two different coverages, 0.25 ML and 1ML (non-wetting conditions for the nanoparticle were considered, thus the Young angle was set to 180°). We found that $\Delta G_{\text{NP}}^{\text{dis}}$ decreases at smaller R and crosses the reference level (-2.7 eV, adsorption energy of an isolated Ni atom in the cavity) at about 40 Å for 1 ML and 20 Å for 0.25 ML. The lower value found for lower coverage is a direct consequence of the simplicity of the model and shall be taken with caution as kinetic constraints related to the need of concentrating CO molecules to extract the surface atoms could play an important role. At high coverage our results imply that Ni NPs smaller than 4nm will disintegrate into Ni(CO)₂ when supported on the NC₂ at 300 K and 0.1 mbar.

5.4.3 The mechanism for Ni NPs reconstruction

The dense layers that we have identified in section Nanoparticle shape show that after eCO₂RR, CO extensively covers the surface and over long times has the potential to increase the fraction of Ni atoms in low-coordinated sites as stepped Ni(211) surfaces. Moreover, the thermodynamic models (see Figure 5.1) indicate that below a certain size, Ni(CO)_x species can be stable in some particular motifs of the N-doped carbons. To merge these two scenarios, a viable route for decomposition needs to be described. Ni can form volatile species with CO in the form of Ni(CO)₄ that can be soluble under the reaction conditions. We have calculated the energy needed for the formation of these species. On nanoparticles, CO₂ is reduced to CO at -0.6 V vs. RHE and pH 7, and the resulting CO drives the reconstruction toward Ni(211). Once the reaction is stopped at open circuit potential (OCP), Ni(CO)₄ species can be formed as the CO coverage is very high. Thermodynamically, the formation energy of a metal vacancy (with respect to Ni bulk) in Ni(100), Ni(111) and Ni(211) surfaces are 0.59, 1.04 and 0.33 eV, respectively. Strong kinetic limitations can likely appear due to geometric constraints on Ni(111) and Ni(100), but for the fully CO adsorbed Ni(211) surface the step edges can present very high CO concentration. Thus Ni(CO)₄ can be formed in this facet with concomitant elimination of Ni atoms. Figure 5.10a presents the average adsorption energy for CO at different coverages (red curve) on the Ni(211) facet and it is compared to the formation-elimination of the Ni(CO)₄ complex (gray curve).

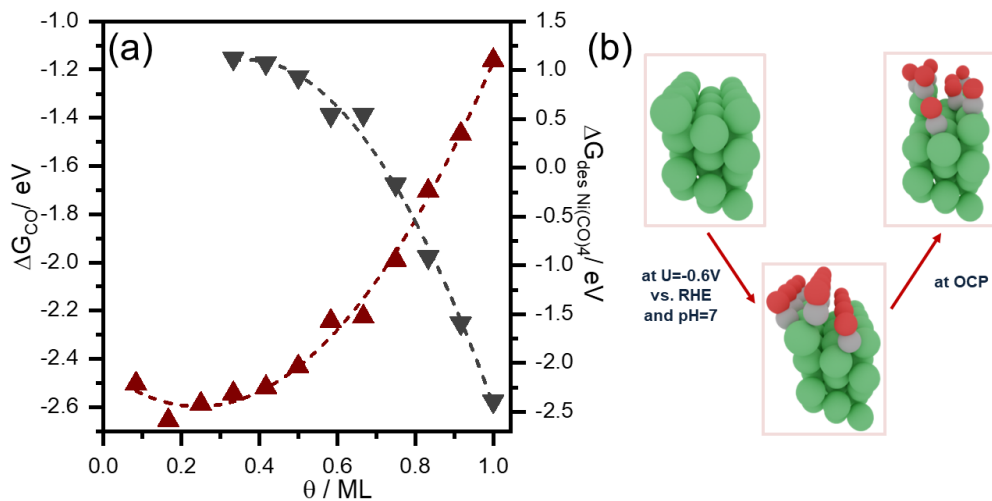


Figure 5.10: Reconstruction of Ni(211) metal surface.(a) CO coverage dependence of the average Gibbs free energy (red curve) of CO adsorption at -0.6 V vs. RHE (pH 7) and (gray curve) desorption of Ni(CO)₄. (b) Schematic process of clean Ni(211) metal surface fully adsorbing CO at electrochemical conditions. When the reaction is stopped at the open circuit potential (OCP) to assess the material recyclability, Ni(CO)₄-like complexes can desorb.

The crossing between these two lines can be interpreted as the coverage at which the elimination of Ni atoms in the form of the species becomes viable. Therefore, high CO coverages would affect the stability of the edge sites starting around 0.78 ML coverage promoting the formation of the species that can then redisperse on the cavities of the N-doped carbon. The process of reconstruction is sketched in Figure 5.10b. All the data is presented in Appendix Chapter B. Once the reaction is under electrochemical potential again, these complexes react with non-occupied defects on the unsaturated N-doped carbon (NiNP-NC₂) and effectively redisperse. This process could potentially lead to higher activity, as shown in the experiments.¹³⁷

5.5 Conclusions

Electrocatalysts based on Ni supported on N-doped carbon materials have been computationally analyzed for the reduction of CO₂ considering both nanoparticles and single atoms. The optimal activity and selectivity are found for NiN₃ model, in which Ni is in Ni¹⁺ oxidation state. NiN₃ exhibits theoretically low onset potential in agreement with experiments and high rate towards CO. However, rather saturated N-doped materials are more selective towards CO formation. As single atoms can be prepared from nanoparticles, we show through the Oswald formalism that disintegration of Ni(CO)_x can happen if Ni nanoparticles are supported on N-doped unsaturated support such as NC₂. We present a mechanism of Ni(211) surface reconstruction that is possible only at high CO coverages around 8/12 ML by the formation of Ni(CO)₄ species, which then redisperse into active single atoms. The complete understanding of the Ni-N-C materials points towards a very dynamic behavior where considering the full complexity of the material in terms of Ni nuclearity and environment is needed. This is likely a common feature for such type of electrocatalysts.

Chapter 6

Temperature dependent distribution of CO₂ electrochemical reduction on CoTPP/MWCNT catalyst

Among the various classes of catalytic materials, molecular catalysts that consist of ligands bound to metal ions have chemically modifiable structures that allow their activity to be tuned quite precisely. In this chapter DFT was used to understand experimental findings on the selectivity of molecular catalyst, Cobalt(II)-tetraphenyl porphine (CoTPP)/multiwalled carbon nanotube (MWCNT) towards different eCO₂RR products. Cobalt(II)-phthalocyanine (CoPc) deposited on carbon nanotube (CNT) and FeN₄C models were used to benchmark the obtained results. Calculations for HER were performed for all three materials while intermediates toward CH₂OH, and CH₄ were performed for CoTPP/MWCNT, toward CH₃OH for CoPc/CNT and toward CH₄ for the FeN₄C.

6.1 Methodology

6.1.1 Computational details

To model CoTPP/MWCNT and CoPc/CNT catalysts, a graphene layer expanded in a (9×9) supercell was used and embedded complexes. These layers were interleaved by 18 Å vacuum along the Z-direction. In order to model atomically-dispersed FeN₄C carbon structures, graphene layer expanded in a (6×6) supercell replacing four C atoms by nitrogens and removing two C atoms. These layers were interleaved by 12 Å vacuum along the Z direction. The kinetic energy cutoff was set at 500 eV for all the models. We investigated competitive eCO₂RR and HER, where structures were relaxed with a force threshold of 0.05 eV/Å. The Brillouin zone was sampled using a 2×2×1 k-point mesh for CoTPP/MWCNT and CoPc/CNT, and 3×3×1 k-point mesh for FeN₄C generated with the Monkhorst–Pack method.⁵⁵

6.2 Characterization of CoTPP/MWCNT catalyst

The CoTPP/MWCNT catalyst was characterized by high-resolution transmission electron microscopy (HRTEM) and scanning electron microscopy (SEM). Both methods suggest a strong interaction existing between MWCNT and CoTPP as intended. Raman spectroscopy was employed for further characterization of the CoTPP/MWCNT catalyst. The Raman spectra of MWCNT exhibit two characteristic peaks the D-band (ca. 1314 cm⁻¹) is assigned to the sp³ hybridized carbon or the disordered nanotube structure^{183,184} and the G-band (ca. 1575 cm⁻¹) correspond to the sp² hybridized carbon atoms which are generated by tangential vibrational mode of the nanotubes. CoTPP peaks in a range of 300-1600 cm⁻¹ are visible and originate from the symmetric and asymmetric stretching and bending molecular vibrational modes of chemical species.¹⁸⁵ The existence of such interactions can also be experimentally proven with 1H NMR.¹⁸⁶ In that study a Ru complex, with aromatic ligands, was deposited on MWCNT, where the experimental and DFT results show existence of CH- π interactions, therefore we suggest a similar behavior for our system. The formation of the catalyst is observed from the change in the intensity ratio (I_D/I_G), which are 0.40 and 0.46 for pristine MWCNT and the catalyst, respectively. The intensity ratio of the D band to the G band (I_D/I_G) corresponds to the amount of sp³-hybridized carbon atoms on the sp² conjugated carbon materials.¹⁸⁷

Co in CoTPP is in a square planar configuration, while phenyl groups are in perpendicular configuration, where hydrogen atoms of the phenyl rings interact with the π system of the MWCNT forming multiple CH- π interactions^{186,188,189} (see Figure 6.1a). Each of the four aromatic rings form two CH- π interactions, where the H...C distances go from 2.65Å to 3.33Å. In order to benchmark experimental findings over CoTPP/MWCNT catalyst, CoPc/CNT¹⁹⁰(see Figure 6.1b) and FeN₄C (see Figure 4.1) are modelled. CoPc complex is in a square planar configuration; where phenyl rings interact

with the π system of the MWCNT forming multiple π - π interactions. Interaction between both π systems are around 3.5\AA , adsorption energies of clean CoTPP and CoPc complexes on the graphene layer are -1.94 eV and -2.28 eV , respectively.

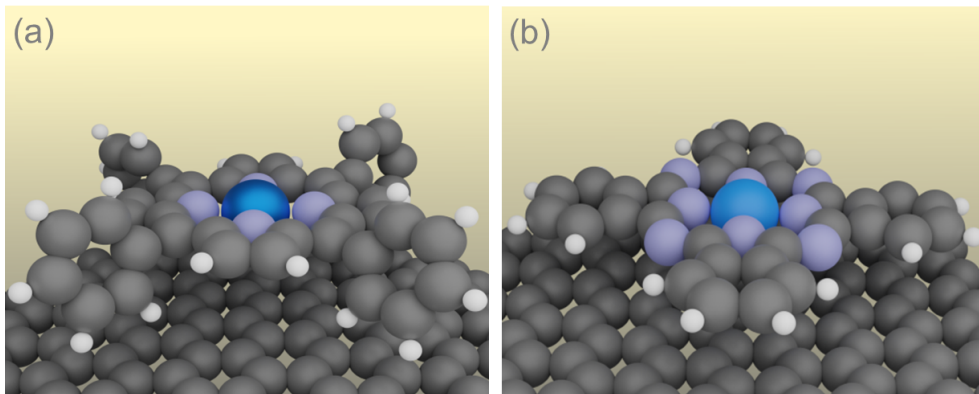


Figure 6.1: Computational models. Structures of the (a) CoTPP/MWCNT catalyst and (b) CoPc/CNT, where C is in grey, N in purple, Co in blue, and H in white.

6.3 Catalytic performance of CoTPP/MWCNT catalyst

From the experimental results we can conclude that the CoTPP is active toward CO₂ reduction reaction when deposited on MWCNT. From Cyclic voltammograms is evident that water reduction is occurring at elevated temperature and increased potential. As a result of simultaneous electrocatalytic reduction of CO₂, H₂O, and subsequent reduction of CO the gas chromatography (GC) detected H₂, CO and CH₄ as gaseous products. At the end of the controlled potential electrolysis, catholyte was collected and analyzed with high-performance liquid chromatography (HPLC). In these measurements, HCOO⁻ and CH₃OH were detected as liquid products. HCOO⁻ is generated with a low Faradaic efficiency, and thus low partial current density, therefore formate is excluded from the discussion. Furthermore, CV experiments were performed at room temperature (20°C), 30°C, 40°C and 50°C in N₂ and CO₂ saturated electrolytes to elucidate the direct proton and H₂O reduction mechanism. At 40°C and 50°C H₂O reduction becomes predominant.¹⁹¹

Regarding the selectivity in terms of Faradaic efficiencies, we noticed two trends in the product distribution influenced by the temperature over CoTPP/MWCNT catalyst. With increasing temperature, the formation of H₂ is increasing, while for the CH₃OH, CO, and CH₄ the opposite trend is observed. An exception is observed at -0.6 V vs. RHE , where FE of CO is increasing with the temperature until 40°C and at 50°C when it decreases. Figure 6.2 shows the FE of detected products, while partial current densities are shown in Figure 6.3. FE and partial current density of H₂ both increase

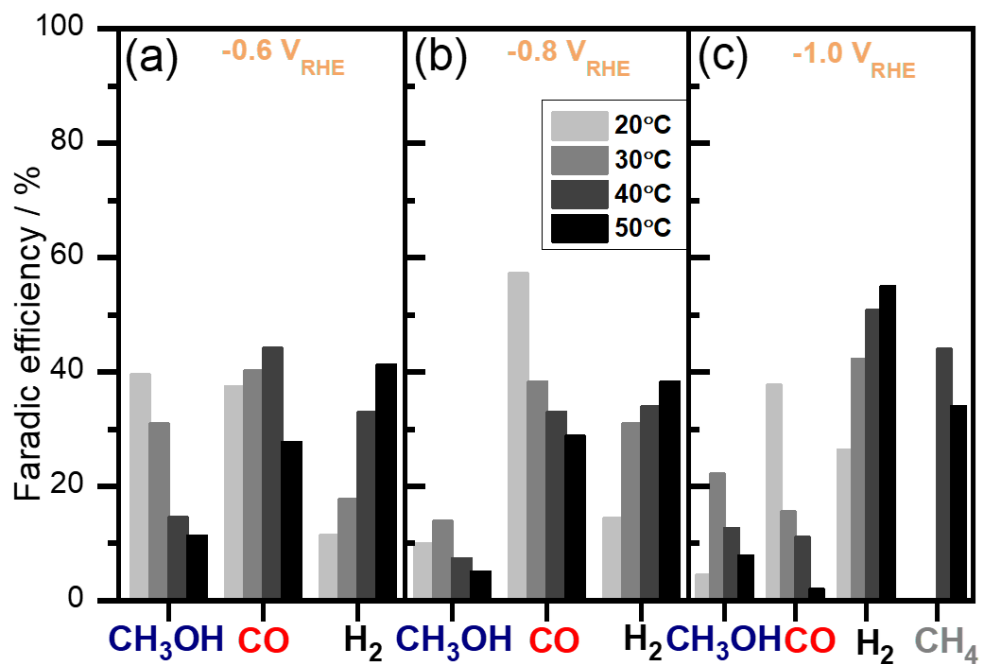


Figure 6.2: Product analysis. Faradaic efficiency of the detected products at studied temperatures at (a) -0.6 V, (b) -0.8 V, and (c) -1.0 V vs. RHE over the CoTPP/MWCNT catalyst, plotted against the obtained products. The data were obtained by CA for 4 hours, performed in 0.1 M KHCO₃ aqueous electrolyte saturated with CO₂ (pH 6.8).

with temperature and potential, with FE being close to 60% at 50°C and -1.0 V vs. RHE. One of the reasons could be solubility which is a function of the temperature, where CO₂ with increasing temperature is getting less soluble in water, while the solubility of H₂ is constant and low.¹⁹²

Consequently, the H₂ formation is increasing because of the low CO₂ solubility in water. A large contribution to the current generation comes from the direct reduction of H⁺ and/or H₂O to H₂. The highest CO formation was observed at room temperature and -0.8 V vs. RHE with FE of 57%. Syngas was produced at 30°C and higher temperatures and at -0.8 V vs. RHE, where the ratios of CO and H₂ are close to 1:1. The temperature influence could be a way to control ratios for syn-gas formation. Significantly in the design of metgas plants, ratios of CO and H₂ are accommodated to be 1:2 to drive CH₃OH synthesis.^{7,193} At room temperature and -0.6 V vs. RHE the main product is CH₃OH, while at the same potential production of CH₃OH is decreasing with increasing temperature. eCO₂RR proceeds in such a way that adsorbed CO₂ reduces via parallel pathways to liquid and gaseous products CH₃OH and CO. Therefore, a considerable amount of CO is desorbed from the electrocatalyst surface before further reduction and it is suggested to be an intermediate species during the formation of the liquid products. An interesting trend occurs with decreasing the potential, at -0.6 V the FE of CH₃OH and CO is relatively the same, 40% and 38% respectively. Once the potential decreases to -0.8 V vs. RHE, the FE for CH₃OH reduces to 10%, while for CO enlarges to 57%. The formation of CH₄ was observed only at -1.0V vs. RHE and elevated temperatures (40°C and 50°C), with FE being more than 35%.

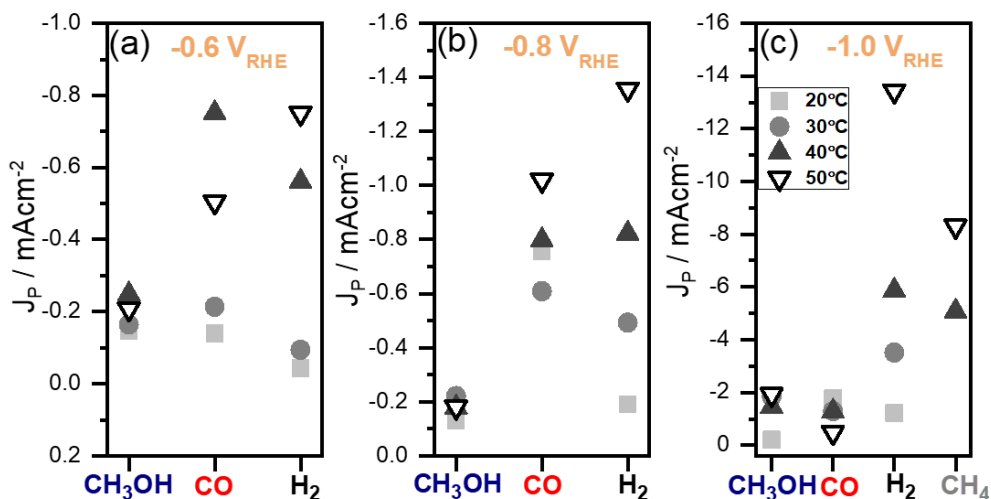


Figure 6.3: Partial current analysis. Partial current densities of the detected products at studied temperatures at (a) -0.6 V, (b) -0.8 V, and (c) -1.0 V vs. RHE over the CoTPP/MWCNT catalyst, plotted against the obtained products. The data were obtained by CA for 4 hours, performed in 0.1 M KHCO₃ aqueous electrolyte saturated with CO₂ (pH 6.8).

6.4 CO₂ reduction pathway

Upon adsorption and reduction of CO₂, the first main intermediate is *CO (Figure 6.4a). Releasing CO would regenerate the catalyst; otherwise, *CO can further hydrogenate to *CHO. This is a key intermediate from which different products can be obtained. On the one hand, hydrogenation of the C atom yields formaldehyde *CH₂O through a O–M bond, which can either desorb or keep reacting to eventually form CH₃OH (Figure 6.4b). On the other hand, hydrogenation of the O atom produces *CHOH via a C–M bond, which can evolve to CH₃OH (two proton and electron transfers) or CH₄ and H₂O (four proton and electron transfers) depending on the hydrogenation pathway (Figure 6.4c). Alternatively, the initial *CO could also hydrogenate to *COH (Figure 6.4d). Additional hydrogenation with concomitant water elimination could produce *C.¹⁹⁴ This intermediate can accept up to four protons and electrons to finally yield CH₄. However, this mechanism is not chemically possible for single atom metals as carbon needs to form four bonds.

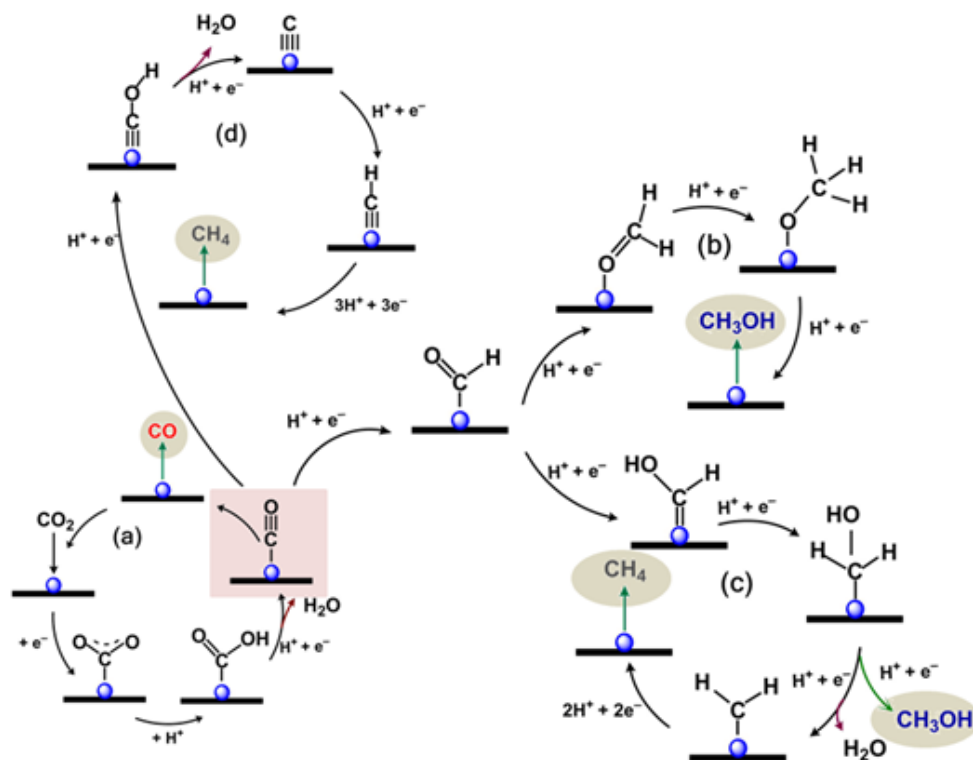


Figure 6.4: Electrochemical CO₂ reduction network. Reaction mechanism towards (a) CO, (b) CH₃OH, (c) CH₄ and side product CH₃OH, and (d) CH₄ on CoTPP/MWCNT.

6.5 Understanding product selectivity

There are ways to understand product selectivity based on the DFT results. One way would be to determine the oxophilicity of the catalyst, which correlates with the binding energy of O_{ads}, and is the main descriptor for selectivity between CH₄ and CH₃OH.¹⁹⁵ In that work, the authors studied selectivity over transition metals. CH₃OH was detected on the poorly oxophilic Au metal surface, while CH₄ was found for the strongly oxophilic Fe. To benchmark this approach we applied this theory to our case, CoTPP/MWCNT and a reference FeN₄ catalyst, see Figure 6.5. The bonding energy of *O to CoTPP/MWCNT is -2.65 eV, and to the reference FeN₄ model is 1.18 eV more exothermic. Indeed, on FeN₄ active sites the formation of CH₄ was detected¹⁹⁶ thus reinforcing the validity of the approach. A second descriptor for selectivity corresponds to energy of oxygen bonding *OCH₂, where the most exothermic energy has CoTPP/MWCNT. All three materials have exothermic *H binding energy with CoTPP/MWCNT having the lowest -1.07 eV at -0.6 V vs. RHE. From all of this findings we can conclude that CoTPP/MWCNT and CoPc/CNT are able to produce CH₃OH, while FeN₄ forms CH₄.

Table 6.1: Gibbs free energies of *OCH₂, *OCH₃, *CHOH, and *CH₂OH at -0.6 V vs. RHE, where energies were obtained with the respect to reactants (CO₂ and support(*)). OCH₂ binding energy was calculated respect to OCH₂ molecule and *(colored row)

	CoTPP/MWCNT	CoPc/CNT	FeN ₄
*OCH ₂	-4.26	-2.26	-2.98
*OCH ₃	-4.52	-3.36	-4.55
*CHOH	-3.09	-1.47	-3.87
*CH ₂ OH	-5.41	-4.31	-4.68
*OCH ₂	-1.21	-0.32	-0.37

Furthermore to investigate the product formation pathway, electrolysis of HCOOH acid and CH₂O was performed. Interestingly, when HCOOH was reduced only H₂ formation is observed, while CH₂O is producing both CH₃OH and CH₄. From CH₂O reduction, we get *OCH₃, where oxygen bonding gives alcohols while *CH₂OH the carbon bonding leads to fully reduced hydrocarbons.¹⁹⁷ The difference in these systems can be noticed between oxygen bonding *OCH₂ and carbon bonding *CHOH intermediates. For CoTPP/MWCNT and CoPc/CNT the *OCH₂ is more exothermic than *CHOH, while for FeN₄ the opposite trend is observed (see Table 6.1). Furthermore the electrochemical activity and selectivity were analyzed by calculating the binding energies of the possible intermediates. The Gibbs free energy network for the different

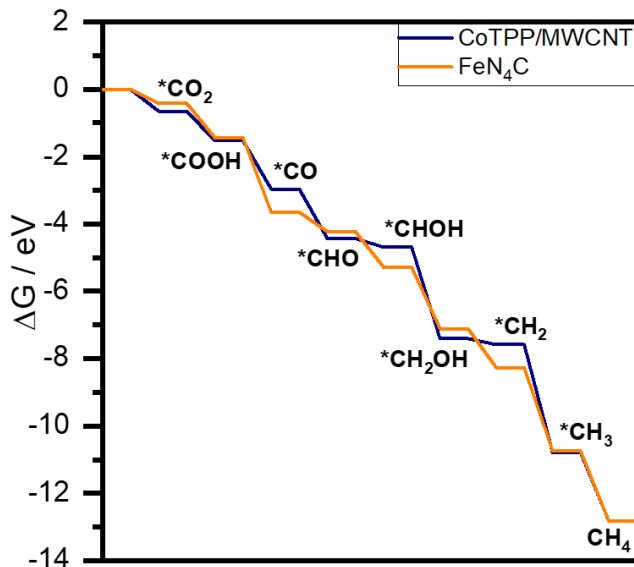


Figure 6.5: Gibbs free energy profiles for eCO₂RR over CoTPP/MWCNT and FeN₄ towards CH₄ at -1.0 V vs. RHE (pH 6.8).

product pathways is presented in Figure 6.6. Compared to previous results in the literature, at room temperature different Co complexes predominately produce CO, except for the CoPc/CNT¹⁹⁰ catalyst which produces CH₃OH. When the CoTPP was used but supported in a different material, CNT instead of MWCNT, CoTPP/CNT¹⁹⁸ 90% of CO production at -0.68 V vs. RHE were reported.

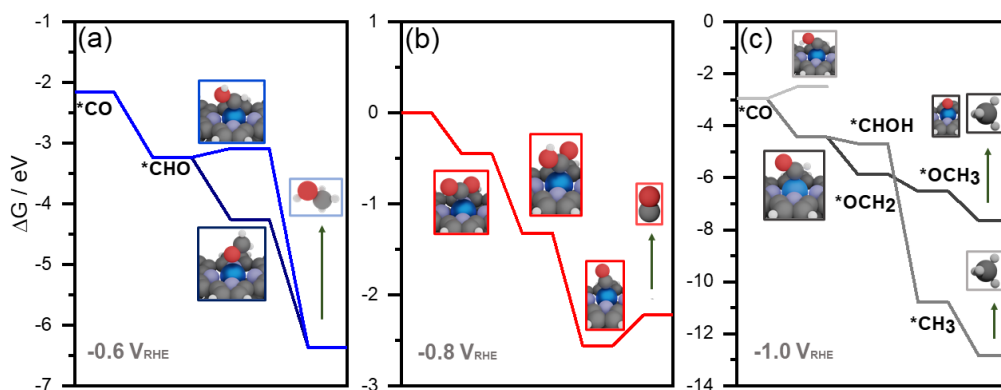


Figure 6.6: Gibbs free energy profiles. Energy diagrams for eCO₂RR over CoTPP/MWCNT. (a) Towards CH₃OH at -0.6 V vs. RHE, (b) CO and at -0.8 V vs. RHE, and (c) CH₄ at -1.0 V vs. RHE (pH 6.8).

In contrast, in this study at room temperature (20°C) and -0.6 V vs. RHE we observe the production of CH₃OH. From simulations, *CO is determining intermediate to produce higher value products and in Figure 6.6b the pathway towards CO is shown at -0.8 V. In the Figure 6.5 we can see that CoTPP/MWCNT binds CO weaker than FeN₄C does.

At the lower potential region, desorption of CO is accelerated by higher temperature as indicated by the increasing total yield of CO (total current times FE) with the temperature.^{185,199} The protonation of *CO to *CHO is exothermic for CoTPP/MWCNT catalyst, which makes CO easier to protonate towards CH₃OH. As mentioned at room temperature with reducing potential, FE and J of CO increases at -0.8 V comparing to -0.6 V vs. RHE. According to the DFT at -0.8 V vs. RHE, more exothermic, which means it would make CO harder to desorb (Figure 6.7). The reason for a difference in results could be a structural change of an active site during the reaction, porphyrin complexes have a reductive decomposition of the macrocycle.^{200,201} The influence of structural features of catalytic sites, meaning the number of coordinated N atoms, and configuration of phenyl groups, plays an important role. The results for CO₂ reduction obtained over CoPc/CNT catalyst¹⁹⁰ show that at -0.82 V around 80 % of CO gets produced, and at -0.94 V vs. RHE around 44% of CH₃OH is formed (see Figure 6.8). Figure 6.7 presents Gibbs free energy diagram for CoTPP/MWCNT and CoPc/CNT molecular catalysts at -0.6 V vs. RHE. Comparing *CO binding energies between systems, a difference of 2.43 eV. CoTPP/MWCNT catalyst binds CO stronger than CoPc/CNT, that is the reason why at -0.6 V CH₃OH gets produced. According to binding energy of *CO to CoPc/CNT at -0.8 V potential. CO gets desorbed rather than going further to produce CH₃OH.

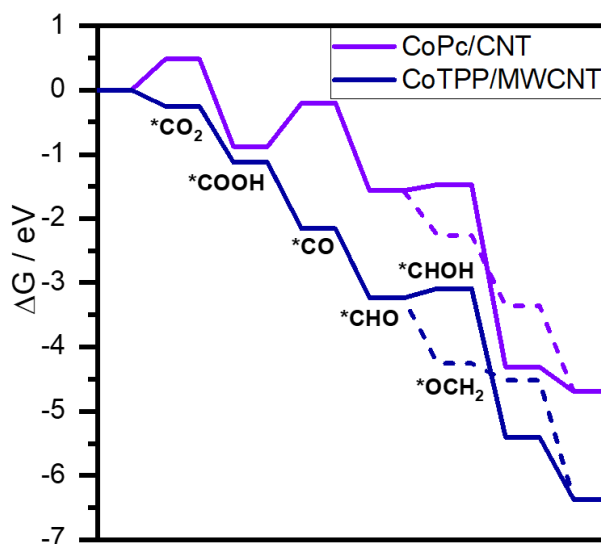


Figure 6.7: Gibbs free energy profile for eCO₂RR over CoTPP/MWCNT and CoPc/CNT towards CH₄ at -0.6 V vs. RHE (pH 6.8).

In the Figure 6.6a the pathways towards forming CH₃OH are presented for CoTPP/MWCNT. As we confirmed the oxygen bonding of CH₂O is more favourable for CoTPP/MWCNT catalyst, the reduction will continue through *OCH₃ intermediate. Based on our result the *OCH₃ and *CH₂OH are similar in energy, thereby both possibilities would be possible. In the case of CH₄ formation (Figure 6.4d), we can go through *COH or *CHO intermediate, where the *COH intermediate is higher in energy by 1.94 eV, therefore the pathway is excluded. As CH₄ is traced at elevated temperatures and low potential, there is a way it can be produced from reduction of CH₃OH or can be reduced to CH₄ through *CH₂OH and *CH₃ intermediates.

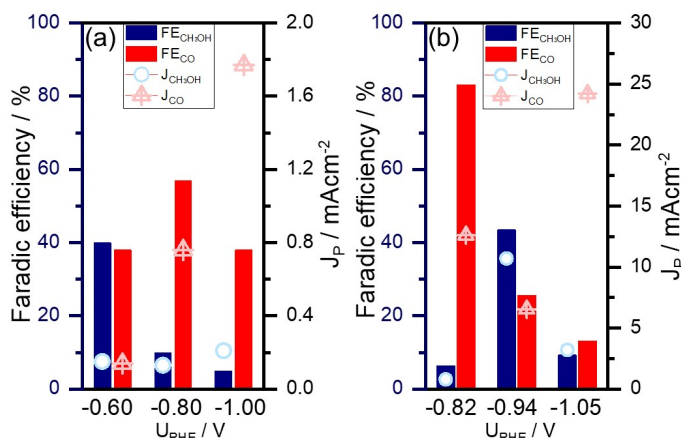


Figure 6.8: Comparison of catalytic performance. Product selectivity (FE) and partial current densities of CH₃OH and CO over (a) CoTPP/MWCNT and (b) CoPc/CNT (Ref.¹⁹⁰) at different potentials.

Increasing temperature to 30°C seems to favor the production of CO and CH₃OH, because the mass transport of CO₂ can be improved by increasing its solubility by adjusting the pressure and the temperature.²⁰² Mass transfer can be enhanced with the modulation of the H₂O surface tension that occurs at elevated temperatures. The modulation of surface tension may help to increase the wettability of the electrode surface, which may contribute to decrease of the diffusion layer thickness and hence the mass transfer.²⁰³

6.6 Conclusions

In this chapter, study supported with both experimental and theoretical results is presented which helped to unravel two interesting features (i) the improved activity for the liquid product formation compared to earlier published studies, and (ii) the temperature dependent products distribution favouring hydrogen production at elevated temperatures. It has been shown that improved activity of the hybrid catalyst is due to the interaction between the CoTPP catalytic active centers and electron conducting MWCNTs. Furthermore comparing the structure of the CoTPP and CoPc complex plays an important role in activity of the catalyst. Temperature affects both the gaseous and liquid product formation, making differences in the balance of energy of the different routes in the mechanism that impact product distribution. Notably, FE of CH₃OH decreases with temperature and in contrary FE of CO increases which indicates that the temperature enhances the desorption of CO before further reduction. CoTPP/MWCNT is poorly oxophilic comparing to FeN₄, indeed it has an affinity to produce CH₃OH. Also, the efficiency of the competing HER increases with the increasing temperature and overpotential. It is suggested from CVs that the H₂ is produced from both direct proton and H₂O reduction. The outstanding methanol production observed at room temperature and low potential points toward a interface design to improve product selectivity.

Chapter 7

General conclusions

Electrocatalysis plays an important role in clean energy conversion, because it would allow sustainable production of a broad range of chemicals and fuels. Electrochemical CO₂ reduction reaction over single site catalysts was investigated using DFT simulations coupled with Computational Hydrogen Electrode formalism. In Chapter 3 I described the importance of understanding the fundamental mechanisms of single-atom catalysis to design systems with improved performance and stability. In addition the crystal field theory within coordination chemistry is explained for anchored homogeneous-like MN_x centers. In Chapter 4 I establish experimental and theoretical correlations between physiochemical and catalytic properties for the CO₂ reduction towards CO for MNC catalyst. In Chapter 5 I explained the process of Ni nanoparticles reconstruction by the disintegration of Ni(CO)₂ complex, by applying Oswald formalism, on N-doped carbon materials. Finally, in Chapter 6 I unravelled the CO₂ reduction product influenced by changing potential and the temperature over modelled CoTPP/MWCNT catalyst. The conclusions specified by Chapter are summarized below.

With respect to the volcano trend in electrocatalytic CO₂ reduction activity over atomically dispersed metal sites on Nitrogen-doped Carbon:

- Among the best performing materials, Fe²⁺N₄-H₂O and Co²⁺N₄-H₂O catalyst were theoretically found as the most likely active catalytic centers at studied potentials, whereas Ni¹⁺N₄ was predicted as the most active Ni-based one.
- The Gibbs free energy change at the RDS accounted for the experimentally determined activity volcano, with Co-, Fe-, and Ni-based systems showing the best compromise between *CO and *CO₂⁻ binding energies.
- Regarding selectivity, the difference between the binding energies for *CO₂⁻ and *H, directly related to the number of d-antibonding electrons, described successfully the high selectivity observed on Fe and Ni catalysts and the more favoured HER observed over the Co based one.

With respect to the stability and redispersion of Ni nanoparticles supported on N-doped carbons for the electrochemical CO₂ reduction

- The optimal activity and selectivity are found for NiN₃ model, in which Ni is in Ni¹⁺ oxidation state.
- NiN₃ exhibits theoretically low onset potential in agreement with CO high rates.
- Saturated N-doped materials are selective towards CO formation.
- Redispersion of Ni nanoparticles happens by disintegration of Ni(CO)₂-like complex, when Ni nanoparticles are supported on N-doped unsaturated support such as NC₂.
- A mechanism of Ni(211) surface reconstruction was developed. Reconstruction is possible only at high CO coverages around 8/12 ML by the formation of Ni(CO)₄ species.
- At 0.92ML CO coverage ($\Theta=0.92\text{ML}$) Ni surfaces bind *H weaker, which means the selectivity for HER is decreasing.

With respect to the temperature dependent selectivity of CO₂ electrochemical reduction on CoTPP/MWCNT catalyst

- The improved activity of CH₃OH at lower overpotential (-0.6 V vs. RHE) is achieved.
- Increasing temperature to 30°C favours production of CO and CH₃OH, while H₂ formation gets produce at elevated temperatures.
- Employing temperature could be a way of controlling the ratio of CO and H₂ for syngas production.
- The CoTPP/MWCNT catalyst is poorly oxophilic, which correlates with the binding energy of O_{ads}, thus making the catalyst selective towards CH₃OH.
- The *OCH₂ intermediate is for -1.17 eV more exothermic than *CH₂OH.

Appendices

UNIVERSITAT ROVIRA I VIRGILI
THEORETICAL STUDIES OF SINGLE-SITE CATALYSTS FOR EFFICIENT ELECTROCHEMICAL CO₂ REDUCTION
Paulina Pršlja

Appendix A

Structural characterization data

A.1 Structural parameters for studied models

Tables A.1-A.3 list structural parameter such as distance between metal and local environment for MN₄⁻, Ni saturated and unsaturated N-doped, CoTPP, and CoPc materials. The label R_{M-N} or R_{M-C}, stands for distance between certain metal and Nitrogen or Carbon atom.

Table A.1: Structural metal-nitrogen/oxygen distances obtained for MN₄⁻ calculated model structures (in Å).

	R _{M-N}	R _{M-O}
MnN₄-H₂O	1.91	2.40
FeN₄-H₂O	1.89	2.33
CoN₄-H₂O	1.89	2.31
ZnN₄-H₂O	1.97	2.35
FeN₄-OH	1.90	1.81
MnN₄-O	1.95	1.58

Table A.2: Structural Ni-nitrogen/carbon distances obtained for Ni saturated and unsaturated N-doped calculated model structures (in Å).

	$R_{\text{Ni-N/C}}$
NiN₄	1.88
NiN₃	1.80
NiN₃C	1.89
NiC₄	1.88
NiNC₃	1.89
NiC₃	1.74
NiNC₂	1.77

Table A.3: Structural Co-nitrogen distances obtained for CoTPP and CoPc calculated model structures (in Å).

	$R_{\text{Co-N}}$
CoTPP	1.97
CoPc	1.91

A.2 Oxidation state of embedded metals

A.2.1 Magnetization and Bader charge

Tables A.4-A.6 list magnetization and Bader charges of the embedded metal in studied materials such as MN₄⁻, Ni saturated and unsaturated N-doped, CoTPP/MWCNT, and CoPc/CNT.

Table A.4: Magnetisation and Bader charges for MN₄-like catalysts.

		*	*CO ₂ ⁻	*COOH	*CO	*H ⁺
MnN₄	μ_n	2.89	2.89	2.08	1.14	2.12
	q(Mn)(e ⁻)	1.36	1.36	1.35	1.35	1.33
FeN₄	μ_n	1.82	0.00	0.81	0.00	0.92
	q(Fe)(e ⁻)	1.13	1.11	1.09	1.06	1.07
CoN₄	μ_n	0.82	0.57	0.00	0.00	0.00
	q(Co)(e ⁻)	0.94	0.95	0.94	0.96	0.93
Ni(I)N₄	μ_n	0.78	0.00	0.38	0.89	0.37
	q(Ni)(e ⁻)	0.75	0.82	0.93	0.92	0.89
NiN₄	μ_n	0.00	0.00	0.38	0.00	0.37
	q(Ni)(e ⁻)	0.86	0.82	0.93	0.88	0.89
CuN₄	μ_n	0.54	0.44	0.00	0.54	0.52
	q(Cu)(e ⁻)	0.95	0.93	0.90	0.96	1.05
ZnN₄	μ_n	0.00	0.00	0.00	0.00	0.00
	q(Zn)(e ⁻)	1.14	1.08	1.17	1.16	0.97
MnN₄-H₂O	μ_n	2.96	1.58	2.04	1.12	2.11
	q(Mn)(e ⁻)	1.42	1.33	1.39	1.32	1.41
FeN₄-H₂O	μ_n	0.00	0.00	0.01	0.00	0.90
	q(Fe)(e ⁻)	1.19	1.11	1.19	1.06	1.14
CoN₄-H₂O	μ_n	0.88	0.61	0.00	0.40	0.00
	q(Co)(e ⁻)	1.05	0.96	1.02	1.03	0.99
MnN₄-O	μ_n	0.46	0.45	2.07	2.03	0.02
	q(Mn)(e ⁻)	1.57	1.54	1.46	1.64	1.49
FeN₄-OH	μ_n	1.04	0.29	0.86	0.00	0.92
	q(Fe)(e ⁻)	1.28	1.26	1.30	1.24	1.22

Table A.5: Magnetisation and Bader charges for MN₄-like catalysts.

		*	*CO ₂ ⁻	*COOH	*CO	*H ⁺
NiN ₄	μ_n	0.00	0.00	0.00	0.96	0.34
	$q(\text{Ni})(e^-)$	0.86	0.85	0.93	0.92	0.90
NiN ₃ C	μ_n	0.02	0.01	0.00	0.10	0.12
	$q(\text{Ni})(e^-)$	0.75	0.72	0.86	0.82	0.87
NiN ₃	μ_n	1.27	1.23	0.74	1.10	0.73
	$q(\text{Ni})(e^-)$	0.84	0.93	1.05	0.72	0.88
NiC ₄	μ_n	0.00	0.00	0.00	0.00	0.10
	$q(\text{Ni})(e^-)$	0.64	0.61	0.65	0.64	0.61
NiNC ₃	μ_n	0.36	0.19	0.00	0.00	0.00
	$q(\text{Ni})(e^-)$	0.67	0.68	0.72	0.70	0.65
NiC ₃	μ_n	0.00	0.00	0.00	0.00	0.10
	$q(\text{Ni})(e^-)$	0.58	0.63	0.49	0.48	0.45
NiNC ₂	μ_n	1.26	0.00	0.00	0.00	0.00
	$q(\text{Ni})(e^-)$	0.72	0.56	0.72	0.73	0.53

Table A.6: Magnetisation and Bader charges for CoTPP/MWCNT catalysts.

	$\mu_{\mathbf{n}}$	$q(\text{Co})(e^-)$
*	0.98	1.12
* H	0.00	1.01
* CO₂	0.84	1.09
* COOH	0.00	1.02
* CO	0.70	1.10
* CHO	0.00	0.99
* OCH₂	0.97	1.14
* OCH₃	1.34	1.23
* CHOH	0.26	1.07
* H₂CHO	0.00	1.03
* COH	0.00	0.98
* C	0.79	0.99
* CH	0.39	1.06
* CH₂	0.00	1.14
* CH₃	0.00	1.06
* O	1.75	1.36
* OH	1.41	1.28

A.2.2 Determining oxidation state

The computed magnetization of embedded metal and geometry of chosen defect becomes useful to determine oxidation state of the metal using crystal field theory. Figures A.1-A.3 present the crystal field energy splitting for the studied materials.

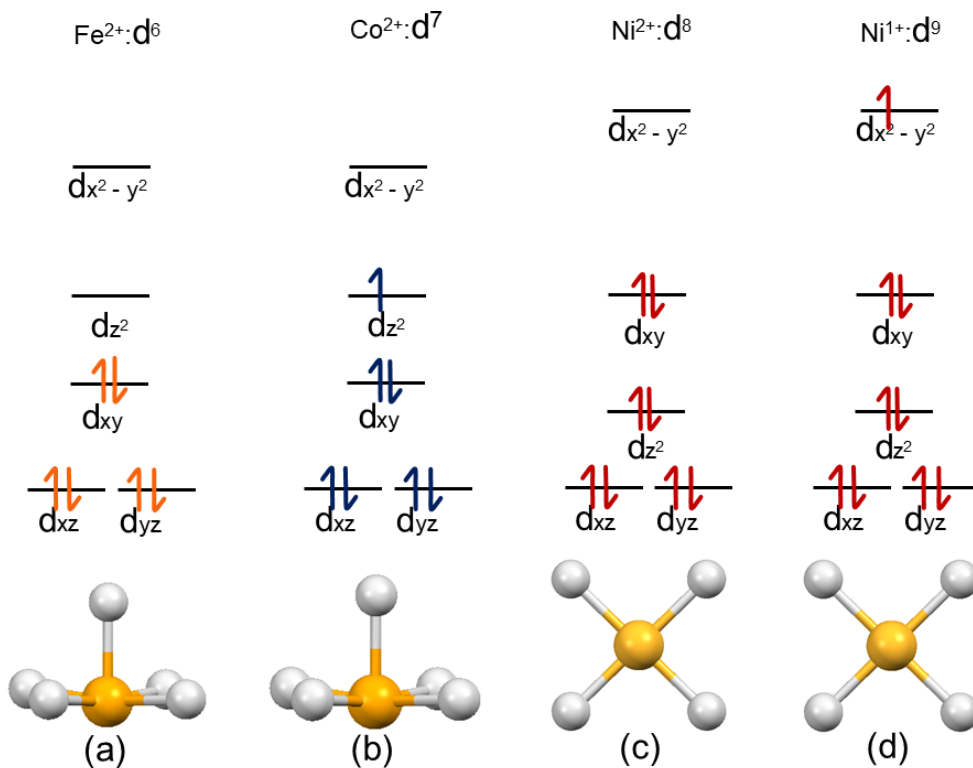


Figure A.1: The crystal field theory splitting of the metal d block for (a) $\text{FeN}_4\text{-H}_2\text{O}$ (square-based pyramidal) (b) $\text{CoN}_4\text{-H}_2\text{O}$ (square-based pyramidal) (c) Ni(II)N_4 (square planar) and (d) Ni(I)N_4 (square planar). Ligands are represented with white spheres, and metal with yellow sphere.

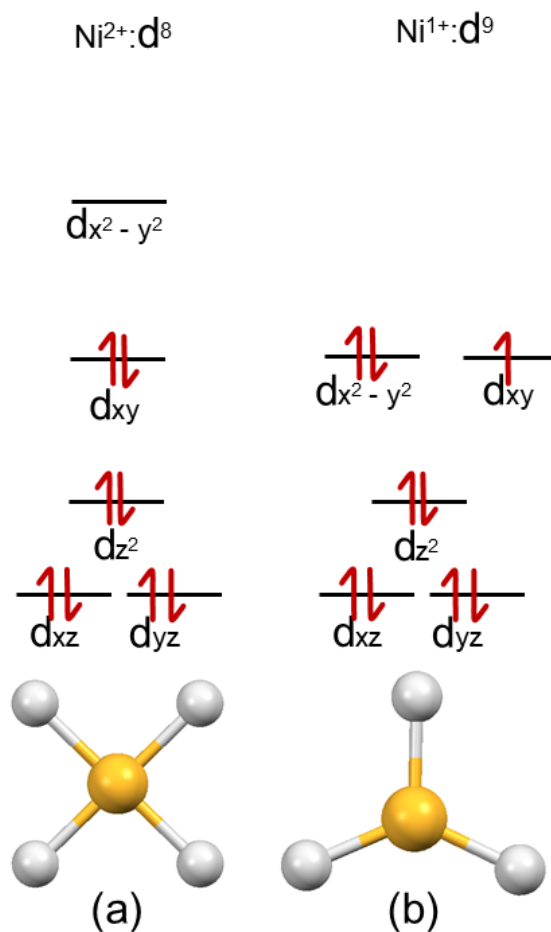


Figure A.2: The crystal field theory splitting of the metal d block for (a) Ni(II)N₄ (square planar) and (b) Ni(I)N₄ (trigonal planar). Ligands are represented with white spheres, and metal with yellow sphere.

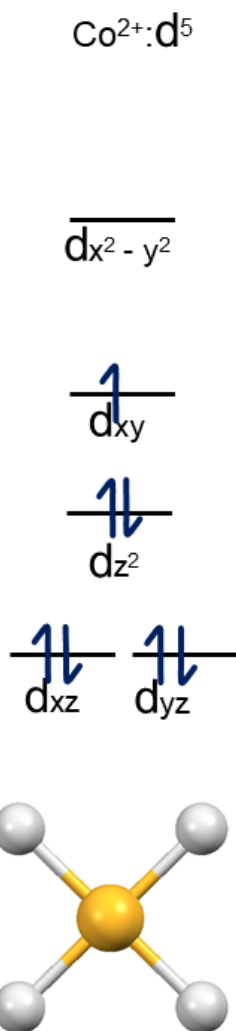


Figure A.3: The crystal field theory splitting of the metal d block for CoTPP and CoPc complexes (square planar). Ligands are represented with white spheres, and metal with yellow sphere.

Appendix B

Additional methodological details

B.1 Construction of Ni nanoparticle at different CO coverages

In order to calculate surface energy of Ni NPs at certain CO coverage, than CO adsorption should be considered on all three facets separately. Therefore I calculated average binding energies E'_{CO} for all three facets at 0.25 ML, 0.50 ML, 0.75 ML, and 1 ML CO coverage (see Table B.1). Table B.2 lists the reduction of surface energy $\Delta\gamma_i$ and effective surface energy γ'_i of supported Ni particles with CO adsorbate. Once the the reduction of surface energies is calculated, than the surface energy of Ni NPs at certain CO coverage can be calculated by

$$\gamma'(T, p) = \sum_i f_i [\gamma_i + \Delta\gamma_i(T, p)]. \quad (\text{B.1})$$

Table B.3 lists the surface energy and ratios at certain CO coverage of the facetes and surface energy of the Ni NP. Figure B.1 presents Wulff construction of Ni NP at CO coverages 0.25 ML, 0.50 ML, 0.75 ML, and 1.00 ML.

Table B.1: Average adsorption energies of CO (in eV/CO molecule) for different coverages and surface unit area of the facet "i" (in Å²)

	Ni(100)	Ni(111))	Ni(211)
A	74.44	85.95	60.84
A_i	6.20	5.37	5.07
E_{CO-0.25 ML}	-1.34	-1.37	-1.28
E_{CO-0.50 ML}	-1.29	-1.21	-1.13
E_{CO-0.75 ML}	-0.98	-0.71	-0.69
E_{CO-1.00 ML}	-0.36	0.55	0.14

Table B.2: Calculated surface energies of Ni particle before γ_i and after CO adsorption γ'_i and ratio f_i , f'_i of considered facets of infinite size Ni based on Wulff construction. Ratios are presented in (%) and γ_i in (eV/Å²)

θ	0.25 ML	0.50 ML	0.75 ML	1.00 ML
$\Delta\gamma_{100}(\mathbf{T},\mathbf{p})$	-0.024	-0.042	-0.027	0.064
$\Delta\gamma_{111}(\mathbf{T},\mathbf{p})$	-0.028	-0.042	0.007	0.243
$\Delta\gamma_{211}(\mathbf{T},\mathbf{p})$	-0.024	-0.036	0.011	0.178
$\gamma'_{100}(\mathbf{T},\mathbf{p})$	0.016	0.015	0.068	0.201
$\gamma'_{111}(\mathbf{T},\mathbf{p})$	0.078	0.067	0.050	0.000
$\gamma'_{211}(\mathbf{T},\mathbf{p})$	0.000	0.000	0.000	0.004

Table B.3: Calculated surface energies of Ni facets after CO adsorption γ'_i and ratio f'_i of considered facets of infinite size Ni based on Wulff construction. Ratios are presented in (%) and γ_i in (eV/Å²)

θ	0.25 ML		0.5 ML		0.75 ML		1.00 ML	
	γ_i	f_i	γ_i	f_i	γ_i	f_i	γ_i	f_i
Ni(100)	0.116	14	0.102	15	0.149	61	0.203	99
Ni(111)	0.092	86	0.078	85	0.127	39	0.363	0
Ni(211)	0.114	0	0.102	0	0.149	0	0.316	1
γ' (T,P)	0.095		0.081		0.118		0.204	

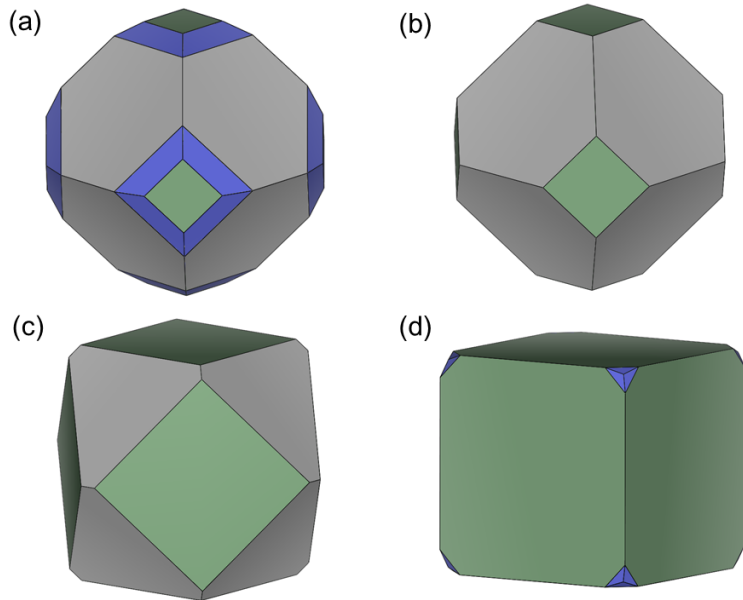


Figure B.1: Wulff construction based on the surface energies calculated by first principle theory. (a) Ni nanoparticle with 0.25 ML coverage of CO. (b) Ni nanoparticle with 0.5 ML coverage of CO. (c) Ni nanoparticle with 0.75 ML coverage of CO. (d) Ni nanoparticle with 1.0 ML coverage of CO. Ni(100) surface is represented in green color, Ni(111) surface with gray color, and Ni(211) surface with blue color.

B.2 Reconstruction of Ni(211)

Tables B.4 and B.5 present Gibbs free energies of CO molecule adsorption and desorption, while desorption energies of Ni(CO)₄ compound are presented in Table B.6. CO adsorption follows as $* + n\text{CO} + n2\text{H}^+ + n2\text{e}^- \rightarrow *\text{CO} + n\text{H}_2\text{O}$ and is calculated with applied CHE formalism as:

$$\Delta G_{\text{DFT}} = nG(*\text{CO}) + nG(\text{H}_2\text{O}) - nG(*) - nG(\text{CO}_2) - nG(\text{H}_2) \quad (\text{B.2})$$

Table B.4: Gibbs free energies of CO molecule adsorption depending on the coverage of Ni(211) metal surface at -0.6 V vs. RHE and pH 7. All energies in eV.

	θ (ML)	$\Delta G_{\text{a-nCO}}$	$\Delta G_{\text{a-nCOave}}$
$\Delta G_{\text{a-1CO}}$	0.08	-2.50	-2.50
$\Delta G_{\text{a-2CO}}$	0.17	-5.31	-2.66
$\Delta G_{\text{a-3CO}}$	0.25	-7.76	-2.59
$\Delta G_{\text{a-4CO}}$	0.33	-10.18	-2.54
$\Delta G_{\text{a-5CO}}$	0.42	-12.59	-2.52
$\Delta G_{\text{a-6CO}}$	0.50	-14.59	-2.43
$\Delta G_{\text{a-7CO}}$	0.58	-15.70	-2.24
$\Delta G_{\text{a-8CO}}$	0.67	-17.80	-2.23
$\Delta G_{\text{a-9CO}}$	0.75	-17.90	-1.99
$\Delta G_{\text{a-10CO}}$	0.83	-17.04	-1.70
$\Delta G_{\text{a-11CO}}$	0.92	-16.13	-1.47
$\Delta G_{\text{a-12CO}}$	1.00	-13.94	-1.16

Table B.5: Gibbs free energies of CO molecule desorption depending on the coverage of Ni(211) metal surface at -0.6 V vs. RHE and pH 7. All energies in eV.

	θ (ML)	ΔG_{d-nCO}	$\Delta G_{d-nCOave}$
ΔG_{d-1CO}	0.08	1.15	1.15
ΔG_{d-2CO}	0.17	2.22	1.11
ΔG_{d-3CO}	0.25	3.13	1.04
ΔG_{d-4CO}	0.33	4.00	1.00
ΔG_{d-5CO}	0.42	4.87	0.97
ΔG_{d-6CO}	0.50	5.33	0.89
ΔG_{d-7CO}	0.58	4.90	0.70
ΔG_{d-8CO}	0.67	5.45	0.68
ΔG_{d-9CO}	0.75	4.01	0.45
ΔG_{d-10CO}	0.83	1.61	0.16
ΔG_{d-11CO}	0.92	-0.84	0.08
ΔG_{d-12CO}	1.00	-4.58	-0.38

Table B.6: Gibbs free desorption energies of Ni(CO)₄ (in eV).

θ (ML)	$\Delta G_{Ni(CO)_4}$	$\Delta G_{Ni(CO)_4}^{ava}$
0.09	4.29	1.07
0.18	3.73	0.93
0.27	2.16	0.54
0.36	2.19	0.55
0.45	-0.69	-0.17
0.55	-3.66	-0.91
0.64	-6.37	-1.59
0.73	-9.57	-2.39

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Publications

UNIVERSITAT ROVIRA I VIRGILI
THEORETICAL STUDIES OF SINGLE-SITE CATALYSTS FOR EFFICIENT ELECTROCHEMICAL CO₂ REDUCTION
Paulina Pršlja

Volcano Trend in Electrocatalytic CO₂ Reduction Activity over Atomically Dispersed Metal Sites on Nitrogen-Doped Carbon

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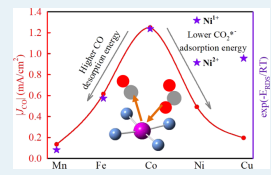
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Supporting Information

ABSTRACT: The development of catalysts for electrochemical reduction of carbon dioxide (eCO₂RR) with high activity and selectivity remains a grand challenge to render the technology useable. As promising candidates, metal–nitrogen–carbon (MNC) catalysts with metal atoms present as atomically dispersed metal–N_x moieties (MN_x, M = Mn, Fe, Co, Ni, and Cu) were investigated as model catalysts. The distinct activity for CO formation observed along the series of catalysts is attributed to the nature of the transition metal in MN_x moieties because of otherwise similar composition, structure, and morphology of the carbon matrix. We identify a volcano trend between their activity toward CO formation and the nature of the transition metal in MN_x sites, with Fe and/or Co at the top of the volcano, depending on the electrochemical potential. Regarding selectivity, FeNC, NiNC, and MnNC had Faradaic efficiency for CO >80%. To correctly model the active sites in operando conditions, experimental operando X-ray absorption near edge structure spectroscopy was performed to follow changes in the metal oxidation state with electrochemical potential. Co and Mn did not change the oxidation state with potential, whereas Fe and Ni were partially reduced and Cu largely reduced to Cu(0). Computational models then led to the identification of M²⁺N₄–H₂O as the most active centers in FeNC and CoNC, whereas Ni¹⁺N₄ was predicted as the most active one in NiNC at the considered potentials of –0.5 and –0.6 V versus the reversible hydrogen electrode. The experimental activity and selectivity could be rationalized from our density functional theory results, identifying in particular the difference between the binding energies for CO₂^{*–} and H* as a descriptor of selectivity toward CO. This in-depth understanding of the activity and selectivity based on the speciation of the metals for eCO₂RR over atomically dispersed MN_x sites provides guidelines for the rational design of MNC catalysts toward eCO₂RR for their application in high-performance devices.

KEYWORDS: CO₂ electroreduction, CO₂ activation, MNC catalyst, single-atom catalyst, metal–N_x electrocatalysis



1. INTRODUCTION

Global climate change triggered by anthropogenic emission of CO₂ is a growing threat to society. Direct electrochemical reduction of CO₂ (eCO₂RR) to produce value-added chemicals and fuels using renewable electricity is a promising strategy for closing the carbon cycle. However, this early-stage technology is challenged by the low activity and selectivity because of the extremely stable C=O bond (806 kJ mol⁻¹) in CO₂ and the competing hydrogen evolution reaction (HER) in aqueous electrolytes.^{1,2} Thus, it is of paramount importance to

identify robust catalysts with economically viable current density at low overpotential and high selectivity toward eCO₂RR.

Large ensembles of crystallographically ordered metal atoms (reduced metals, metal-oxides, chalcogenides, metal-oxide@metal core@shell, etc.) have hitherto been the most studied

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family of eCO₂RR catalytic sites. They can be categorized into four types, according to their main eCO₂RR products and/or selectivity:^{3–5} (1) formate or formic acid for catalytic sites with low binding energy for CO₂* (typical examples are ensembles based on Sn or Pb); (2) CO for catalytic sites with strong COOH* binding but weak CO* binding (typical examples are metallic surfaces such as Au, Ag); (3) products corresponding to 4–12 electron reduction (HCHO, CH₃OH, CH₄, C₂H₄, C₂H₅OH) for catalytic sites with less weak CO* binding energy (metallic Cu is an outstanding material for achieving this); (4) catalytic sites with low selectivity toward eCO₂RR, with H₂ as the main product (such as Pt or Ni metallic surfaces).

A second family of eCO₂RR catalysts involves active sites based on a single metal-atom coordinated only with nonmetal elements. In this family, the subgroup of molecular catalysts with a well-defined metal coordination have been studied since the 1970s and have shown promising properties for eCO₂RR.⁵ In particular, some transition-metal phthalocyanines and porphyrins with well-defined square-planar MN₄ structure presented high selectivity toward CO formation.^{6–9} The electroreduction of CO₂ to CO (CO₂ + 2H⁺ + 2e⁻ → CO + H₂O) with standard electrochemical potential of -0.11 V versus a reversible hydrogen electrode (RHE) is attractive as it requires the transfer of only two electrons, leading to lower activation barriers than for hydrocarbon formation.⁹ Together with formic acid, CO has thus been identified as the economically most appealing product of eCO₂RR.^{2,4} Moreover, the simultaneous production of CO and H₂ (syngas) in appropriate composition during eCO₂RR on surfaces with multiple catalytic sites may be used to prepare synthetic fuels via a downstream Fischer–Tropsch process.¹⁰ Molecular catalysts however have several disadvantages, such as cost related to their synthesis, lower operando stability because of metal leaching or chemical degradation, and poor electrical connection with the conductive support, impeding reaching high current densities in a CO₂/water electrolyzer.^{11,12}

As a second subgroup of catalysts with single metal-atom sites for eCO₂RR, metal–nitrogen–carbon (MNC) materials exclusively comprising atomically dispersed transition-metal cations covalently bonded with nitrogen atoms (MN_x moieties) hold promise to overcome the above disadvantages of molecular catalysts while retaining high activity and selectivity toward CO formation.^{13–21} The metal-ion coordination in MNC materials resembles the local patterns existing in metal–N₄ macrocycles. Such materials were initially developed for catalyzing the oxygen reduction reaction (ORR), in particular in acidic medium, with focus on the FeN_x and CoN_x as other MN_x motives have a much lower ORR activity.^{22–25} MNC materials are prepared via the self-assembly of metal, nitrogen, and carbon atoms at a high temperature. However, the elevated temperature and excess of metal often lead to the formation of clusters of metal atoms (metallic, metal-oxide, metal-carbide, etc.) along with the MN_x sites.^{26–28} This has for a long time impeded the elucidation of the MN_x site structures and their ORR catalytic activities.²⁵ In the context of eCO₂RR, the presence of metal clusters along with MN_x moieties may, depending on the nature of the metal, catalyze the HER^{10,29–31} and/or other pathways of eCO₂RR. This has prevented establishing clear experimental structure-to-property relationships between the nature of the metal in MN_x and their activity and selectivity toward eCO₂RR. In addition, the presence of even a minute amount of metallic

clusters in MNC catalysts also impedes a reliable comparison between experimental eCO₂RR activity/selectivity trends for such materials and density functional theory (DFT) results obtained on MN₄ sites.^{32–34} Validation of DFT predictions on a reasonably large series of model materials is highly important to speed up the rational identification of more promising MNC candidates for eCO₂RR,⁵ given that the library of metals demonstrated to form MN_x sites in carbon has been recently expanded also to heavier elements such as Sn,³⁵ Ru,³⁶ Pt,³⁷ Pd,³⁷ and so on. The investigation of model MNC catalysts exclusively comprising atomically dispersed MN_x sites is thus critical to understand their structure–performance relationships and to identify the proper activity/selectivity descriptor(s). Whereas numerous studies have now reported the eCO₂RR selectivity of some atomically dispersed MNC catalysts prepared by different synthetic approaches,^{13,16,20,38} a single study has, to the best of our knowledge, compared the trends in activity and selectivity for a narrow family of three atomically dispersed MNC catalysts prepared identically except for the nature of the metal element.²¹ The synthesis leaned toward silica templating of *o*-phenylenediamine and Fe, Co, or Ni salts. Whereas metal particles formed in parallel to MN_x sites during the first pyrolysis, they could be removed during subsequent NaOH etching and acid-wash steps. However, the formation of metallic particles during the first pyrolysis leads to markedly different pore size distributions, in particular for CoNC. This rendered the comparison of the materials more complex because of expectedly different utilizations of the MN_x sites for different pore size distributions.

Herein, we applied a robust synthesis approach involving a sacrificial Zn-based metal–organic framework (MOF) and a metal salt (Mn, Fe, Co, Ni, or Cu) that resulted in the complete integration of the metal atoms as atomically dispersed MN_x moieties in the N-doped carbon matrix derived from the MOF, without formation of metallic particles. With this series of five model MNC catalysts, a volcano trend in their catalytic activity toward CO₂ electroreduction to CO could be revealed for the first time, with the CoN₄ moieties exhibiting the highest intrinsic activity at -0.6 V versus RHE. Our operando spectroscopy characterization and DFT models further reveal the importance of the speciation of the metal center under reaction conditions in order to correctly interpret and predict the competition between CO₂ activation, CO and H adsorption, thus providing a more accurate understanding of the activity of eCO₂RR toward CO over MN_x motives.

2. EXPERIMENTAL METHODS

2.1. Catalyst Synthesis. All catalyst precursors were prepared via dry ball-milling of the Zn(II) zeolitic imidazolate framework (Basolite Z1200 from BASF, labeled ZIF-8), M²⁺ acetate (M = Mn, Fe, Co, Ni, and Cu), and 1,10-phenanthroline. 1,10-phenanthroline (200 mg), 800 mg of ZIF-8, and a certain amount of M²⁺ acetate corresponding to 0.5 wt % transition metal in the total precursor were weighed and poured into a ZrO₂ crucible with 100 zirconium-oxide balls of 5 mm diameter. The ZrO₂ crucible was then sealed under air and placed in a planetary ball-miller (Fritsch Pulverisette 7 Premium, Fritsch, Idar-Oberstein, Germany). The powders were milled for four cycles of 30 min at 400 rpm milling speed. The catalyst precursors resulting from the milling were pyrolyzed at 1050 °C in Ar for 1 h. The obtained catalysts were labeled MNC. A control sample was also

synthesized according to this process, except that no separate metal acetate was added (labeled (Zn)NC).

2.2. Scanning Transmission Electron Microscopy. For the scanning transmission electron microscopy (STEM) investigations, the material was dispersed in ethanol and a few drops of the suspension deposited onto a perforated carbon foil supported on a copper grid. After evaporation of the ethanol, the grid was mounted on the single tilt holder of the microscope. STEM images were recorded on an aberration-corrected HD2700CS (Hitachi) with a high-angle annular dark-field detector (HAADF). CuNC was measured on a Mo grid to avoid the Cu stray radiation from the TEM grid. STEM combined with energy-dispersive X-ray (EDX) spectroscopy was performed on a Talos F200X microscope (ThermoFisher) with a high brightness field emission gun operated at an acceleration potential of 200 kV. The EDX system of this microscope consists of four silicon drift detectors, which enables one to record EDX maps with good signal to noise ratio in a relatively short collection time (here 10–20 min).

2.3. X-ray Diffraction. X-ray diffraction (XRD) patterns were obtained using a PANalytical X'Pert Pro powder X-ray diffractometer with Cu K α radiation.

2.4. X-ray Absorption Spectroscopy. Metal K-edge X-ray absorption spectroscopy (XAS) was conducted at room temperature at SAMBA beamline (Synchrotron SOLEIL). The beamline is equipped with a sagittally focusing Si 220 monochromator and two Pd-coated mirrors that were used to remove X-rays harmonics. The catalysts were pelletized as disks of 10 mm diameter with 1 mm thickness using Teflon powder (1 μ m particle size) as a binder. The detailed analysis method for the extended X-ray absorption fine structure (EXAFS) data can be found elsewhere.^{24,25}

For operando XAS measurements, catalyst inks were prepared by mixing 10 mg of catalyst with 50 μ L of deionized water and 100 μ L of 5 wt % Nafion solution with ultrasound. A 50 μ L aliquot was then pipetted on a \sim 3 cm² circular area of a 100 μ m thick graphite foil (Goodfellow cat. C 000200/2), resulting in a catalyst loading of \sim 1 mg cm⁻². The graphite foil then served as a working electrode, and was installed in an electrochemical cell²⁴ (PECC2, from Zahner) with an Ag/AgCl (saturated KCl) reference electrode, a Pt counter electrode, and CO₂-saturated 0.1 M KHCO₃ electrolyte. Operando measurements were performed by recording the K α X-ray fluorescence of Fe with a Canberra 35-elements monolithic planar Ge pixel array detector.

2.5. X-ray Photoelectron Spectroscopy. X-ray photoelectron spectroscopy (XPS) spectra were acquired on a Kratos Axis DLD Ultra-X-ray photoelectron spectrometer using an Al K α source monochromatic operating at 150 W with no charge compensation. The base pressure was about 2×10^{-10} Torr and the operating pressure was around 2×10^{-9} Torr. Survey and high-resolution spectra were obtained at pass energies of 80 and 20 eV, respectively. Acquisition times were 2 min for survey spectra, 3 min for C 1s spectra, and 40 min for N 1s and M 2p spectra. Data analysis and quantification were performed using CasaXPS software. A linear background subtraction was used for quantification of C 1s, O 1s, and N 1s spectra, whereas a Shirley background was applied to the M 2p spectra. Sensitivity factors provided by the manufacturer were utilized to obtain atomic percentages of Fe, N, C, and O present in samples, and the elemental percentage reported was the average of five spots per sample. A 70% Gaussian/30%

Lorentzian line shape was utilized in the curve fitting. Averages from three areas per sample were presented. N 1s spectra were fitted with a set of components with a fixed energy position, in accordance with a previous study.³⁹ The fitting components are assigned to pyridinic N (398.8 eV), N in MN_x moieties (398.8–399.9 eV), pyrrolic and hydrogenated N (400.7 eV), graphitic N (401.8 and 402.7 eV), and N-oxide (404.3 and 405.6 eV).

2.6. Scanning Electron Microscopy. Scanning electron microscopy (SEM) micrographs were obtained with a Hitachi S-4800 apparatus (Hitachi, Tokyo, Japan).

2.7. N₂ Physisorption. N₂ physisorption analysis was performed at liquid nitrogen temperature (77 K) with a Micromeritics ASAP 2020 instrument. Prior to the measurements, MNC materials were degassed at 200 °C for 5 h in flowing nitrogen to remove guest molecules or moisture.

2.8. Raman Spectroscopy. Raman spectra were collected using a LabRAM ARAMIS Raman microscope with a 473 nm laser.

2.9. Linear Scan Voltammetry. Catalyst powders were deposited on glassy carbon as the working electrode. Catalyst inks were prepared by dispersing 10 mg of catalyst in a mixture of Millipore water (36.5 μ L, 18.2 M Ω cm) and ethanol (300 μ L, Sigma-Aldrich, 99.8%), into which 5 wt % Nafion solution (108.5 μ L, Sigma-Aldrich) was added as a binder. The ink was sonicated for 60 min, and an aliquot of 7 μ L was dropcast onto glassy carbon (0.196 cm², Pine instrument), to reach 800 μ g cm⁻² loading. The resulting working electrode was used in a three-electrode cell setup connected to a bipotentiostat (Biologic SP 300) and rotator (Pine Instruments). A graphite rod and Ag/AgCl (saturated KCl) were used as counter and reference electrodes, respectively. eCO₂RR activity was measured in CO₂-saturated 0.1 M KHCO₃ (pH = 6.7), and a 0.1 M N₂-saturated KH₂PO₄/K₂HPO₄ buffer solution (pH = 6.7) was used as CO₂-free blank. Linear scan voltammetry (LSV) was carried out from -0.6 to -1.4 V versus Ag/AgCl with a scan rate of 5 mV s⁻¹ at 1600 rpm. The change of potential scale versus RHE was done according to the following equation

$$E(\text{V vs RHE}) = E(\text{V vs Ag/AgCl}) + E(\text{V of Ag/AgCl vs NHE}) + 0.059 \times \text{pH} \quad (1)$$

where $E(\text{V of Ag/AgCl vs NHE}) = 0.199 \text{ V}$.

2.10. Chronoamperometry and Product Analysis. Catalysts were deposited on carbon gas diffusion layers (GDLs) by airbrushing. The ink was prepared by dispersing 50 mg of catalyst powder in a mixture of Millipore water (4 mL, 18.2 M Ω cm) and 2-propanol (4 mL, Sigma-Aldrich, 99.8%), into which 5 wt % Nafion solution (50 μ L, Sigma-Aldrich) was added as a binder. The ink was sonicated for 15 min, and then sprayed onto the GDL (Sigracet 39BC, SGL Group) with an airbrush (Iwata Eclipse HP-SBS) at 353 K. The loading was 0.64 mg cm⁻² in average, with a standard deviation of 0.03 for all prepared electrodes.

A custom gas-tight glass cell with two compartments separated by a Nafion 212 membrane (Alfa Aesar, 0.05 mm thickness) was employed for the electrochemical study. Both catholyte and anolyte chambers were filled with 40 cm³ of a 0.1 M KHCO₃ solution (Sigma-Aldrich, 99.95% trace metals basis) prepared with ultrapure water. Before (20 min) and

during the measurement, CO₂ (Messer, purity 4.8) was supplied to the catholyte at a flow rate of 20 cm³ min⁻¹, resulting in a pH level of 6.7. A Pt wire and an Ag/AgCl (3.0 M KCl) electrode were used as counter and reference electrodes, respectively. The geometric surface area of the working electrodes was in the range of 1–6 cm², accurately quantified for each electrode with the ImageJ image processing and analysis software (Wayne Rasband, National Institutes of Health). All measurements were performed using an Autolab PGSTAT302N potentiostat at room temperature. The double layer capacitance of the electrodes was assessed by cyclic voltammetry (CV) at varying scan rates in the potential range of open circuit potential (OCP) ±15 mV and performing a linear fit only in the region of scan rates where a straight line between electric charge and scan rate is observed. The performance of the prepared electrodes for eCO₂RR was investigated by chronoamperometry (CA) at –0.5 and –0.6 V versus RHE for 1.5 h. These potentials were chosen because they result in sufficient currents to quantify the formed eCO₂RR products with high precision while at the same time avoiding excessive Faradaic reduction currents, which would otherwise lead to mass-transport limitations and local pH changes, implying that the measured current would not only be controlled by electrochemical kinetics but by other factors as well.⁴⁰ CAs were performed with the *iR* compensation function set at 85% of the uncompensated resistance (*R_u*), which was determined before the start of the CAs and updated every 10 min by potentiostatic electrochemical impedance spectroscopy at the electrolysis potential (100 kHz, 10 mV amplitude). The remaining 15% of the *iR* loss was corrected manually by the *R_u* measured by the impedance spectroscopy. Current densities were normalized by the geometric surface area.

Gaseous products were analyzed using an on-line gas chromatograph (SRI 8610C, Multi-Gas #3 configuration) with Ar as a carrier gas, which was equipped with a HayeSep D column and a Molecular Sieve 13X column. Gas samples were injected through a sampling loop and analyzed 10 min after the start of the electrolysis and thereafter every 15 min. The Faradaic efficiency (FE) for the gaseous products was calculated with the following equation

$$\text{FE} = \frac{\text{gas flow through the cell} \times \text{concentration of the product}}{\text{electric current at sampling time}/nF} \times 100 \quad (2)$$

Liquid products were examined with a high-performance liquid chromatograph in a Merck LaChrom system equipped with a Bio-Rad Aminex HPX-87H column heated at 333 K and a refractive index detector (Hitachi Chromaster S450), using 5 mM H₂SO₄ as eluent. The FE for the liquid products was quantified by analyzing the catholyte after the test using the following equation

$$\text{FE} = \frac{\text{catholyte volume} \times \text{concentration of the product}}{\text{total charge passed during the electrolysis}/nF} \times 100 \quad (3)$$

In the above equations, *n* is the number of electrons transferred and *F* is the Faraday constant.

2.11. Computational Details. All calculations were performed using spin-polarized DFT as implemented in the Vienna Ab initio Simulation Package (VASP).^{41,42} The GGA

PBE-D3 was the functional of choice.^{43,44} Core electrons were described by the projector augmented wave,^{42,45} and valence electrons were expanded by plane wave basis sets with a kinetic energy cut-off of 500 eV. To model atomically dispersed metal–N_x carbon structures, we used graphene layer expanded in a (6 × 6) supercell replacing four C atoms by nitrogens and removing two C atoms to compensate for the valence. These layers were interleaved by 12 Å vacuum along the *z* direction. Various transition metal atoms (M = Mn, Fe, Co, Ni, Cu, Zn) were placed in the cavity left on the carbon sheet. In this configuration, two axial positions are empty. However, most of these atoms prefer octahedral coordinations,⁴⁶ and thus water-derived ligands (L = none, H₂O, OH, O) typically occupy one of the axial positions in aqueous electrolytes. The empty coordination site is then employed to investigate the competitive eCO₂RR and HER. Structures were relaxed with a force threshold of 0.05 eV/Å. The Brillouin zone was sampled using a 3 × 3 × 1 *k*-point mesh generated with the Monkhorst–Pack method.⁴⁷ To obtain the energies of the relevant intermediates in the electrochemical environment, the computational hydrogen electrode (CHE) was used.^{48,49} The speciation of the different oxidation states of the metals in the carbon lattice was performed via a Pourbaix diagram derived with Phonopy.⁵⁰ The optimized structures can be retrieved from ioChem-BD.^{51,52}

3. RESULTS

3.1. Structural Characterization of Pristine Catalysts.

Leaning on the synthetic route developed by our group for the preparation of FeNC and CoNC materials that exclusively comprise MN_x moieties,^{24,25} we expanded it to prepare MnNC, NiNC, and CuNC materials. A control sample was synthesized by the same approach without the addition of divalent metal acetate, and was labeled (Zn)NC, with (Zn) indicating the possible presence of residual Zn from ZIF-8, as discussed later. The MNC materials were first characterized by HAADF-STEM. Multiple areas of the samples were examined, and no metallic particles could be observed (Figure 1). The absence of crystalline metallic species in as-prepared MNC materials was further confirmed by powder XRD (Figure S1), exhibiting only two broad reflection peaks typical for nanometric graphite-like domains present in amorphous N-doped carbon materials. Elemental mapping with EDX spectroscopy showed a homogenous distribution of the metal, C and N atoms in the catalysts, revealing also the presence of residual Zn from ZIF-8 in all five MNC catalysts (Figure S2). These STEM and EDX analyses indicate that the metals are probably atomically dispersed in the nitrogen-doped carbon matrix. However, because of limited resolution of the EDX mapping, nano- or sub-nano-sized metal clusters cannot be excluded at this stage.

In order to unambiguously demonstrate the absence of metal clusters and to identify the metal coordination and bond distances in atomically dispersed MN_x sites, we resorted to EXAFS spectroscopy (Figure 2). The Fourier transform (FT) of the ex situ EXAFS spectra [*k*² weighted $\chi(k)$] of all MNC catalysts at their respective metal *K*-edge energy exhibited a major peak at ~1.5 Å, attributed to the backscattering by light atoms (C, N, or O) situated in the first coordination shell of the absorbing metal. The secondary peak(s) at 2–3 Å is associated with the contribution of carbon (or N, O) atoms in the second coordination sphere of the metal center. The FT-EXAFS spectra of the MNC catalysts agree well with those of

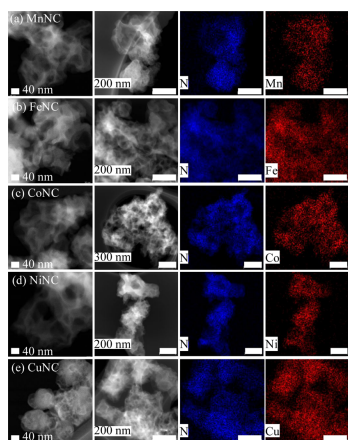


Figure 1. MNC catalyst morphology: high-resolution TEM characterization. TEM images and EDX elemental mappings for N and the respective metal derived from the metal-acetate precursor of (a) MnNC, (b) FeNC, (c) CoNC, (d) NiNC, and (e) CuNC.

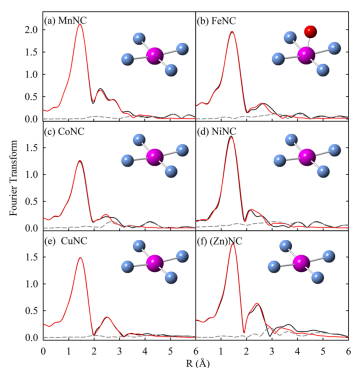


Figure 2. Atomically dispersed MN_x sites: metal K-edge EXAFS analysis. (a) MnNC, (b) FeNC, (c) CoNC, (d) NiNC, (e) CuNC, and (f) ZnNC. The black curves represent the experimental spectra, whereas the red curves represent calculated spectra based on the depicted structures. The dashed gray curves are the difference between experimental and calculated spectra. Metal, nitrogen, and oxygen atoms are represented in purple, blue, and red, respectively (carbon atoms in the second coordination sphere are not shown). No phase-shift correction was applied to the FTs.

some atomically dispersed MNC materials previously reported by other groups.^{13,2b}

The metal K-edge EXAFS spectra were then fitted with the model structures depicted in the insets of Figure 2. For CoNC, NiNC, and CuNC, a variable number of in-plane nitrogen

atoms binding the metal–cation center were allowed for the fitting, while the spectra of FeNC and MnNC were fitted assuming the presence of four in-plane nitrogen atoms, and either one or two oxygen atoms as axial ligands. This approach was used to reduce the number of EXAFS parameters in the fit, as FeN₄ and MnN₄ moieties have a higher affinity to oxygen compared to other metal–N₄ moieties.⁵³ The structural parameters obtained from the fittings are shown in Table 1, including the coordination number (CN) and M–N bond distances. The best-fit results show a CN-value of ~4 for MnNC, CoNC, NiNC, and CuNC, suggesting that the vast majority of metal cations in these catalysts could be involved in MN₄ moieties. EXAFS alone however cannot distinguish between N, C, and O atoms, which means that the CN value of 4 may also be acceptable, from a structural viewpoint, with a mix of N, C, and O atoms. From a chemical viewpoint however, metal–nitrogen bonds in the first coordination sphere are more likely than metal–carbon or metal–oxygen ones. For Fe and Co, time-of-flight secondary-ion mass spectroscopy revealed the major presence of MN₂C₂⁺, MN₂C₂⁺.⁵⁴ Experimentally, it has also been shown that nitrogen is critical for the formation of single Fe atom site in a carbon support during pyrolysis, only metallic Fe particles being observed in the absence of nitrogen in the synthesis.⁵⁵

For other metals of the 3d row, although less experimental characterization has been reported yet, DFT predicts higher energy stability for metal–N₄ defects in graphene than metal–C₄ defects.⁵⁶ For those reasons, we started by investigating in-plane MN₄ sites in our DFT study to explain the reactivity, as reported in Section 4.2. With the starting hypothesis of the in-plane FeN₄ site, the higher average CN-value of 5–6 for FeNC^{22,25} strongly suggests that one or two axial O atoms are adsorbed on top of FeN₄ moieties, resulting in coordinatively saturated iron cations, in line with the high oxophilicity of Fe.

It should be noted that ZnN_x moieties with an average CN of ~4 are present not only in the control sample (Zn)NC (Figure 2f and Table 1) but also in all five MNC samples. The Zn K-edge X-ray absorption near edge structure (XANES) and EXAFS spectra of the five samples are identical to that recorded for (Zn)NC (Figure S3), implying that the ZnN_x coordination is unaffected by the presence of a second metal. The coexistence of ZnN_x and MN_x moieties is another indication that the metal content (from metal acetate) is not saturating the material, explaining the absence of metallic particles. In summary, the EXAFS spectra were fitted with in-plane MN₄ structures, with an axial oxygen atom needed to obtain a best-fit only for FeNC. Most important, no M–M backscattering signal was needed to get an excellent EXAFS fit, unambiguously indicating the isolated nature of the metal centers.

The metal–nitrogen bond distance and K-edge EXAFS spectra were also calculated ab initio for MN₄C₁₀ model sites (10 in-plane carbon atoms in the second coordination sphere) in graphene sheets, with or without axial oxygen adsorbates (Figure S4). Considering water as an oxygen ligand, MnN₄ (2.40 Å), FeN₄ (2.33 Å), CoN₄ (2.31 Å), and ZnN₄ (2.35 Å) exhibited similar M–O bond distance between 2.31 and 2.40 Å (Table S1). No water could be stabilized on the CuN₄ and NiN₄ models. The DFT-calculated EXAFS spectra shown in Figure S5 have characteristic features analogous to those of the experimental FT-EXAFS spectra in Figure 2 with a main peak assigned to N or O back-scattering from the first coordination sphere. The calculated and experimental M–N distances are

Table 1. Structural Parameters Obtained from the Analysis with MN_x Moieties of the Metal K-Edge EXAFS Spectra and DFT Calculations^a

	MnNC	FeNC	CoNC	NiNC	CuNC	(Zn)NC
R_{M-N} (Å)	1.89(1)	1.99(2)/2.01(2)	1.96(2)	1.92(1)	1.94(1)	2.00(1)
σ^2 (Å ²)	0.0033(4)	0.008(2)	0.008(2)	0.007(1)	0.0072(7)	0.0056(3)
CN	4	5/6	3.8(6)	4.3(3)	3.9(3)	4.1(4)
R_{M-N}^D (Å)	1.91	1.89	1.89	1.88	1.93	1.96

^a R_{M-N} is the metal–nitrogen bond distance, σ^2 the Debye–Waller factor, and CN the coordination number, obtained from EXAFS fitting, whereas R_{M-N}^D is the DFT simulated metal–nitrogen bond distance for MN_xC₁₀ model sites. Errors are given in brackets, for example, 1.89(1) means 1.88–1.90.

also in good agreement, except for Fe, for which the EXAFS fitting is complicated by the presence of axial oxygen ligands (Tables 1 and S1).

The oxidation state of the atomically dispersed metal centers was then investigated with XPS. No signal related to metal(0) state could be observed in the M 2p spectra. The narrow-scan spectra in the binding energy regions corresponding to the M 2p_{3/2} spin–orbit couplings are shown in Figure 3. They

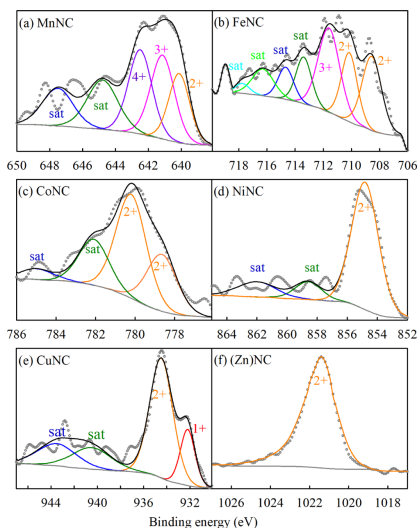


Figure 3. XPS characterization. Metal 2p spectra of (a) MnNC, (b) FeNC, (c) CoNC, (d) NiNC, (e) CuNC, and (f) (Zn)NC.

identify the coexistence of M²⁺ and M³⁺ species for FeNC, and the coexistence of M⁺ and M²⁺ species for CuNC.^{57,58} Only the 2+ oxidation state is observed in CoNC, NiNC, and (Zn)NC, whereas for MnNC, the three different oxidation states 2+, 3+, and 4+ seem to be present.^{59–61} The existence of different oxidation states of the metal in a same catalyst reveals a structural and/or a topological heterogeneity of MN_x moieties. For example, higher oxidation states of the metal–carbon might be triggered by oxygen (O₂, OH) adsorbates, which is possible only for MN_x moieties located on the top surface, whereas those located in the bulk of the nitrogen-

doped carbon matrix are not accessible. Chemical heterogeneity might also lead to different oxidation states for a same structure. For example, whereas an average MN₄ structure was identified from our EXAFS analysis for all materials except FeNC (with additional oxygen adsorbates on top of the FeN₄ site itself), EXAFS cannot distinguish between N, O, and C atoms in the first coordination sphere, implying that, for example, MN₂O₂ moieties or any other combination with a total of four in-plane ligating atoms among C, N, and/or O are also possible structural candidates from an EXAFS viewpoint. Whereas the energetic stability of such moieties is usually lower than MN₄,^{62,63} their existence as minor species cannot be excluded. For MnNC in particular, the presence of high valence-state metal (4+) suggests that there might be some O atoms replacing N atoms in MnN₄ moieties. This will be re-discussed later.

The elemental content of the catalysts was quantified from XPS signals. All MNC catalysts have similar C, N, and O absolute contents, whereas the metal contents are in the range 0.16–0.27 at. %, except Zn in the range of 0.11 to 0.38 at. % (Table S2). The average nitrogen chemistry is similar for all five MNC catalysts with similar relative concentrations of different N species observed, as shown in Figure S6 and Table S3. Nitrogen–transition metal interactions are being expressed as a singular peak in the N 1s fine resolution spectra. Interactions of transition metal with 2, 3, or 4 co-planar N atoms are indistinguishable within the resolution of the peak composition analysis, as well as different transition metals studied here produce a de-facto identical shift, and the spectral component is commonly addressed here as M–N_x (Figure S6). Nevertheless, a positive correlation was found between the N content and the total metal content of selected samples (Figure S7), with the linear fitting results for selected samples CuNC, ZnNC, NiNC and CoNC. MnNC was not included as we demonstrate that the first sphere coordination of Mn probably involves both O and N atoms (details will be discussed in Section 4.2), and FeNC was excluded from the linear fitting as it is, for an unknown reason, an outlier. Most interestingly, the fitted line in Figure S7 has a slope of 3.34 N atoms per metal atom, nicely supporting the fixation of close to four nitrogen atoms during pyrolysis per each single metal atom, and the major presence of MeN₄ sites for the (Zn)NC, CoNC, NiNC, and CuNC materials. The position of MnNC data point below the line also supports the mix of O and N in the first coordination sphere of Mn.

SEM (Figure S8) and N₂ physisorption isotherms (Figure S9) indicate a similar porous carbon structure in the six materials. The isotherm shape indicates a high microporosity (vertical rise at low P/P₀) and the existence of small mesopores that are responsible for the hysteresis. Similar porous structures had been observed on analogous ZIF-8-derived materials,^{22,23}

and micropores were shown to be important to host MN_x active sites.^{64,65} The Brunauer–Emmett–Teller surface areas of the catalysts are similar, in the range of 220–330 m² g⁻¹ (Table S4). Similar double layer capacitances according to CV further confirm that MNC and (Zn)NC materials possess a similar electrochemical surface area (Figure S10 and Table S4). As shown in Figure S1, the XRD patterns of the MNC catalysts revealed only two broad reflection peaks attributed to the (002) and (101) facets of nanometric graphitic structures.^{22,66} In the Raman spectra (Figure S11), all samples showed the same ratio of D to G band intensity ($I_D/I_G \approx 0.8$), suggesting the same degree of disorder of the carbon phase. In summary, the five MNC catalysts have similar structure and composition, with the only significant difference being the nature of the metal and its exact coordination and oxidation state in MN_x moieties.

3.2. Selectivity and Activity of MN_x Sites for eCO₂RR

The catalytic activities of MNC and (Zn)NC toward eCO₂RR were first screened with LSV. The five MNC materials and also (Zn)NC exhibit higher current densities in CO₂-saturated 0.1 M KHCO₃ than in N₂-saturated 0.1 M K₂HPO₄/KH₂PO₄ buffer electrolyte (Figure S12). The onset potential is highest for FeNC, whereas the current density reached with NiNC surpassed that reached with FeNC below -0.8 V versus RHE (the reasons for the high onset of FeNC will be discussed later).¹⁴ CoNC shows even higher current density than NiNC at a low potential but, because of similarity in the polarization curves in CO₂-saturated and CO₂-free electrolytes (Figure S12c), this might mostly be attributed to HER. These trends agree with previous reports.^{14,15,20,21} The identification and quantification of the eCO₂RR products were then carried out under potentiostatic control. The FEs obtained from CA at -0.5 and -0.6 V versus RHE vary with the embedded metal center, as shown in Figure 4. Although CO was the major

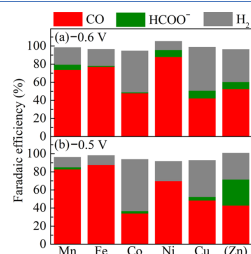


Figure 4. Product analysis. FE at (a) -0.6 V and (b) -0.5 V vs RHE obtained over the MNC catalysts (M = Mn, Fe, Co, Ni, and Cu), plotted against the embedded metal center sorted by the atomic number. The data were obtained by CA for 90 min, performed in 0.1 M KHCO₃ aqueous electrolyte saturated with CO₂ (pH 6.7) at room temperature.

eCO₂RR product over all MNC, the highest FEs were observed over the Fe-, Ni- and Mn-NC materials. High FEs have been reported previously for other FeNC,^{10,17,20,87} and NiNC^{13,16,18,19,38} materials synthesized via different approaches but sharing a key common structural feature, namely, the atomically dispersed MN_x moieties. Also, in agreement with other recently published results, the major

product over CoNC is H₂.^{14,15,20,21} On the contrary, high selectivity (more than 90%) toward CO₂ reduction to CO in a wide potential range (-0.57 to -0.88 V vs RHE) over atomically dispersed CoN_x was reported as well,^{68,69} but with slightly different coordination environment. As the present series of MNC and (Zn)NC catalysts has, except for the nature of the metal element in MN_x moieties, negligible difference in morphology and structure, the different activity and selectivity toward eCO₂RR to CO can be mainly ascribed to the nature of the metal in MN_x sites. Last, (Zn)NC shows a selectivity for CO being in-between the low value of CoNC and the high values of MnNC, FeNC, and NiNC.

Although Figure 4 indicates the relative fraction of all products formed during CA on each of the catalysts, it does not convey any information on the rate at which these products formed. Therefore, the partial current densities for each of the products (CO, HCOO⁻, and H₂) were calculated from the CA data, and are reported in Figure 5. They were

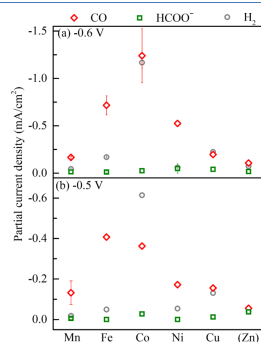


Figure 5. Volcano trend in eCO₂RR activity. Partial current densities at (a) -0.6 V and (b) -0.5 V vs RHE over the MNC catalysts for CO, HCOO⁻, and H₂ products. The data were obtained by CA for 90 min, in 0.1 M KHCO₃ saturated with CO₂ (pH 6.7) at room temperature.

calculated from the knowledge of (i) the average current density on each catalyst during the CA experiment and (ii) the product quantification averaged during the CA (see Experimental Methods). Looking at the results obtained at -0.6 V versus RHE (Figure 5a), it should first be noted that, whereas (Zn)NC shows non-zero activity toward CO₂ reduction, it is negligible compared to the activity observed for FeNC, CoNC, and NiNC. The low activity of (Zn)NC is attributed to ZnN₄ moieties^{8,70} and/or nitrogen groups not bonded to any metal.^{71,72} Given that MNC catalysts and (Zn)NC share a similar carbon matrix with similar amounts of ZnN₄ moieties and nitrogen functionalities, we assign the higher partial current densities for CO formation (j_{CO}) of MNC samples to the presence of MN_x moieties (Fe-, Co-, and NiN_x moieties in particular). A catalytic synergy between MN_x and ZnN₄ moieties is also unlikely when the main considered product of CO₂ reduction is CO (only two electrons transferred, with COOH* as the only possible intermediate species). There are generally three steps for the eCO₂RR to CO: (i) CO₂ + H⁺ + e⁻ → COOH* (this first electron and proton transfer can be

further decoupled into two steps: (i-a) $\text{CO}_2 + \text{e}^- \rightarrow \text{CO}_2^{\bullet-}$, and (i-b) $\text{CO}_2^{\bullet-} + \text{H}^+ \rightarrow \text{COOH}^*$, as discussed in the following DFT simulations), (ii) $\text{COOH}^* + \text{H}^+ + \text{e}^- \rightarrow \text{CO}^* + \text{H}_2\text{O}$ and (iii) CO^* desorption. If a synergistic effect between MN_x and ZnN_4 moieties occurs, there are two possibilities: (a) MN_x catalyzes step (i) and ZnN_4 catalyzes step (ii); or (b) the opposite. Assuming the synergy case (a), the FE for formate on FeNC, CoNC and NiNC should be significantly higher than that on (Zn)NC, which is not the case (Figure 4). Assuming the synergy case (b), (Zn)NC should have a J_{HCOO^-} value comparable to J_{CO} observed on the active catalysts NiNC, FeNC, and CoNC. However, this is also in contrast with experimental results, where J_{HCOO^-} on (Zn)NC is negligible compared to J_{CO} on NiNC, FeNC, and CoNC (Figure 5). Thus, we can rule out any significant catalysis synergy between MN_x and ZnN_4 moieties.

Interestingly, with increasing atomic number from Mn to Co, J_{CO} monotonically increases, whereas with further increasing atomic number of the metal, the J_{CO} value decreases (Figure 5a). A volcano-shaped trend is thus identified in the electrocatalytic activity for the reduction of CO₂ to CO. A similar volcano trend is also apparent at -0.5 V versus RHE, but with FeNC exhibiting a slightly higher J_{CO} value than CoNC at this potential (Figure 5b). The same volcano trends are observed if J_{CO} is normalized by the electrochemical double layer capacitance (Figure S13b) or by the metal content obtained from XPS (Figure S13c). These observations indicate that MN_x moieties embedded in the nitrogen-doped carbon matrix exhibit not only different selectivity but also distinct activities toward eCO₂RR. Herein, we for the first time reveal a volcano-shaped trend in the electrocatalytic activity for the CO₂ to CO reduction over a series of pyrolyzed MNC catalysts, with CoNC and FeNC located at the volcano peak at -0.6 and -0.5 V versus RHE, respectively. The underlying reasons for this volcano-trend are discussed in the next section.

The DFT simulations were performed on the most likely oxidation state and environment for each of the metals. Although it was characterized via ex situ XPS, the oxidation state under operating conditions might be very different.^{24,67} We therefore resorted to operando XANES spectroscopy to track possible changes in the oxidation state of MN_x sites during eCO₂RR. The overlapping XANES spectra at OCP before and after CA at 0.0, -0.5 , -0.6 and -1.0 V versus RHE demonstrate that no irreversible structural change occurred for the MN_x sites during eCO₂RR (Figure 6a), in agreement with operando XAS results on other MNC catalysts previously reported.^{24,67} The position of the absorption edge at the metal K-edge XANES spectra is generally recognized to correlate, for a given coordination environment of the metal, to the metal oxidation state. The absorption edge position was therefore determined as the maximum of the first derivative of the XANES spectra, and is shown as a function of the electrochemical potential in Figure 6b. The figure indicates that a reduction occurred for Fe, Ni, and Cu sites when decreasing the potential from OCP to -0.5 V versus RHE, whereas the oxidation state of Mn and Co sites remained the same. Combined with the ex situ XPS oxidation state identification, we can infer that the oxidation state of Fe was modified from 2+/3+ to only 2+, that of Ni reduced from 2+ to 1+, and that of Cu reduced from 2+/1+ to 0.^{14,73} The oxidation state of Mn remained a mixture of 2+/3+/4+, and that of Co remained as 2+, even at -1.0 V versus RHE. It can further be observed that Fe was reduced from a mix of 2+/3+

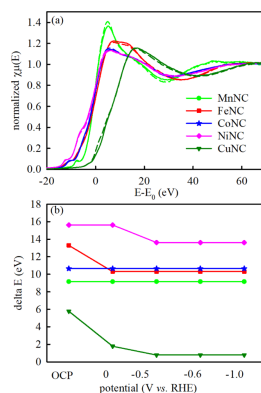


Figure 6. (a) Metal K-edge XANES spectra of MNC catalysts at OCP before and after the CA measurements at different potentials (solid curve: before; dashed curve: after), and (b) change in the absorption edge of metal K-edge XANES spectra of MNC catalysts with applied potential. The electrolyte was CO₂-saturated 0.1 M KHCO₃, and the OCP was around 0.7 V vs RHE. In (a), the x-axis is the difference between the X-ray energy and the absorption edge of MNC (E_0). In (b), the y-axis (ΔE) is ($E_0 - E_{0,\text{metal}}$).

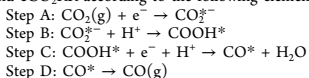
to only 2+ already at 0 V versus RHE, in line with the redox potential identified for FeNC at ca. 0.75 V versus RHE in acidic electrolyte.^{22,74} This may explain the high onset potential for eCO₂RR over FeNC, whereas NiNC is possibly inactive for eCO₂RR as long as it is in 2+ oxidation state (discussed later in combination with DFT). The changes in oxidation states of the three most active MNC catalysts (namely Fe, Co, and Ni) with applied potential agree well with the calculated Pourbaix speciation diagrams, which shows that in a pH = 7 electrolyte, the oxidation state of Fe and Co in MN_x corresponds to Co^{2+} at any potential < 0 V versus RHE, whereas there is a reduction from Ni^{2+} to Ni^{1+} below 0 V versus RHE at pH = 7 (Figure S14), in line with the Ni-reduction experimentally observed between 0 and -0.5 V versus RHE experimentally (Figure 6b). It should be noted that XAS is a bulk technique and therefore the signal and values reported in Figure 6 are an average of all metal atoms of the same nature in the sample, regardless of whether they are on the surface (sensitive to the electrochemical potential) or in the bulk (possibly insensitive to the electrochemical potential).

4. DISCUSSION

4.1. Explanation for the Slightly Different Volcano Trends at -0.5 and -0.6 V Versus RHE. Figure 5 reveals that the most active catalyst for eCO₂RR to CO depends on the electrochemical potential. In the general case, such a potential-dependent volcano trend can be expected, as the electrochemical kinetics of even a simple one-electron reaction is governed by at least two parameters, namely, the exchange current density (i_0) and the Tafel slope. Only if the Tafel slope for eCO₂RR to CO is the same on all MNC catalysts can one expect a potential-independent activity trend. This is clearly shown in Figure S15, constructed from Tafel laws and

assuming different i_0 values and either a common Tafel slope (Figure S15a) or slightly different Tafel slopes (Figure S15b). The plot of the kinetic current density versus potential show the same trends at -0.5 and -0.6 V versus RHE when the Tafel slope is assumed to be the same for all three catalysts (Figure S15c), whereas the trends differ when different Tafel slopes are assumed (Figure S15d). Thus, the simplest explanation to account for the slightly different volcano trends in Figure 5a,b is a variation of the eCO₂RR Tafel slope from one catalyst to another. This may, in turn, be related to different rate-determining steps of eCO₂RR on those different catalysts. From the theory of electrokinetics of multielectron reactions, higher Tafel slopes are generally expected when the rate-determining step (RDS) is the first electron transfer than when it is the second electron transfer or any reaction step after the second electron transfer.⁷⁵

4.2. Underlying Reason for the Volcano Trends. The observation of volcano trends in the electrochemical activity for various reactions among a series of metal-based catalysts with similar site structures is generally explained by Sabatier's principle, with the binding energy of the metal sites being tuned by the metal's atomic number, and leading to too weak, optimum, and too strong binding of the key intermediate-adsorbed species when moving in a given direction in the series of transition metal.^{76,77} We then performed DFT calculations with MN₄C₁₀ model (M = Mn, Fe, Co, Ni, and Cu, with axial oxygen adsorbates for some of the metals) with the CHE model at $U = -0.6, -0.5,$ and 0.0 V versus RHE, to investigate the binding energy of the key reaction intermediates for HER and eCO₂RR according to the following elementary steps



The reaction energy profiles for each elementary step of eCO₂RR (Figure 7 ($U = -0.6$ and -0.5 V vs RHE), Figure S16 ($U = 0.0$ V vs RHE) and Tables S5–S7) show that the particular transition metal present in the N₄ lattice has a tremendous effect on reactivity. As eCO₂RR does not occur at $U = 0.0$ V versus RHE, we focus in our discussion on the free energy diagrams at $U = -0.6$ and -0.5 V versus RHE (Figure 7), potentials at which we experimentally measured the eCO₂RR activity and selectivity of MNC catalysts. Overall, the calculated Gibbs free energy diagrams demonstrate that MnN₄ and FeN₄ bind CO too strongly and their CO production rate is determined by the CO* desorption step. In contrast, NiN₄ and CuN₄ have weak CO binding energy, and the CO₂ activation and first electron transfer then becomes the RDS (Table S8). In this series of model sites, CoN₄ shows the optimum balance between the energy barriers for CO₂ activation and CO desorption (steps A and D in Table S8, respectively).

We also investigated how axial adsorbates (H₂O, OH, and O) affect the reaction energies on MnN₄, FeN₄, and CoN₄ (no oxygen adsorbates could be stabilized on CuN₄ and NiN₄). Particularly, the H₂O molecule axially adsorbed on FeN₄ lowers the reaction barrier for the first electron transfer to form CO₂^{•-} (Figures 7b and S14) and thus promotes the eCO₂RR to CO over FeNC. In this case, FeN₄ sites are approachable from both axial directions (one side for adsorbing H₂O, and the other side for activating CO₂ gas molecule). Moreover, for the case of NiNC, NiN₄C₁₀ with oxidation states of 1+ and 2+ are both used in the DFT

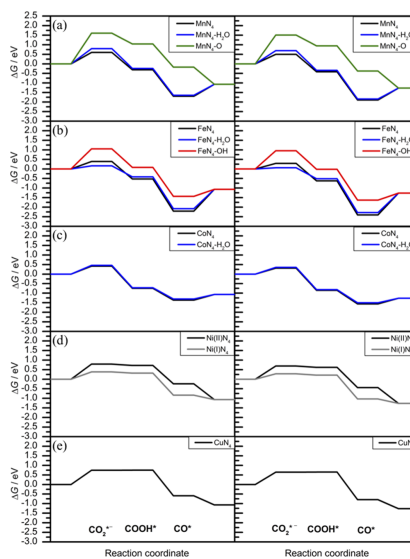


Figure 7. Gibbs free energy profiles for eCO₂RR over MN₄ model sites ((a–e) M = Mn, Fe, Co, Ni, and Cu) at $U = -0.5$ V (left) and -0.6 V (right) vs RHE (pH = 7) according to the different metal speciation (oxidation state and ligand coordination).

simulations, as a reduction of Ni²⁺ to Ni¹⁺ was observed from operando XAS (Figure 6). In contrast to the high energy barrier for the CO₂ activation on Ni²⁺N₄C₁₀ (the first step for eCO₂RR), Ni¹⁺N₄C₁₀ binds CO₂ more strongly (Figure 7d). This explains that the onset potential of eCO₂RR over NiNC is highly correlated to the Ni^{2+/1+}N₄ redox potential (Figures 6b and S11d). The CO₂ activation energies are also highly correlated to the geometries of the CO₂ coordination over MN₄ sites, as depicted in Figure S17 showing how CO₂ is activated in the Fe, Co, and Ni sites. Whereas both a M–C and a M–O bonds can simultaneously form with the FeN₄ active site, only a M–C bond appears for the CoN₄ and NiN₄ sites. In addition, the O–C–O angle departs from the 180° value in the gas phase to only 145° when adsorbed on FeN₄. This explains the lower activation energy of CO₂ over FeNC versus other MNC catalysts and the higher onset potential for FeNC.

The activity trends for CO₂ reduction to CO can be estimated from DFT calculations on the basis of the exergonic value of the reaction energy at the RDS. Table S8 reports the reaction energy calculated for each elementary step for MN₄(–H₂O) moieties according to the Gibbs free energy profiles for eCO₂RR over MNC catalysts (M = Mn, Fe, Co, Ni, and Cu) at $U = -0.6$ V versus RHE. For each MN₄ site, the structure (with/without H₂O adsorbate) with lowest reaction energy at the RDS is considered in Table S8. To compare DFT results to the experimental eCO₂RR activity trend, we assumed an Arrhenius type law of the form $A \exp(-E_{\text{RDS}}/RT)$, where E_{RDS} is the free energy change at the RDS and A is a pre-

exponential factor, assumed the same for all MNC catalysts investigated because of similar metal content and similar carbon porosity, chemistry, and morphology. Figure 8 shows

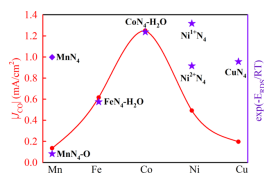


Figure 8. Comparison of experimental CO partial current density at -0.6 V vs RHE for pyrolyzed MNC materials and DFT calculated trends ($U = -0.6$ V vs RHE). E_{RDS} is the calculated free reaction energy at the specific RDS of eCO₂RR to CO for each active site (see Table S8).

that the trend of reaction rate predicted from DFT according to this approach matches well the experimentally observed volcano trend in the present work, except for MnN₄ and CuN₄, which will be discussed later. Among all the MN₄(-H₂O) active sites, Ni¹⁺N₄ is predicted to show the highest activity toward eCO₂RR to CO, with lowest maximum for the free reaction energy at any of the elementary steps. However, the experimental oxidation states of NiN₄ sites are most likely in a mix of 1+/2+ at -0.6 V versus RHE (Figures 6 and S12). Thus, the experimental activity of NiNC should fall between those predicted for Ni¹⁺N₄ and Ni²⁺N₄. Experimentally, the eCO₂RR activity of NiNC is even lower than that estimated

from DFT for Ni²⁺N₄ however. We assign this to a lower onset potential for NiNC than for CoNC, but once the reaction has started, recent studies have shown that NiNC catalysts with single metal atom sites are the most active and selective MNC catalysts reported to date.^{13,16,18,19,38} The calculated activity trends for MN₄ sites are as following FeN₄-H₂O < MnN₄ < CoN₄-H₂O < Ni¹⁺N₄ > Ni²⁺N₄ ≈ CuN₄, with metals on the left-hand side of Co or Ni binding CO too strongly, and metals on the right-hand side binding CO₂ too weakly.

One discrepancy between the present experimental results and theoretical results is MnNC, for which our experimental results show much lower CO partial current density than FeNC, whereas the DFT calculations predicted that MnN₄ and FeN₄ sites are both limited by the CO* desorption step, but with lower desorption barrier for MnN₄. This results in a predicted eCO₂RR activity higher for MnN₄ than FeN₄ moieties (Figure 8).

Given that a higher oxidation state of Mn (a mixture of 4+, 3+, and 2+) was observed by ex situ XPS and operando XANES than can be assumed with a Mn₄ model site, and given that MnN₄ has a higher predicted affinity to oxygen,⁵³ we re-evaluated the possible active site structure of MnNC. We assumed MnN₂O₂ as an alternative model site (Figure S18a). Both MnN₄ and MnN₂O₂ structures converged and optimized under the same conditions, confirming similar energy stabilities of these two structures. Moreover, the experimental FT-EXAFS spectrum of MnNC could be properly fitted with the MnN₂O₂ model site (Figure S18b), and its DFT calculated EXAFS spectrum (Figure S18c) exhibits the same characteristic peak at ~ 1.5 Å assigned to both N and O back-scattering from the first coordination sphere. According to our DFT

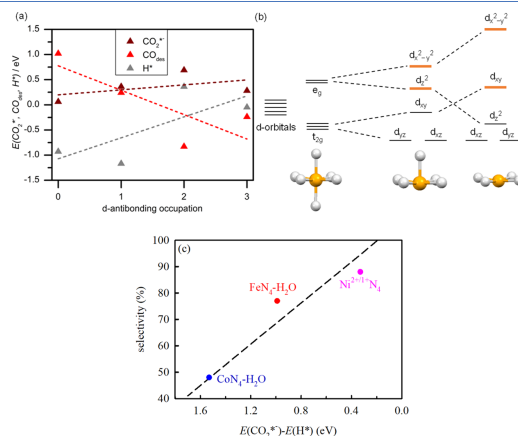


Figure 9. Relationship between the adsorption/desorption energies and the local electronic configurations on the Fe, Co, and Ni (1+ and 2+). (a) CO₂* (wine), H* (gray) adsorption energies and CO (red) desorption (at pH = 7, U = -0.6 V) as a function of the d-antibonding occupation in the active site. In (b), the symmetry descent from the octahedral configuration to either square pyramidal or square planar is shown (orange indicates the antibonding states). For the most common Fe species under potential Fe and Co are in MN₄-H₂O (square pyramid with 0 and 1 antibonding electrons). Instead, Ni is in the square planar configuration with already two electrons in the antibonding state. At high negative potential the highest antibonding d_{x²-y²} would be occupied, which could trigger the distortion of the planar pocket (Ni¹⁺ state). (c) The experimental selectivity at -0.6 V vs RHE over FeNC, CoNC, and NiNC vs the DFT-simulated $E(\text{CO}_2^*) - E(\text{H}^*)$.

calculations, this MnN₂O₂ model can however not coordinate CO₂ and is therefore predicted to be inactive toward eCO₂RR. Even with a single oxygen adsorbate, the MnN₄ site becomes poorly active toward eCO₂RR, as indicated by the very high energy barrier for CO₂ adsorption calculated by DFT for the MnN₄-O model (Figure 7a). Of all the Mn models investigated, only Mn²⁺ without oxygen ligands can effectively coordinate CO₂. However, as identified by ex situ XPS and operando XANES (Figures 3 and 6), this species is uncommon at the investigated potentials, explaining why no activity is experimentally found for MnNC. Last, for CuNC, our operando XANES data show that Cu switches from 2+ to 0 oxidation state early on (Figure 6), explaining the mismatch between DFT predictions and the experimental CuNC activity. In summary, with the MnN₄-O model site for MnNC and MN_xC₁₀(-H₂O) model sites for other MNC catalysts (M = Fe, Co, Ni), the experimental volcano trend in the eCO₂RR to CO activity of MNC catalysts can be nicely explained by our DFT-calculated Gibbs free energy for the RDS (Figure 8).

The volcano-shaped trend in eCO₂RR activity to CO over various MNC catalysts could not be identified in other reports, where CoNC exhibited lower *J*_{CO} values relative to other MN_x moieties (Ni and Fe in particular) in a broad potential range.^{14,15,21} The reason for this seems to be that previous reports either lacked evidence for the exclusive atomically dispersed nature of transition metals (the XRD for some of the MNC catalysts exhibited sharp peaks arising from metallic nanoparticles, which might alter the carbon structure),^{14,15} or studied the activity of MNC catalysts varying not only in the nature of the transition metal but also in the carbon matrix structure (as shown in XRD and/or N₂ adsorption/desorption isotherms).²¹

4.3. Selectivity toward CO over MN_x. As another critical factor for a good eCO₂RR catalyst, selectivity is highly correlated to the H* and CO₂* binding energies over MN_x active sites. Thus, we investigated the Gibbs free energy profiles of HER over MN_xC₁₀ model sites (M = Mn, Fe, Co, Ni, and Cu) at -0.5 and -0.6 V versus RHE (Figure S19).

These binding energies can be rationalized from simple coordination chemistry rules. The adsorption behavior for the metals in the first transition series depends on the occupation of antibonding d-orbitals coming from the crystal field splitting for the resting square pyramidal structure. Therefore, we use the number of d-antibonding electrons (0 for FeN₄-H₂O, 1 for CoN₄-H₂O, 2 for Ni¹⁺N₄, and 3 for Ni²⁺N₄) to explain the selectivity toward eCO₂RR to CO over MN_x sites (Figure 9). The population of the antibonding states results from counting the number of electrons in the d-orbitals of the ground-state configuration according to the crystal field splitting theory. Figure 9a indicates that there is a linear scaling relationship between the CO₂* and CO* binding energies (Figure S20), which determines the CO activity. Thus, the CO binding energy can serve as a key descriptor for eCO₂RR to CO activity. The volcano trend with CO binding energy as a descriptor (*J*_{CO} measured experimentally vs CO binding energy calculated from DFT simulations) is shown in Figure S21, with Fe located at the strong binding branch, whereas Ni and Cu are on the weak binding branch. CoN₄-H₂O sites possess the optimum binding energies with CO*, which are neither too strong nor too weak. On the other hand, good selectivity toward CO requires high CO₂* binding energy and low H* binding energy (the more negative value in Figure 9a, the higher the binding energy). Thus, we utilized the difference of

CO₂* and H* binding energies ($E(\text{CO}_2^*) - E(\text{H}^*)$) as the descriptor for selectivity (Figure 9c). In that figure, we show only the experimental data for FeNC, NiNC, and CoNC and did not include the other MNC materials, whose eCO₂RR activity and selectivity can originate in whole or in part from metal-free nitrogen functional groups.⁷⁸ The CoN₄-H₂O site has an optimum binding energy with CO₂*, but, unfortunately, the binding energy of H* is much higher than that of CO₂*, which leads to a low selectivity toward CO, in line with the experimental result on CoNC (48%, Figure 9c). The high selectivity of eCO₂RR to CO over NiN₄ benefits from the lower H* binding energies compared to other MN_x sites (Figure 9c). For FeN₄-H₂O, even though the high binding energy to H* suggests a high HER activity (Figure 9a), the reduction of CO₂ to CO* also occurs with zero energy barrier, and the binding energy of CO* on FeN₄-H₂O is much stronger than that of H* on FeN₄-H₂O. Thus, the majority of FeN₄ sites are occupied by CO* in CO₂-saturated electrolyte. This site-blocking effect probably suppresses the binding of such active sites with HER-related species, resulting in high experimental selectivity toward CO over FeNC. As can be seen in Figure 9c, the selectivity of eCO₂RR toward CO over MN_x moieties can be explained with the DFT-calculated CO₂* and H* binding energies, that in turn are highly correlated to the number of d-antibonding electrons. The latter descriptor can thus be used to screen for more promising eCO₂RR candidate MNC materials.

5. CONCLUSIONS

The careful integration of experimental and theoretical methods allowed us to establish clear correlations between physicochemical and catalytic properties for the eCO₂RR of atomically dispersed MN_x centres (M = Mn, Fe, Co, Ni, and Cu). The ex situ characterization demonstrates the incorporation of the metal exclusively in MN_x centers over otherwise similar nitrogen-doped carbon matrices. Ex situ XPS and in situ XANES were applied to identify and monitor changes of the metal oxidation state at rest and under operating conditions, respectively. Mn and Co did not change oxidation state down to -1.0 V versus RHE, whereas Fe and Ni were partially reduced and Cu was largely reduced to a metallic state. Over these highly controlled systems, catalytic tests revealed a volcano-like dependency between eCO₂RR activity and the atomic number of the transition metal, with Fe and Co as the most active centers. In contrast, no clear trend was observable for the FE toward CO, for which Fe, Mn, and Ni-based materials showed the highest values (>80%). Computational models led to the identification of active centers and finally to descriptors rationalizing the distinct activity and selectivity patterns observed. To this end, different configurations were studied according to the speciation of the metal atom, its charge, spin, and coordination sphere. Among the best performing materials, Fe²⁺N₄-H₂O and Co²⁺N₄-H₂O centers were found as the most likely active catalytic centers at considered potentials (-0.5 and -0.6 V vs RHE), whereas Ni¹⁺N₄ was predicted as the most active Ni-based one. The Gibbs free energy change at the RDS accounted for the experimentally determined activity volcano, with Co-, Fe-, and Ni-based systems showing the best compromise between CO* and CO₂* binding energies. Regarding selectivity, the difference between the binding energies for CO₂* and H*, directly related to the number of d-antibonding electrons, described successfully the high selectivity observed on Fe- and

Ni catalysts and the more favored HER observed over the Co-based one. All in all, these results describe the influence of the particular metal species existing at relevant potentials on the catalytic activity and selectivity, laying down the path for a rational optimization of single-atom catalysts for eCO₂RR.

■ ASSOCIATED CONTENT

5 Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acscatal.9b02594.

Physical (XRD, (S)TEM, EDX, XAS, XPS, SEM, N₂ physisorption, and Raman) and electrochemical (CV, LSV) characterizations of MNC catalysts, the calculated electrochemical potential versus logarithm of the reaction rate with simplified Tafel laws, and the DFT calculated Gibbs free energy diagrams of eCO₂RR to CO and HER (PDF)

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Author Contributions

J.L. and F.J. designed, synthesized, and physically characterized the materials. T.S., A.J.M.F., and J.P.-R. performed the eCO₂RR electrochemical measurements and product analysis. P.P., Y.Z., R.G.-M., and N.L. performed the DFT simulations. F.K. conducted the (S)TEM and EDX characterizations. A.Z. conducted the XAS analysis and computation. K.A. and P.A. conducted the XPS measurement and analysis. J.L. and F.J. wrote and edited the paper with input from all authors. All authors read and approved the final version of the paper.

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Notes

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UNIVERSITAT ROVIRA I VIRGILI
THEORETICAL STUDIES OF SINGLE-SITE CATALYSTS FOR EFFICIENT ELECTROCHEMICAL CO₂ REDUCTION
Paulina Pršlja

Stability and Redispersion of Ni Nanoparticles Supported on N-Doped Carbons for the CO₂ Electrochemical Reduction

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1. INTRODUCTION

Single-metal nitrogen carbon materials (M–N–C) have been proposed as one of the most attractive catalysts for electrocatalytic CO₂ reduction reaction (eCO₂RR). Furthermore, with maximum atom efficiency and low coordination metal centers, they show high selectivity toward producing CO^{1–4} even with the competing hydrogen evolution reaction (HER). Combined products or syngas can be used to produce liquid hydrocarbons via Fischer–Tropsch synthesis. Moreover, the M–N–C materials can further reduce CO to hydrocarbons.^{3,5,6} The experimentally most studied example of these types of materials, even proposed to be industrially efficient, comes from Ni–N–C, which is highly selective to producing CO.^{7–20}

The efficiency of these materials can be controlled by synthetic and postsynthetic treatments. Their nitrogen content can be controlled in several ways, such as changing the doping temperature.²⁰ All these approaches change the nature of the cavities present both in terms of size and coordination environment. Therefore, we can classify the hosts according to these cavities as saturated and unsaturated N-doped carbon materials. Homogenous-like confinement of the metal in the cavities plays an important role in preventing the migration as the material goes through high-temperature carbonization.²¹ The vacancies can be synthetically tailored by employing MOF scaffolds,^{4,22} by oxidizing carbon materials and further doping with nitrogen to form undefined motifs,^{20,23} and by depositing porphyrin-like complexes.^{5,24–26} The ultimate (i) size and structure of the defect in which the metal is embedded; (ii) the nitrogen content which can be controlled with the doping temperature,²⁰ and (iii) the reaction conditions determine the nature of the metal sites, their coordination environment, and the size of the metal nanoparticles, thus affecting the electronic structure of the active metal. Targeting activity and selectivity require adapting the anchoring sites to the metal electronic properties so that they can synergistically act.

To understand the efficiency and selectivity of the materials, different types of catalysts were synthesized and tested under electrochemical eCO₂RR.²⁰ To be concrete, these catalysts include (i) Ni nanoparticles supported on carbon material, (ii) single metal atom in saturated N-doped carbon material, (iii) sample containing both particles and single atoms, and (iv) Ni nanoparticles supported on unsaturated N-doped carbon materials. Importantly, significant changes in efficiency and product distribution were observed when repeating the electrochemical reaction five times in a row. This effect was

particularly acute for the Ni-containing colloidal nanoparticle samples on the unsaturated N-doped materials. With the colloidal particle method, the size and particle loading can be controlled, and this preparation method is crucial for the stability of the catalyst and depends on the support. High efficiency can also be seen in the Ni single atoms induced by thermal atomization in N-doped carbon, but it does not occur over pure carbon support.¹⁹ Ni particles and Ni single active sites supported on carbon nitrides²⁷ featured high selectivity toward CO. These authors demonstrated that the materials can boost the conversion of CO₂ into CO.¹⁰ All these works highlight the role of the nature of the support and how its detailed characterization allows fine-tuning the chemical properties of these materials.

To shed light and rationalize such catalytic behavior, we turn to simulations at the atomic level to explore different N-doped carbon materials, both saturated and unsaturated, as reported in previous experiments.²⁰ Then we compute the activity of Ni catalysts for eCO₂RR, both as a nanoparticle and as single atoms embedded in different coordination environments. Furthermore, we report a model nanoparticle to link the possible reconstruction and redispersion under electrochemical conditions when supported on unsaturated N-doped carbon materials.

2. COMPUTATIONAL DETAILS

All simulations were performed using spin-polarized density functional theory (DFT) as implemented in the Vienna ab initio simulation package (VASP)^{28,29} using the GGA PBE-D3 density functional.^{30,31} Core electrons were described by projector augmented wave (PAW),^{29,32} while valence electrons were expanded by plane wave basis sets with a kinetic energy cutoff of at least 500 eV.

N-doped carbon materials are represented with a graphene layer expanded in a (6 × 6) supercell and replacing carbon atoms by nitrogen atoms to design saturated and unsaturated N-doped carbon models. These layers were interleaved by a 12 Å vacuum along the z direction. A nickel atom was placed in the cavity left on the carbon sheet. Structures were relaxed with

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a force threshold of 0.050 eV/Å. The Brillouin zone was sampled using a $(3 \times 3 \times 1)$ k -point mesh generated with the Monkhorst–Pack method.³³

Ni nanoparticles were represented using three metal surfaces with the lowest Miller indexes (and thus lowest surface energies) Ni(100), Ni(111), and Ni(211). High-quality surface-energy calculations ($(10 \times 10 \times 1)$ k -point mesh) were performed to build the nanoparticle structure using the Wulff model^{34,35} with VESTA.³⁶ The vacuum between slabs was larger than 10 Å, and the structures were relaxed with a force threshold of 0.025 eV/Å. They contained five layers, where the two uppermost were fully relaxed and the rest were fixed to the bulk structure.

CO can adsorb in very dense phases on Ni,^{37,38} and thus, different coverages were considered: from 1/12 to 12/12 ML for Ni(100) and Ni(211) and from 1/16 to 16/16 ML for Ni(111). The Brillouin zone was sampled using a $(3 \times 4 \times 1)$, $(3 \times 3 \times 1)$, and $(5 \times 3 \times 1)$ k -point meshes for Ni(100), Ni(111), and Ni(211) respectively. The Wulff construction was also employed on the CO-dense phases to explain the restructuring of the nanoparticles. Ni(CO)_{*x*} species can be formed in rich CO environments, which can then disperse and reconstruct the Ni nanoparticle. The Ni(211) surface edge has enough space to accommodate more than one CO molecule, allowing the formation of Ni(CO)₄ complexes. We thus calculated the formation and elimination of Ni(CO)₄ with different CO coverages on that surface, where the elimination of this species leaves a vacancy at the edge of the Ni(211) steps. The Ni(CO)₄ formation enthalpy computed with PBE starting from CO is more exothermic than that experimentally reported, and the implications for this deviation are evaluated in the text.

For all the systems described before, we obtained the energies of the relevant intermediates in the electrochemical environment under eCO₂RR conditions by considering the computational hydrogen electrode (CHE).^{39,40} The intermediates' energies were corrected with calculated zero point and entropic contributions. Notice that for the first step where CO₂ adsorbs on the metals, it drags charge from the system, thus the notation then implies that it is activated. Notice that the calculations were done in boxes with zero net charge. Gas-phase molecules used as a reference were calculated with PBE functional. Zero-point, entropic, and heat capacity contributions are shown in Table S1.⁴¹ Solvation effects were added to CO and COOH intermediates following the values employed in ref 41. We have tested and compared the results of our in-house implicit solvent model (VASP-MGCM)⁴² to the values in the literature, Table S5–S6. Solvation of N-doped carbon systems are known to depend on the N-concentration, and we have evaluated these effects in our case.⁴³ However, in our models the N-doping occurs close to a single point in the cavity, and thus, the effect by changing the stoichiometry of the cavity is reduced, for the NiN₃ and NiN₄ models –0.39 and –0.35 eV, respectively. Further details on thermodynamic corrections are discussed in the Supporting Information. All structures can be retrieved from the database ioChem-BD⁴⁴ via the following link,⁴⁵ which contains a unique and labeled list of the DFT calculations in this project.

3. RESULTS AND DISCUSSION

3.1. Computational Characterization of the Materials.

To understand the selectivity and activity of N-doped or N-free carbon materials, we modeled them under reaction

conditions. The DFT simulations were performed on several N-doped defects where the Ni atoms can be anchored and Ni nanoparticles can be supported, to represent the species characterized in experiments.²⁰ In Figure 1, we show three

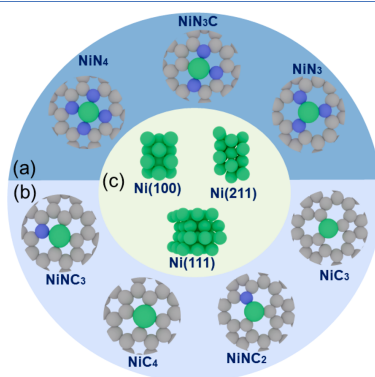


Figure 1. Computational models for different systems. (a) Saturated N-doped carbon material. (b) Unsaturated N-doped or N-free carbon material. (c) Metal surfaces. C in gray spheres, N in blue, and Ni in green.

families of models: (a) saturated N-doped carbon materials (NiN_{*x*}C) representing NiN₄ and similar undefined N-doped materials made at low doping temperature; (b) unsaturated N-doped (NiNC_{*x*}) and N-free materials (NiC_{*x*}) typically occurring for carbons at high doping temperatures, when the nitrogen content decreases; and (c) Ni metal surfaces composing the supported nanoparticles on the unsaturated N-doped supports (NiN_{*NP*}–NC_{*x*}). This approximation allows us to provide a holistic chain of models that accounts for the complexity of the N-doped carbon support and the metallic nature of the nanoparticle. The ultimate aim is to show the role of the speciation (charge, oxidation state, and coordination environment) on the electrochemical activity and the dynamic behavior or the different synthetic preparations. We simulated EXAFS spectra for all catalysts (Figure S1), where NiN_{*x*}C (saturated N-doped materials) agreed well with experimental data.²⁰ DFT structural parameters are presented in Table S4.

3.2. Activity and Selectivity of the Materials.

The activity and selectivity of the modeled catalysts can be discussed in terms of the corresponding Gibbs free-energy profiles, presented in Figure 2. The electrochemical activity was analyzed by calculating the binding energies of the intermediates according to the most agreed mechanism for CO₂ to CO reduction (Figure 2a–c). The process requires only a two-electron transfer, and the steps are (i) CO₂ + e[–] → *CO₂, (ii) *CO₂ + H⁺ → *COOH (also, the first two steps in the literature are combined into: CO₂ + H⁺ + e[–] → *COOH), (iii) *COOH + H⁺ + e[–] → *CO + H₂O and (iv) *CO → CO_(g) + *. Notice that once CO₂ is adsorbed, activated, and polarized with partial charge transfer (see Table S10) from the host material, step (i) does not correspond to any extra electron added to the simulation box.

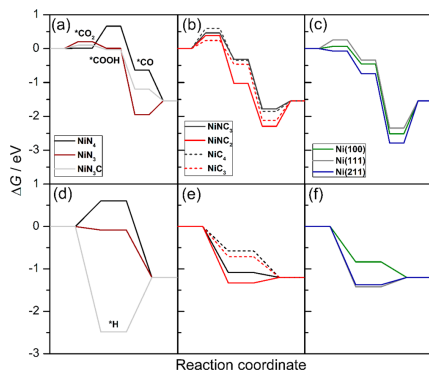


Figure 2. Gibbs free-energy profiles for (a, b, and c) eCO₂RR and (d, e, and f) HER over different model sites at $U = -0.6$ V vs RHE (for reference the experimental pH was 7). (a,d) Saturated nitrogen-doped carbon materials. (b,e) Unsaturated or free N-doped carbon materials. (c,f) Metal surfaces.

In agreement with previous results, we can see that the highest activity toward CO is shown by unsaturated N-doped materials⁷ and metal surfaces, as they bind CO strongly. Regarding selectivity, it can be discussed by comparison to the competing HER (Figure 2d–f). Saturated N-doped materials reach the highest selectivity toward CO, where NiN₄ and NiN₃ are representative models, while NiN₂C model binds *H strongly. Therefore, based on the experimental selectivity³⁰ of single-atom catalysts, the NiN₂C model is not representative. For a catalyst to be selective toward CO, the binding of *COOH and *CO₂ intermediates need to be strong to reduce the onset potential. Also, selectivity can be compared with *H binding energy and CO desorption energy. The correlation between the adsorption energies of *COOH and *H is shown in Figure 3a. We can see that NiN₄ and NiN₃ bind *H weakly, and therefore, saturated N-doped materials are selective. NiN₄ binds *CO₂ and *COOH weakly, which means it has a higher onset potential, while the opposite trend is seen for NiN₃. The computed onset potentials are presented in Table S11, where

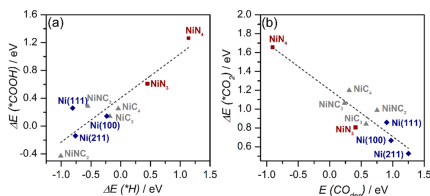


Figure 3. Linear scaling relationships. (a) Linear dependence for the COOH binding energy scaling as a function of the adsorption of a H atom. (b) CO₂ binding energy and corresponding linear-scaling relationships for desorption of the CO molecule. Saturated N-doped carbon materials are represented with red symbols, unsaturated N-doped carbon materials with gray symbols, and metal surfaces with blue symbols.

NiN₃ has a lower onset potential than NiN₄ because it binds *COOH stronger. Ni(211) and NiN₂C bind strongly *H and *COOH, which means they will not be selective toward CO. Ni(100) would be the best candidate among the metal surfaces as it binds *H weakly and *COOH relatively strongly. Figure 3b represents a correlation between adsorption energies of *CO₂ and desorption of CO, where CO₂ activation can relate to Bader charges of the adsorbed CO₂ molecule (see Table S10). NiN₃, NiN₂C, and Ni(100) bind *CO₂ relatively strongly, where Ni(100) binds *CO strongly (see Figure S3), which means the rate of CO formation would be low. The strong adsorption energy of *H for Ni(100) and NiN₃ can be related to low selectivity, where HER and eCO₂RR compete. The most interesting system is NiN₃ because of the optimal *CO binding and the weakly *H binding. From an electronic structure point of view, the Ni atom in the NiN₄ square planar confinement is in the Ni²⁺ oxidation state, while that in the NiN₃ trigonal planar configuration is in the Ni⁺ oxidation state (see Figure S4). This special role of Ni⁺ is long known in organometallic chemistry⁴⁶ and even in thermal catalysis.⁴⁷ The change in the oxidation state has an impact on the geometries of the activated CO₂ intermediate as both the O–C–O angle and Ni–C(CO₂) distance indicate the degree of activation of the CO₂ molecule in the N-doped saturated materials. In Figure S5, the linear scaling relationship of CO₂ geometry and adsorption of *CO shows that the activation in NiN₃ correlates with the medium binding of *CO, while the lowest activation appears in NiN₄ with a C–O–O angle of 170.2°. All these fingerprints point out the relevance of understanding the organometallic structures when trying to identify the most suitable defects for a given metal and the importance of the charge state for the encapsulated metal ions.

4. RECONSTRUCTION OF NI NPS ON N-DOPED CARBON MATERIALS

Electrocatalysts can be affected by the presence of different variables. In our case, the reconstruction of Ni NPs and the disintegration of Ni(CO)₂ complexes change the activity and selectivity. Once the reconstruction happens experimentally, the partial current density of H₂ drastically drops compared with that of CO, thus favoring the disintegration process. Such dynamic behavior of electrocatalysts is crucial for their performance, and thus, we here address these aspects by simplified models.

To investigate the role of the reaction products in the stability of the Ni NPs, we model a nanoparticle-based on the lowest-energy surfaces through a Wulff construction. This is straightforward for metal nanoparticles and can be extended when adsorbates, like CO, are present.^{47,48} In addition, we apply the atomistic version of Ostwald ripening theory^{49,50} by the formation of Ni(CO)₂ complexes to show when the disintegration of nanoparticles can occur on such N-doped carbon materials.

4.1. Nanoparticle Shape. To address the role of nanoparticles, we first focus on their structure. The size of the synthesized Ni NPs was 2.5 to 7 nm, and thus, they can be represented via the Wulff construction (Figure 4), which corresponds to the equilibrium structure for a medium-to-large nanoparticle and is obtained from the surface energies. We calculated the surface energies for three orientations Ni(111), Ni(100), and Ni(211), and, employing the Wulff construction, we then extracted the corresponding areas for each facet. The average surface energy for the nanoparticle is 0.126 eV/Å. All

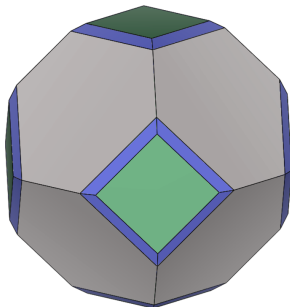


Figure 4. Wulff construction based on the surface energies calculated by DFT. Surface legend: Ni(100) in green, Ni(111) in gray, and Ni(211) in blue.

surface energies and ratios of the corresponding facets are presented in Table S17.

To assess the change in the nanoparticle that is induced by the environment, we consider CO adsorption on all three facets separately.³⁷ We calculated average binding energies E_{CO}^{f} for a range of coverages 1/12 to 12/12 ML for Ni(100) and Ni(211) facets, and Ni(111) from 1/16 to 16/16 ML (Figure S7). In all cases, at low to medium coverages, all the modified surface energies (γ_s') are smaller, but the reduction in surface energies depends on the particular facet orientation. The detailed balance between adsorption sites and CO–CO repulsion contributes to the surface energy in a nonlinear way, as to achieve higher CO densities the molecules need to go to rather unstable sites. At 0.25 ML, we are left with Ni(100) and Ni(111) surfaces, while at full coverage of 1 ML the extent of Ni(211) facet in the nanoparticle increases and Ni(111) reconstructs (see Figure S6). In summary, the average surface energy for the decorated nanoparticle (γ') at CO coverage 0.25 ML is 0.095 eV/Å², which is 0.031 eV/Å² less than for the bare one (Table S16). Therefore, the nanoparticle under product-rich conditions has more surface Ni atoms with less Ni–Ni average coordination. This could be the first step toward disintegration. In addition, this high product coverage influences the selectivity as it limits the HER reaction. The effect can be explained with the Gibbs free energies of the *H intermediate for the clean and on the CO covered ($\theta = 0.92$ ML) Ni surfaces (see Figure S10). Once Ni surfaces are covered with CO, *H binds more weakly, thus lowering the selectivity for HER (see Figure S11). This agrees with the experimental observations showing a drastic drop of H₂ partial current density. Thus, selectivity is triggered by minimizing the HER path.

4.2. Gibbs Free Energy of Disintegration Ni(CO)_x Complexes on N-Doped Carbon Materials. We applied the energetic formalism from Ouyang and co-workers⁵⁰ to evaluate whether the disintegration of Ni(CO)_x from Ni NPs is possible on N-doped carbon materials. We followed the criteria of reactant induced dynamics, where feasibility can be explained by the Gibbs free energy of disintegration ($\Delta G_{\text{NP}}^{\text{dis}}$). First, we compute the formation energy of isolated metal adatom (E_{Ni}^{f}) with respect to Ni bulk reference supported on the different carbon defects (Table S18). From these results,

we can see that formation energies are exothermic, as Ni atoms are anchored to a defect. Second, we calculate the Gibbs free energy for the CO-decorated nanoparticle, $\Delta G_{\text{NP}}^{\text{dis}}$, following the equation:

$$\Delta G_{\text{NP}}^{\text{dis}}(R, T, p) = E_{\text{surf-nCO}}^{\text{f}} - n \times \Delta \mu_{\text{CO}}(T, p) - \Delta E'_{\text{NP}}(R) - TS \quad (1)$$

where $E_{\text{surf-nCO}}^{\text{f}}$ is the energy of the Ni atom coordinated to the N-doped carbon substrate and n remaining CO molecules, $\Delta E'_{\text{NP}}(R)$ is the energy of the CO-covered nanoparticle with radius R , and $\Delta \mu_{\text{CO}}(T, p)$ is the excess chemical potential of CO (T, p). In our case, the configurational entropy S was not included because the number of cavities is difficult to assess; however, it will favor dispersed species. Therefore, the values obtained here for the transition correspond to a maximum threshold. The excess chemical potential was set to $\Delta \mu_{\text{CO}}(T, p) = -0.76$ eV,⁵¹ corresponding to experimental conditions of 300 K and 0.10 mbar. The binding energy of nCO ($E_{\text{ICO}}^{\text{f}}$) needs to be negative and lower than $\Delta \mu_{\text{CO}}(T, p) = -0.76$ eV. When both requirements are fulfilled, nanoparticle disintegration is thermodynamically possible.

The nature of the cavities in the support provides a large versatility regarding the stability against disintegration. In our investigations, we have considered all 7 types of defects in Figure 1. The anchored Ni atom in saturated N-doped models is coordinated either to pyridinic or pyrrolic N atoms, which are strong ligands able to easily trap metal atoms. This is in line with experiments where they do not observe nanoparticles on those materials. We observe an opposite trend for the unsaturated N-doped materials, where Ni NPs can redisperse into single-atom cavities. The reason for that is the decreased N-content and increased C-content in the materials, which makes Ni dynamic and prone to aggregate and disintegrate in the presence of dense CO phases and electrochemical conditions. Therefore, single Ni metal atoms with little aggregation are expected in agreement with experiments.¹⁰ Finally, we show that the redispersion of Ni NPs on the NC₂ defect (unsaturated N-doped model) is thermodynamically possible. NiNC₂ is also one of the most active and nonselective eCO₂RR configurations, with strong bindings for *CO, *COOH, and *H intermediates. Once Ni NPs are supported on NC₂, they can disintegrate, as the experimentally partial current density drops drastically compared with CO partial current densities. That is how the combination of the two results (activity and stability against dispersion) can support the fact that enhanced Faradaic efficiency was observed when the reaction was performed five times in a row for nanoparticles supported on unsaturated N-doped carbon materials.²⁰ Redispersion is thus possible on unsaturated materials; however, depending on the nature of the carrier that would imply a different performance. As the performance of the single atoms depends on the properties of the cavity, redispersion would only be more favorable provided that the right cavity exists, and it is occupied. The electrochemical cycling from the nanoparticles does enhance the selectivity in unsaturated hosts but not up to the levels found for the saturated materials.

To investigate the size (R) dependence of the disintegration induced by CO, we have used Ni_{NP}-NC₂ as cavity and have calculated the $\Delta G_{\text{NP}}^{\text{dis}}(R, T, p)$ considering the formation of Ni(CO)₂ sitting at these pockets at experimental temperature and pressure. Figure 5 presents the $\Delta G_{\text{NP}}^{\text{dis}}$ corresponding to

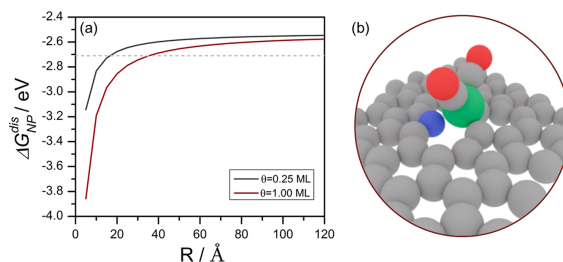


Figure 5. (a) Size dependence of the Gibbs free energy (ΔG_{NP}^{dis}) of $Ni(CO)_2$, when Ni_{NP} is supported on NC_2 at 300 K and 0.1 mbar. The horizontal dashed line indicates the limit for stability. (b) Schematic representation of $Ni(CO)_2$ bound to the NC_2 cavity.

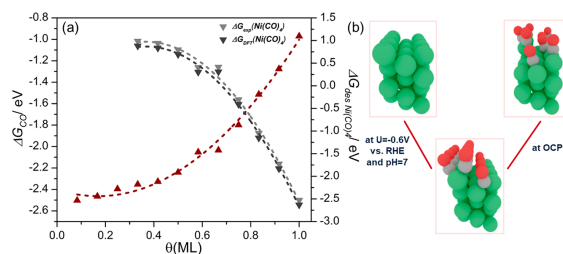


Figure 6. Process of $Ni(211)$ metal surface reconstruction. (a) CO coverage dependence of the average Gibbs free energy (red curve) of CO adsorption at $U = -0.6$ V vs RHE ($pH = 7$) and (gray curves) desorption of $Ni(CO)_x$. (b) Schematic process of clean $Ni(211)$ metal surface fully adsorbing CO at electrochemical conditions. When the reaction is stopped at the open circuit potential (OCP) to assess the material recyclability, $Ni(CO)_x$ -like complexes can desorb.

$Ni(CO)_2$ versus the radius R of a Ni nanoparticle at two different coverages, 0.25 and 1 ML (nonwetting conditions for the nanoparticle were considered, thus the Young angle was set to 180°). We found that ΔG_{NP}^{dis} decreases at smaller R and crosses the reference level (-2.7 eV, energy of an isolated Ni atom in the cavity) at about 40 Å for 1 ML and 20 Å for 0.25 ML. The lower value found for lower coverage is a direct consequence of the simplicity of the model and shall be taken with caution as kinetic constraints related to the need of concentrating CO molecules to extract the surface atoms could play an important role. At high coverage, our results imply that Ni NPs smaller than 4 nm will disintegrate into $Ni(CO)_2$ when supported on the NC_2 at 300 K and 0.1 mbar.

4.3. Mechanism for Ni Nanoparticle Reconstruction.

The dense layers that we have identified in Section 4.1 show that after eCO_2RR , CO extensively covers the surface and over long times has the potential to increase the fraction of Ni atoms in low-coordinated sites as stepped $Ni(211)$ surfaces. Moreover, the thermodynamic models in Section 4.2 indicate that below a certain size, $Ni(CO)_x$ clusters can be stable in some particular motifs of the N-doped carbons. To merge these two scenarios, a viable route for decomposition needs to be described. Ni can form volatile species with CO in the form of $Ni(CO)_4$ that can be soluble under the reaction conditions. We have calculated the energy needed for the formation of these species.

On nanoparticles, CO_2 is reduced to CO at -0.6 V vs RHE and pH 7, and the resulting CO drives the reconstruction

toward $Ni(211)$. Once the reaction is stopped at open circuit potential (OCP), $Ni(CO)_4$ species can be formed as the CO coverage is very high. Thermodynamically, the formation energy of a metal vacancy (with respect to Ni bulk) in $Ni(100)$, $Ni(111)$, and $Ni(211)$ surfaces are 0.59, 1.04, and 0.33 eV, respectively. Strong kinetic limitations can likely appear due to geometric constraints on $Ni(111)$ and $Ni(100)$, but for the fully CO adsorbed $Ni(211)$ surface, the step edges can present very high CO concentration. Thus, $Ni(CO)_4$ can be formed in this face with concomitant elimination of Ni atoms. Figure 6a presents the average adsorption energy for CO at different coverages (red curve) on the $Ni(211)$ facet and it is compared to the formation-elimination of the $Ni(CO)_4$ complex (gray curves). The crossing between these two lines can be interpreted as the coverage at which the elimination of Ni atoms in the form of the species becomes viable. Therefore, high CO coverages would affect the stability of the edge sites starting around 0.78 ML coverage (when using DFT or experimental value for $Ni(CO)_4$), promoting the formation of the species that can then redisperse on the cavities of the N-doped carbon. The process of reconstruction is sketched in Figure 6b. All the data is presented in Tables S21–23. Once the reaction is under electrochemical potential again, these complexes react with nonoccupied defects on the unsaturated N-doped carbon ($Ni_{NP}-NC_2$) and effectively redisperse. This process could potentially lead to higher activity, as shown in the experiments.²⁰

5. CONCLUSIONS

Electrocatalysts based on Ni supported on N-doped carbon materials have been computationally analyzed for the reduction of CO₂, considering both nanoparticles and single atoms. The optimal activity and selectivity are found for the NiN₃ model, in which Ni is in Ni¹⁺ oxidation state. NiN₃ exhibits theoretically low onset potential in agreement with experiments and a high rate toward CO. However, rather, saturated N-doped materials are more selective toward CO formation. As single atoms can be prepared from nanoparticles, we show through the Ostwald formalism that disintegration of Ni(CO)_x can happen if Ni nanoparticles are supported on N-doped unsaturated support such as NC₂. We present a mechanism of Ni(211) surface reconstruction that is possible only at high CO coverages around 8/12 ML by the formation of Ni(CO)₄ species, which then redisperses into active single atoms. The complete understanding of the Ni–N–C materials points toward a very dynamic behavior where considering the full complexity of the material in terms of Ni nuclearity and environment is needed. This is likely a common feature for such type of electrocatalysts.

■ ASSOCIATED CONTENT

Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/acscatal.0c01909>.

Additional notes on the energy corrections and the modeling for the Ni nanoparticles for this article (PDF)

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Notes

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